

ROADMAP • OPEN ACCESS

## 2024 roadmap on magnetic microscopy techniques and their applications in materials science

To cite this article: D V Christensen *et al* 2024 *J. Phys. Mater.* **7** 032501

View the [article online](#) for updates and enhancements.

You may also like

- [LOCALIZATION AND BROADBAND FOLLOW-UP OF THE GRAVITATIONAL-WAVE TRANSIENT GW150914](#)  
B. P. Abbott, R. Abbott, T. D. Abbott et al.
- [Multi-messenger Observations of a Binary Neutron Star Merger](#)  
B. P. Abbott, R. Abbott, T. D. Abbott et al.
- [Fast  \$b\$ -tagging at the high-level trigger of the ATLAS experiment in LHC Run 3](#)  
G. Aad, B. Abbott, K. Abeling et al.

**PRIME**  
PACIFIC RIM MEETING  
ON ELECTROCHEMICAL  
AND SOLID STATE SCIENCE

**HONOLULU, HI**  
October 6-11, 2024

*Joint International Meeting of*  
The Electrochemical Society of Japan (ECSJ)  
The Korean Electrochemical Society (KECS)  
The Electrochemical Society (ECS)

Early Registration Deadline:  
**September 3, 2024**

**MAKE YOUR PLANS NOW!**

## Journal of Physics: Materials



## ROADMAP

## OPEN ACCESS

## 2024 roadmap on magnetic microscopy techniques and their applications in materials science

## RECEIVED

14 July 2023

## REVISED

18 December 2023

## ACCEPTED FOR PUBLICATION

8 March 2024

## PUBLISHED

13 June 2024

Original content from this work may be used under the terms of the [Creative Commons Attribution 4.0 licence](https://creativecommons.org/licenses/by/4.0/).

Any further distribution of this work must maintain attribution to the author(s) and the title of the work, journal citation and DOI.



D V Christensen<sup>1,\*</sup> , U Staub<sup>2</sup> , T R Devidas<sup>3</sup> , B Kalisky<sup>3</sup> , K C Nowack<sup>4,5</sup> , J L Webb<sup>6</sup> , U L Andersen<sup>6</sup> , A Huck<sup>6</sup> , D A Broadway<sup>7,8</sup> , K Wagner<sup>7</sup> , P Maletinsky<sup>7</sup> , T van der Sar<sup>9</sup> , C R Du<sup>10</sup> , A Yacoby<sup>11</sup> , D Collomb<sup>12</sup> , S Bending<sup>13</sup> , A Oral<sup>14</sup> , H J Hug<sup>7,15</sup> , A-O Mandru<sup>15</sup> , V Neu<sup>16</sup> , H W Schumacher<sup>17</sup> , S Sievers<sup>17</sup> , H Saito<sup>18</sup> , A A Khajetoorians<sup>19</sup> , N Hauptmann<sup>19</sup> , S Baumann<sup>20</sup> , A Eichler<sup>21</sup> , C L Degen<sup>21</sup> , J McCord<sup>22,23</sup> , M Vogel<sup>22,23</sup> , M Fiebig<sup>24</sup> , P Fischer<sup>25,26</sup> , A Hierro-Rodriguez<sup>27,28</sup> , S Finizio<sup>2</sup> , S S Dhesi<sup>29</sup> , C Donnelly<sup>30,31</sup> , F Büttner<sup>32,33</sup> , O Kfir<sup>34</sup> , W Hu<sup>35</sup> , S Zayko<sup>36</sup> , S Eisebitt<sup>37,38</sup> , B Pfau<sup>37</sup> , R Frömter<sup>39</sup> , M Kläui<sup>39</sup> , F S Yasin<sup>40</sup> , B J McMorran<sup>41</sup> , S Seki<sup>42</sup> , X Yu<sup>40</sup> , A Lubk<sup>16,43</sup> , D Wolf<sup>16</sup> , N Pryds<sup>1</sup> , D Makarov<sup>44</sup> and M Poggio<sup>7</sup>

- <sup>1</sup> Department of Energy Conversion and Storage, Technical University of Denmark, DK-2800 Kgs Lyngby, Denmark
- <sup>2</sup> Swiss Light Source, Paul Scherrer Institute, 5232 Villigen PSI, Switzerland
- <sup>3</sup> Department of Physics and Institute of Nanotechnology and Advanced Materials, Bar-Ilan University, Ramat Gan 5290002, Israel
- <sup>4</sup> Laboratory of Atomic and Solid State Physics, Cornell University, Ithaca, NY 14853, United States of America
- <sup>5</sup> Kavli Institute at Cornell for Nanoscale Science, Ithaca, NY 14853, United States of America
- <sup>6</sup> Centre for Macroscopic Quantum States (bigQ), Department of Physics, Technical University of Denmark, 2800 Kongens Lyngby, Denmark
- <sup>7</sup> Department of Physics and Swiss Nanoscience Institute, University of Basel, Klingelbergstrasse 82, 4056 Basel, Switzerland
- <sup>8</sup> School of Science, RMIT University, Melbourne, Victoria, Australia
- <sup>9</sup> Department of Quantum Nanoscience, Kavli Institute of Nanoscience, Delft University of Technology, Lorentzweg 1, 2628 CJ Delft, The Netherlands
- <sup>10</sup> School of Physics, Georgia Institute of Technology, Atlanta, GA 30332, United States of America
- <sup>11</sup> Department of Physics, Harvard University, 17 Oxford Street, Cambridge, MA 02138, United States of America
- <sup>12</sup> Quantum Brilliance GmbH, Tullastraße 72, Freiburg im Breisgau, Germany
- <sup>13</sup> Department of Physics, University of Bath, Bath BA2 7AY, United Kingdom
- <sup>14</sup> NanoMagnetics Instruments Ltd, Suite 290, 266 Banbury Road, Oxford OX2 7DL, United Kingdom
- <sup>15</sup> Empa, Materials Science and Technology, Ueberlandstrasse 129, 8600 Dübendorf, Switzerland
- <sup>16</sup> Leibniz Institute for Solid State and Materials Research, Helmholtzstrasse 20, D-01069 Dresden, Germany
- <sup>17</sup> Physikalisch-Technische Bundesanstalt, 38116 Braunschweig, Germany
- <sup>18</sup> Department of Mathematical Science and Electrical-Electronic-Computer Engineering, Graduate School of Engineering Science, Akita University, Akita, Japan
- <sup>19</sup> Scanning Probe Microscopy Department, Institute for Molecules and Materials, Radboud University, Nijmegen, The Netherlands
- <sup>20</sup> University of Stuttgart, Institute for Functional Matter and Quantum Technologies, Stuttgart, Germany
- <sup>21</sup> Laboratory for Solid State Physics, ETH Zürich, 8093 Zürich, Switzerland
- <sup>22</sup> Department for Materials Science, Kiel University, Kiel, Germany
- <sup>23</sup> Kiel Nano, Surface and Interface Science (KiNSIS), Kiel University, Kiel, Germany
- <sup>24</sup> Department of Materials, ETH Zurich, Zurich, Switzerland
- <sup>25</sup> Materials Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States of America
- <sup>26</sup> Department of Physics, University of California Santa Cruz, Santa Cruz, CA 95064, United States of America
- <sup>27</sup> Departamento de Física, Universidad de Oviedo, 33007 Oviedo, Spain
- <sup>28</sup> CINN (CSIC-Universidad de Oviedo), 33940 El Entrego, Spain
- <sup>29</sup> Diamond Light Source, Chilton, Didcot, Oxfordshire OX11 0DE, United Kingdom
- <sup>30</sup> Max Planck Institute for Chemical Physics of Solids, Dresden, Germany
- <sup>31</sup> International Institute for Sustainability with Knotted Chiral Meta Matter (WPI-SKCM2), Hiroshima University, Hiroshima 739-8526, Japan
- <sup>32</sup> Helmholtz-Zentrum für Materialien und Energie, Berlin, Germany
- <sup>33</sup> University of Augsburg, Augsburg, Germany
- <sup>34</sup> School of Electrical Engineering, Tel Aviv University, Tel Aviv 69978, Israel
- <sup>35</sup> National Synchrotron Light Source II, Brookhaven National Laboratory, Upton, NY 11973, United States of America
- <sup>36</sup> Max Planck Institute for Multidisciplinary Sciences, Göttingen, Germany
- <sup>37</sup> Max Born Institute, Max-Born-Str. 2a, 12489 Berlin, Germany
- <sup>38</sup> Institute of Optics and Atomic Physics, Technische Universität Berlin, Straße des 17. Juni 135, 10623 Berlin, Germany
- <sup>39</sup> Institute of Physics, Johannes Gutenberg University Mainz, 55099 Mainz, Germany
- <sup>40</sup> RIKEN Center for Emergent Matter Science (CEMS), Wako, Japan
- <sup>41</sup> Department of Physics, University of Oregon, Eugene, OR, United States of America
- <sup>42</sup> Department of Applied Physics and Institute of Engineering Innovation, University of Tokyo, Tokyo, Japan
- <sup>43</sup> Institute of Solid State and Materials Physics, TU Dresden, 01062 Dresden, Germany
- <sup>44</sup> Helmholtz-Zentrum Dresden-Rossendorf e.V., Institute of Ion Beam Physics and Materials Research, 01328 Dresden, Germany

\* Author to whom any correspondence should be addressed.

E-mail: [dechr@dtu.dk](mailto:dechr@dtu.dk)

**Keywords:** magnetic microscopy, magnetism, electron transport, spin dynamics, 2D materials, multiferroics, superconductors

---

## Abstract

Considering the growing interest in magnetic materials for unconventional computing, data storage, and sensor applications, there is active research not only on material synthesis but also characterisation of their properties. In addition to structural and integral magnetic characterisations, imaging of magnetisation patterns, current distributions and magnetic fields at nano- and microscale is of major importance to understand the material responses and qualify them for specific applications. In this roadmap, we aim to cover a broad portfolio of techniques to perform nano- and microscale magnetic imaging using superconducting quantum interference devices, spin centre and Hall effect magnetometries, scanning probe microscopies, x-ray- and electron-based methods as well as magneto-optics and nanoscale magnetic resonance imaging. The roadmap is aimed as a single access point of information for experts in the field as well as the young generation of students outlining prospects of the development of magnetic imaging technologies for the upcoming decade with a focus on physics, materials science, and chemistry of planar, three-dimensional and geometrically curved objects of different material classes including two-dimensional materials, complex oxides, semi-metals, multiferroics, skyrmions, antiferromagnets, frustrated magnets, magnetic molecules/nanoparticles, ionic conductors, superconductors, spintronic and spinorbitronic materials.

---

## Contents

Introduction	4
1. Scanning SQUID microscopy	7
2. NV magnetometry	13
3. Scanning Hall microscopy	24
4. Magnetic force microscopy	28
5. Spin sensing on the atomic scale with STM and non-contact AFM	34
6. Nanoscale magnetic resonance imaging	38
7. Magneto-optical imaging	41
8. X-ray techniques	48
9. Scanning Electron Microscopy with Polarization Analysis	60
10. Transmission electron microscopy techniques	63
Data availability statement	71
References	72



## Introduction

*Dennis V Christensen*<sup>1</sup>, *Urs Staub*<sup>2</sup>, *Nini Pryds*<sup>1</sup>, *Denys Makarov*<sup>3</sup> and *Martino Poggio*<sup>4</sup>

<sup>1</sup> Department of Energy Conversion and Storage, Technical University of Denmark, DK-2800 Kgs Lyngby, Denmark

<sup>2</sup> Swiss Light Source, Paul Scherrer Institute, 5232 Villigen PSI, Switzerland

<sup>3</sup> Helmholtz-Zentrum Dresden-Rossendorf e.V., Institute of Ion Beam Physics and Materials Research, 01328 Dresden, Germany

<sup>4</sup> Department of Physics and Swiss Nanoscience Institute, University of Basel, Klingelbergstrasse 82, 4056 Basel, Switzerland

Today's devices rely on the flow of charge and the magnetisation of materials for sensing, carrying out computations, and storing data. The continued pressure to make these devices smaller, more efficient, and more powerful is pushing researchers to explore the use of spin currents, quantum transport, nanoscale spin configurations, and various types of quantum bits. Here, magnetic microscopy techniques play a vital role in enabling researchers to visualise, manipulate, and understand magnetic and transport phenomena at the micro- and nanoscale, and eventually optimise materials and devices for specific applications. The broad portfolio of magnetic microscopy techniques allows for spatially resolving a plethora of functional properties, including imaging the magnetic landscape, visualising complex current distributions, mapping the magnetic susceptibility, and detecting ultra-fast spin dynamics. Key applications in material science range from quantitatively mapping the flow in low-dimensional material systems, spatially resolving the superconducting fluid density, and disclosing the dynamics of spin-waves and spin textures. By utilising the properties of magnetic fields to traverse through solid materials, several magnetic microscopy techniques offer the possibility to spatially resolve processes buried underneath protective layers or at the interface of heterostructures.

In this roadmap, experts in diverse areas of magnetic microscopy techniques have been invited to provide their views on the current capabilities of the techniques, forecast future technical advances, and assess how such advances may impact material science. Although magnetic microscopy techniques find use in a broad range of fields, this roadmap focuses on their use within material science and physics. The aim is to provide a single access point comprising forward-looking opinions on selected techniques rather than presenting a comprehensive account of all magnetic microscopy techniques. We hope that the roadmap will serve as a valuable resource for experts and students in magnetic imaging techniques and functional materials, enterprises developing or leveraging on the methods, readers from outside the field as well as those entering the field.

The roadmap divides the magnetic microscopy techniques into three major fields, namely techniques based on scanning probe microscopy (SPM), light, and electrons:

### **Magnetic microscopy based on scanning probes**

Scanning tunnelling microscopy (STM), introduced in 1981, produced the first real-space images of individual atoms on a surface. A boom in SPM followed this, leading to various probes now being routinely used to form diverse images of surfaces by rastering the physical probe in close proximity to a sample while making a measurement at each position. In the area of magnetic imaging, some of the techniques combining the highest spatial resolution with the highest sensitivity are based on scanning probes. These include magnetic force microscopy (MFM), scanning superconducting quantum interference device (SQUID) microscopy, scanning Hall microscopy (SHM), scanning nitrogen-vacancy (NV) centre microscopy, and spin-polarised STM (SP-STM). These five techniques are covered in the subsequent sections of this roadmap. The techniques work under a wide variety of different conditions, and each is best suited to different types of samples. All are restricted to mapping contrast associated with the magnetic field or magnetisation at the surface of a sample. As physically scanning a probe across a sample is typically slow, these techniques are best suited to 100 nm to 100  $\mu\text{m}$  fields of view (FOVs) with resolutions down to the atomic scale in some cases. The spatial resolution is ultimately limited by a combination of the size of the scanning probe tip or sensor and the proximity of the latter to the sample surface. The high sensitivity and excellent spatial resolution enable, e.g. real-space imaging of uncompensated magnetisation at the surface of antiferromagnets and visualisation of the current distribution in the ballistic, hydrodynamic, or superconducting transport regimes. In addition, scanning probes are a convenient choice for measuring weakly magnetic materials with low magnetisation volumes and buried magnetic textures, such as encapsulated two-dimensional (2D) magnets.

While there is an astounding number of different types of SPMs, the list is bound to grow in the coming years. The development of ever more advanced fabrication techniques, including focused ion beam milling and advanced lithography, allows for patterning ever more complex and compact sensing devices at the apex of scanning probes. These include devices such as the sub-micron-scale graphene Hall bars, robust nanometre-scale SQUIDs, and sharp diamond tips with stable, ultra-shallow NVs, as discussed later in this roadmap. This is complemented by the development of multimodal scanning probes where imaging of magnetic susceptibility and spontaneous magnetic order can be combined with the mapping of temperature profiles, electric fields, and current distributions. The new generation of scanning probes promises to enable new types of contrast, including magnetic contrast, that could help provide microscopic information about length scales, inhomogeneity, and interactions in emerging and poorly understood magnetic materials.

### Light-based magnetic microscopy

Conventional imaging based on magneto-optics (MOs) using Kerr or Faraday effects provides easy access to magnetic textures with a spatial resolution of hundreds of nanometres and femtosecond (fs) temporal resolution in stroboscopic measurements. Current activities on the field- or current-driven evolution of magnetic textures in not only planar but also 3D magnetic architectures call for major improvements of the MO microscopies towards enhancement of contrast enabled by plasmon filtering and on-the-fly image analysis using graphics processing unit (GPU) computing assisted by machine learning methods. The improvement in the signal-to-noise ratio (SNR) should allow for single-shot magnetisation dynamics studies of non-repetitive events, which are out of reach at present at short time scales. It is envisioned that combining MO with near-field microscopies will allow for routine measurements at sub-100 nm spatial resolution. These features will broaden the use of MO techniques for magnetic topological materials like Weyl semimetals, collinear and non-collinear antiferromagnets, as well as magnetoelectrics.

Frequency doubling of light in matter, termed ‘second harmonic generation’ (SHG), is a powerful method for studying ordered states of matter based on the symmetry change accompanying the emergence of the order. The greatest potential of SHG microscopy is in the investigation of states with compensated magnetisation, like antiferromagnets, states with coexisting magnetic and electric order, termed multiferroics, and the spatially resolved imaging of the associated domains. Magnetisation processes are resolved with sub-picosecond (ps) accuracy, which is particularly useful in developing new magnetoelectronic or spintronic functionalities and devices. Based on the current development of SHG equipment, this technique could become a standard component of sample preparation equipment like deposition chambers. These so-called *in situ* SHG equipment will provide real-time feedback on growth conditions of magnetically and electrically active materials to facilitate the discovery or optimisation of multiferroic and magnetoelectric materials for low-energy storage and computing devices. It is foreseeable that SHG and other nonlinear optical methods will be extended to terahertz (THz) frequencies, opening possibilities to probe elementary excitations in materials resonantly and thus very directly and effectively.

X-rays are powerful probes for magnetism with the strength of being chemically selective to the ions carrying the magnetic moment. This leads to additional sensitivity in both x-ray scattering and absorption processes, allowing the study of small amounts of materials. In addition, it enables very high spatial resolution down to the 10 nm range with sufficient signal strength to probe magnetism. Current imaging techniques cover direct methods such as x-ray photo-emission-electron microscopes (XPEEMs) or scanning transmission x-ray microscopes (STXMs) as well as indirect methods based on coherent x-ray scattering such as holography, ptychography (PTY), and coherent diffraction imaging (CDI). Current developments are concerned with *in-situ* imaging, imaging of three-dimensional (3D) magnetisation distributions in planar and geometrically curved magnets as well as extending x-ray magnetic imaging to antiferromagnets. Strong future improvements are anticipated from the huge enhancement of the x-ray coherence flux fraction obtained from the upgraded diffraction-limited light sources. In addition, the possibility to study dynamic stochastic processes by imaging on ultrafast timescales based on x-ray free-electron lasers and from laser-driven higher harmonic generation sources will give a completely new view of ultrafast magnetisation with high spatial resolution.

### Electron-based magnetic microscopy

Imaging magnetisation using electron microscopes allows for magnetic characterisation with exceptional spatial resolution down to the atomic scale. The roadmap focuses on three electron-based magnetic imaging techniques, including Lorentz transmission electron microscopy (LTEM), electron holography (EH), and scanning electron microscopy (SEM) with polarisation analysis (SEMPA). LTEM and its reciprocal technique Lorentz scanning transmission electron microscopy (LSTEM) are used to image nanometric magnetic domains or even track the response of a spin texture to external stimuli such as heat, electric currents, or magnetic fields. In addition to further improving the resolution, advances in cryogenic and environmental

transmission electron microscopy (TEM) holders and fast direct detector technology place L(S)TEM at the forefront of magnetic imaging techniques in the decades to come to probe nontrivial spin textures and their dynamics. Recent developments of objective lenses of scanning transmission electron microscopy (STEM) provide access to imaging of magnetic field distribution at the atomic scale, even in antiferromagnetic materials, which opens new avenues for developing and understanding antiferromagnetic spintronics and multiferroics.

EH allows quantitative mapping of projected in-plane magnetic flux densities of magnetic two-dimensional (2D) and 3D textures at spatial resolutions approaching atomic length scales. When combined with electron tomography, it is possible to properly reconstruct all Cartesian components of the flux density in 3D of complex-shaped magnetic architectures at the nanoscale, such as magnetic skyrmions, domain walls, and vortices. Ongoing activities extend this technique towards complex sample manipulation using temperature, currents, electric and magnetic fields, light, and strain to study transient material responses within, e.g. magnetostriction and magnetoionics. In this respect, further development focuses on increasing the time resolution of the EH.

Complementary to the transmission-based electron microscopy imaging techniques, SEMPA allows ultimate surface sensitivity to image two-component magnetisation vector of magnetic textures at surfaces of ferromagnets and synthetic antiferromagnets. The recent method developments enable measurements of antiferromagnetic materials. The new advances in detectors with appropriate correctors for SEMPA are envisioned to improve the time resolution below 100 ps, provide imaging resolution better than 10 nm and enable studies of magnetic texture evolution in magnetic fields.

By addressing a range of key magnetic microscopy techniques where scanning probes, light, and electrons are used to examine and manipulate magnetisation and charge transport in condensed matter, we intend to provide snapshots of the current state and future projections within magnetic imaging techniques. We hope the roadmap will inspire the further use and development of these techniques within material science and physics.

## Acknowledgements

D V C acknowledges the support of Novo Nordisk Foundation NERD Programme: New Exploratory Research and Discovery, Superior Grant NNF21OC0068015. D V C and N P acknowledge the support of Novo Nordisk Foundation Challenge Programme 2021: Smart Nanomaterials for Applications in Life-Science, BIOMAG Grant NNF21OC0066526. N P acknowledges funding from the Danish Council for Independent Research Technology and Production Sciences for the DFF-Research Project 3 (Grant No. 00069B) and the funding from the European Research Council (ERC) Advanced 'NEXUS' Grant 101054572.

## 1. Scanning SQUID microscopy

*T R Devidas*<sup>1</sup>, *B Kalisky*<sup>1</sup> and *K C Nowack*<sup>2,3</sup>

<sup>1</sup> Department of Physics and Institute of Nanotechnology and Advanced Materials, Bar-Ilan University, Ramat Gan 5290002, Israel

<sup>2</sup> Laboratory of Atomic and Solid State Physics, Cornell University, Ithaca, NY 14853, United States of America

<sup>3</sup> Kavli Institute at Cornell for Nanoscale Science, Ithaca, NY 14853, United States of America

### Current state of the technique

Scanning SQUID microscopy (SSM) is a versatile technique to study and image magnetic, superconducting and transport properties at micro and nano-scales. This technique enables sensitive, spatially resolved detection of stray magnetic fields above a sample, across a wide range of temperatures.

A DC SQUID sensor consists of a superconducting loop interrupted by two Josephson junctions (JJs). For bias currents exceeding the JJ's critical current, a DC voltage develops across the SQUID. When magnetic flux threads the SQUID loop a screening current in the loop maintains flux quantisation within the loop. As a result, the critical current and the voltage across the SQUID are a periodic function of the applied flux with a period equal to the flux quantum,  $\phi_0 = h/2e$ , with  $h$  Planck's constant,  $e$  elementary charge [1].

Two key characteristics of a scanning SQUID sensor are flux sensitivity and spatial resolution. However, many other characteristics are of interest as well, including measurement bandwidth, ease-of-use, ease-of-fabrication, and functionality beyond magnetometry. The flux sensitivity is influenced by many factors, such as circuit and junction design, materials and the details of the scanning setup [2, 3]. The spatial resolution is largely determined by geometry, i.e. the size of the sensitive area and the distance between the sample and the sensitive area.

There are currently two main approaches to SSM, which differ in geometry and functionality of the SQUID sensor. One approach is to fabricate planar large-scale superconducting devices in which the sensitive area is a small part of the overall SQUID and is spatially separated from the JJs on the chip (SQUID on chip, SOC (figure 1(a))). These SQUIDs are typically based on Nb/AlO<sub>x</sub>/Nb trilayer JJ technology [4]. SOCs usually have micrometre-scale pickup loops and include on-chip circuitry such as flux modulation coils allowing for a flux-locked feedback loop to linearise the inherently nonlinear SQUID response, and an on-chip field coil concentric with the sensitive area, enabling local magnetic response measurements. Typical spatial resolution of SOC lies in the micrometre range [2].

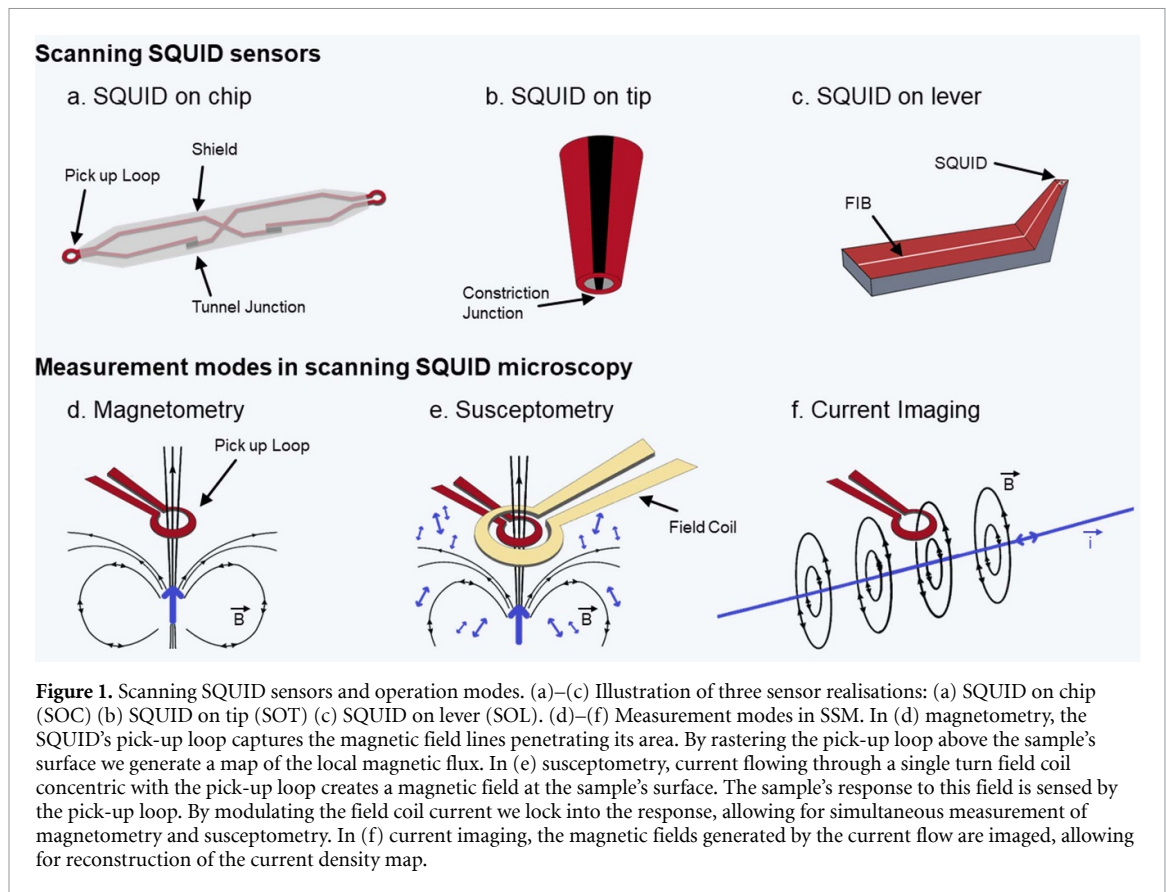
An alternative approach to realise a scanning SQUID is to fabricate a small SQUID loop for example at the apex of a hollow quartz tube (SQUID on tip, SOT (figure 1(b))) [5, 6] or, realised more recently, on an atomic force microscope (AFM) cantilever (SQUID on lever, SOL (figure 1(c))) [7]. The JJs in these SQUIDs are realised in the form of Dayem bridges, which allows the fabrication of loops with sub-100 nm diameter while maintaining  $\mu$ As of critical current. In SOTs, a quartz tube is pulled into a sharp pipette, and superconducting layers are deposited in three steps to define two superconducting leads connected to a superconducting loop. In SOLs, the loop is created via focused ion beam milling of a continuous superconducting thin film deposited on a blunted AFM cantilever. These geometries are suited for scanning in close proximity to a sample allowing for sub-micron spatial resolution. The SQUIDs have been made from different superconductors including Al, Nb, Pb, In, and MoRe which results in different performances in terms of flux sensitivity and continuous operation in magnetic field ranges up to several tesla.

The flux sensitivity of SSM sensors ranges from a few hundred  $n\phi_0/\sqrt{Hz}$  to a few  $\mu\phi_0/\sqrt{Hz}$  depending on the operating frequency and details of the sensor. The sensitivity to different sources of magnetic field then depends on how much flux is coupled into the SQUID, which in turn is determined by the SQUID geometry and the spatial dependence of the magnetic field lines produced by the source [2, 8].

The basic operation mode of SSM is magnetometry (figure 1(d)), which captures the static magnetic flux over a plane close to the sample. Susceptometry map is acquired by recording the local magnetic response of the sample to a varying local magnetic field applied via the on-chip field coil (figure 1(e)). The spatial distribution of electric current can be mapped by SSM by recording the magnetic fields generated by the current flow in a sample (figure 1(f)). Operating the last two with sinusoidal excitations allows measuring all modes simultaneously using lock-in detection. Mechanical vibrations at a different frequency can be used to improve DC magnetometry sensitivity [9].

### Future advances of the technique

SQUID sensors can be fabricated using nanolithography tools and rely on electrical measurement for readout, requiring only a few electrical leads. This makes it straightforward to envision numerous extensions and variations to the technique. Here we give examples of some possible directions.



The current SQUID sensors actively used for SSM have several limitations. SOCs have limited spatial resolution largely due to their geometry and layer structure, limiting the sample-pickup loop distance to around  $\sim 300$  nm. Additionally, the operation of SOCs is restricted to fields smaller than  $\sim 10$  mT due to the properties of the JJs. On the other hand, SOTs achieve high spatial resolution of about  $\sim 50$  nm and can operate in magnetic fields up to a few tesla. However, they do not operate in a flux-locked loop, which can limit the dynamic range and the linear relationship between the measurement signal and the magnetic field at the sample, important for quantitative reconstructions. Furthermore, a powerful capability of SOCs, simultaneous local magnetic response measurements, enabled by an on-chip field coil, is absent in SOTs and SOLs. Mapping the local susceptibility is particularly interesting when studying superconductors.

In terms of fabrication, SOTs are fabricated one by one with limited control over the resulting loop size. The recently developed SOLs achieved high spatial resolution and integration with force microscopy [7]. While the fabrication is more deterministic, each SQUID is fabricated individually with multiple processing steps using a focused ion beam. On the other hand SOCs are fabricated on wafer-scale.

The variety of scanning SQUID sensors covers a wide range of spatial resolution and sensitivity. One desirable future advance is nanoscale scanning SQUIDs that overcome the limitations of existing sensors and combine their advantages allowing for high spatial resolution, high sensitivity, and operation in magnetic fields, while realising advanced functionality through on-chip circuitry. This will likely require fabrication in a planar geometry using nanolithography and integration with a cantilever enabling scanning close to a sample. Using standard lithography tools would allow for wafer-scale fabrication with nominally identical properties on the same chip—a prerequisite for sharing sensors with the wider community and establishing user facilities. This could make SSM an accessible technique for the wider scientific community.

SQUID sensors can benefit from rapid advancements in superconducting, classical, and quantum information technology, to enhance measurement bandwidth, time resolution, sensitivity, and measurement modes. Possibilities range from utilising the same cryogenic amplifiers and read-out schemes to implementing a flux-sensitive scanning superconducting quantum bit. Several advances have already been demonstrated, such as the implementation of a dispersive microwave (MW) read-out of a scanning radiofrequency (RF) SQUID [10] and a sampling SQUID sensor integrating a Josephson pulse generator [11]. The dispersive readout relies on shunting the SQUID with an on-chip capacitor to form an LC resonator with a flux-dependent resonance frequency. This enables in principle higher sensitivity and wider measurement bandwidth. The sampling SQUID has demonstrated a  $\sim 40$  ps time resolution for periodically



triggered signals, limited by instrumentation rather than the intrinsic time resolution. Applying these sensors to quantum materials will open up interesting new dimensions. With increased bandwidth, SQUIDs could be used to probe magnetic noise complementary to the capabilities of NV centres in diamond, particularly in terms of operating temperatures and field range, and achieve unprecedented time resolution for magnetic imaging.

Advancements in SSM can extend beyond traditional magnetometry and magnetic response measurements. One direction is the on-chip integration of SQUID sensors with other on-chip sensors, enabling access to complementary physical properties with a single scanning probe. SQUIDs are well-suited for this due to their compatibility with standard lithography approaches. Another direction is integrating the ability to apply new stimuli or excitations to the sample with the SSM sensor and track magnetically detectable responses. For example, instead of a field coil, a local gate can be included, a nanoscale structure to apply local stress, or high-frequency structures to apply MW excitations. Leveraging the properties of JJs can give access to new signals. The strong dependence of their critical current on temperature can complicate SQUID performance, but can also be leveraged for high-sensitivity local thermometry [12]. Scanning a single JJ can unlock new capabilities. A JJ has flux-sensitive characteristics and emits and absorbs radiation at  $\omega = 2eV/\hbar$  where  $V$  is the voltage across the junction. Leveraging the latter, JJs have been used as on-chip tuneable sources and spectrometers of MW radiation with operating frequencies exceeding 100 GHz in non-scanning circuits [13]. The incorporation of sensing layers is another approach to expand the capabilities of SQUID and leverage their high sensitivity. For instance, SSM can contactless read out a sensing layer with a strong dependence on the local magnetic properties enabling, for example, highly-sensitive thermometry. Lastly, SQUID capabilities are expected to improve with material and fabrication development, for example, high- $T_c$  SQUIDs [14] will expand the accessible temperature range for SSM.

### Current and future use of the technique for material science

For the past three decades, SSM has been used to explore superconductivity, magnetism, and current distributions in a wide range of samples such as bulk crystals, thin films, interfaces, heterostructures, and nanowires. SSM probes the stray magnetic field above a sample that penetrates any non-magnetic top layer. Therefore buried, top-gated, patterned, lithographically defined and even ionic-gated samples can be studied with SSM. Different measurement modes and capabilities enable, for example, studying magnetic properties such as magnetic textures and susceptibility in magnetic materials and trapped flux, vortices, and diamagnetism in superconductors. Magnetic imaging also allows for the visualisation of current distributions in 2D materials. Currents and supercurrents can be mapped simultaneously with other electronic properties [2].

Using SSM for probing magnetic properties offers benefits beyond spatial imaging. Its local nature and high sensitivity allow investigation of weak signatures and small sample volumes, outperforming the sensitivity of bulk measurements by several orders of magnitude. Substrate contributions in thin film samples can be determined *in-situ* via side-by-side measurements in areas where the film is partially removed. This also helps in identification of the substrate's role in thin films and interfaces, where electronic properties of the active layer often stem from the substrate or its structural properties.

In the following, we describe the application of different measurement modes of SSM across different material classes and provide an outlook on future directions throughout. SSM serves as a tool for characterising multiple material properties, assessing their uniformity, detecting defects, and verifying predicted physical phenomena with magnetic signatures. Moreover, SSM can uncover new and unexpected phenomena. Due to the constraints of this roadmap format, we only highlight a few recent experiments in each area as a starting point for further exploration. We refer the reader to existing reviews for more details and for a complete overview of SSM.

**Probing magnetic materials.** SSM has been used to probe a variety of magnetic systems. It can map magnetic textures, and gain information even on resolution-limited disordered domain configurations. Isolated moments can be detected due to the high sensitivity and will appear as micro- or nano-scale objects due to the convolution with the sensor's point spread function. Reconstructing the moment configuration from SSM images is typically challenging since the moment distribution is underconstrained, i.e. no unique solution for the magnetisation exists given an image of the stray magnetic field. Unique reconstruction is possible in specific cases, such as a 2D magnetisation pointing along a fixed axis. Despite these constraints, imaging magnetic structures as a function of temperature, electrostatic gating, strain, and other conditions provides valuable insights into their nature and behaviour.

SSM has been applied to study magnetism at complex oxide interfaces, such as revealing the coexistence of magnetism and superconductivity at the  $\text{LaAlO}_3/\text{SrTiO}_3$  interface [15], strain tuneable magnetism [16], and studying the emergence of magnetic order in  $\text{LaMnO}_3/\text{SrTiO}_3$  above a critical  $\text{LaMnO}_3$  thickness

(figure 2(b)) [17]. Images as a function of magnetic fields can reveal superparamagnetic and ferromagnetic dynamics at the magnetic reversal, for example, in Cr-doped  $\text{Bi}_2(\text{Sb},\text{Te})_3$  [18]. SSM is sufficiently sensitive to study magnetism in few-layer van der Waals (vdW) structures fabricated through mechanical exfoliation, for example, it has enabled imaging of orbital magnetism in twisted BLG associated with a quantum anomalous Hall phase [19].

Local susceptometry can be used to track magnetic phase transition, by mapping the sample's magnetic response as a function of the temperature. This allows estimating the density of responding magnetic moments, and its frequency dependence can probe the dynamics of the moments. For example, an out-of-phase magnetic response in layered antiferromagnet  $\text{EuCd}_2\text{P}_2$  indicated the presence of magnetic fluctuation [20]. As the bandwidth and spatial resolution of SQUID sensors improve, SSM will be able to access more material systems, for example, frustrated magnetic systems. Beyond the traditional use of SSM, local strain tuning of magnetism has been shown by applying a small amount of force to the sample by pressing the SOC's blunt tip into the sample in a controlled manner [16, 21].

**Probing superconducting materials.** SSM has made significant contributions to the study of superconductors [2, 8]. For instance, the direct confirmation of the  $d$ -wave order parameter in cuprates was facilitated by SSM in a suite of elegant phase-sensitive measurements [22, 23] in which spontaneous currents in superconducting rings were probed.

As shown in figure 2(a), SSM can image the magnetic fields generated by Meissner screening currents, map the positions of superconducting vortices and explore their dynamic behaviour. The local magnetic susceptibility in superconducting samples is proportional to the superfluid stiffness, which is related to the superfluid density and the magnetic penetration depth for bulk and 2D superconductors.

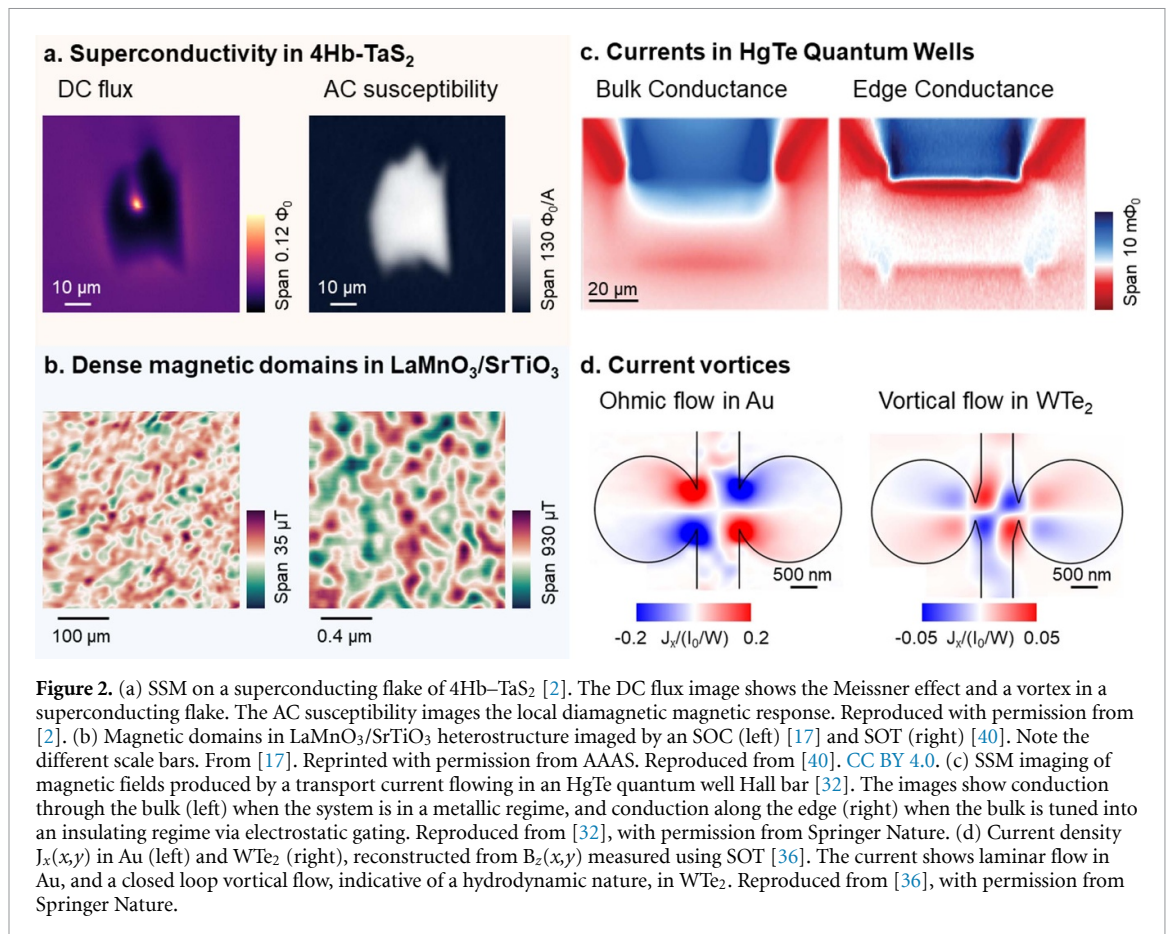
SSM can characterise the uniformity of a superconducting thin film and identify the presence of defects. SSM can distinguish different forms of spatial variation of the superfluid stiffness and critical temperature that may cause a broadened resistive transition. Measurements of the superfluid stiffness as a function of temperature have been performed on the micrometre scale and as a function of tuning parameters such as electrostatic gating. Through modelling, the superfluid stiffness can be extracted quantitatively. The low-temperature value, and the onset of the transition, and its temperature dependence can give insights on whether the superconductor realises the clean or dirty limit, if nodes and multiple gaps exist, and if a clear Berezinskii–Kosterlitz–Thouless transition is present. The imaging of vortices can provide insights into unconventional superconducting states, and advance understanding of flux trapping properties of conventional superconducting films and structures, which is important for applications including classical and quantum superconducting computing.

Again, we highlight a few recent examples. SSM can probe dynamics of superconducting samples, e.g. fast dynamics and flow of vortices was observed in a Pb film [24], whereas fluctuations in superfluid density near the critical temperature were probed in the superconductor NbTiN [25]. In microstructures of heavy-fermion superconductor  $\text{CeIrIn}_5$ , SSM revealed that differential thermal contraction between the structure and the substrate combined with strain sensitivity of  $\text{CeIrIn}_5$  causes a large-scale non-uniform superconducting transition [26]. SSM is sufficiently sensitive to small sample volumes and low superfluid density, such that it allowed measurements of the superfluid stiffness of  $\text{LaAlO}_3/\text{SrTiO}_3$  [15] and of micrometre sized samples of atomically thin ionic gated  $\text{MoS}_2$  [27].

SSM is a valuable tool for designing phase-sensitive measurements of unconventional superconducting order parameters [22, 23] and for determining current–phase relationships (CPRs) of exotic JJs. In these measurements, a JJ is embedded in a superconducting ring formed by an  $s$ -wave superconductor. Magnetic flux through the ring imposes a phase-bias across the JJ causing a supercurrent in the ring, which can be magnetically detected with the SQUID. To couple flux into the ring an on-chip field coil is used, enabling high bandwidth and rejection of background signals. As examples, the CPR of a JJ based on a 2D topological insulator [28] and of a few mode InAs nanowire JJ [29] have been measured.

Finally, SSM is also suitable for the discovery of unexpected phenomena with the discovery of a hidden magnetic phase above the superconducting transition temperature of the vdW superconductor  $4\text{Hb-TaS}_2$  [30] and vortex excitation that carry a temperature dependent fraction of the flux quantum [31] being recent examples.

With improved spatial resolution both for magnetometry and susceptometry, imaging supercurrent distributions in unconventional JJs and the proximitised superfluid stiffness will be feasible. Another direction is that atomic scale measurements of the superconducting gap and vortices by STM could be correlated directly with local SSM measurements.



**Visualising current distributions.** Electric currents produce magnetic fields. A 2D current distribution can be uniquely reconstructed from an image of one component of this stray magnetic field. This principle has been applied to visualise current distributions in various materials. In fact, one of the initial applications of SSM was troubleshooting failing integrated circuits through current imaging. In the context of materials research, visualising current distributions complements electronic transport measurements, which are often the first way to characterise new conducting materials. Transport measurements average over the sample volume sometimes in unexpected ways, in particular, when non-uniformities in a sample lead to real-space variations in the local conductivity. Visualising the current distribution has revealed surprises and explanations for unusual transport behaviour that was previously difficult to understand. A number of electronic phenomena lead to direct signatures in the current distribution with viscous electron flow, bulk versus edge currents in topological phases of matter, and electron guiding and focusing being examples. This measurement mode is not unique for SSM—for example, more recently NV centres in diamond are employed to visualise current distributions based on the same principle.

Here we highlight a few recent examples of using SSM to visualise current distributions. SSM showed that a transport current flows along the edge of a quantum spin Hall insulator (see figure 2(c)) [32], while finding that the current flows through the bulk in quantum anomalous Hall insulator Cr-doped Bi<sub>2</sub>(Sb,Te)<sub>3</sub> [33]. At the LaAlO<sub>3</sub>/SrTiO<sub>3</sub> interface, the role of structural domains was revealed by visualising their impact on the local current distribution [34] and on the universality of the metal–insulator transition in SrTiO<sub>3</sub>-based interfaces [35]. The presence of vortices in the viscous electron flow regime was verified by direct visualisation (see figure 2(d)) [36]. In addition to transport currents, equilibrium currents can exist in samples. Coexisting topological and non-topological currents have been visualised in the quantum Hall regime in graphene [37], as well as equilibrium currents arising from twist-angle disorder in magic-angle twisted bilayer graphene [38]. Imaging current in superconducting wire networks [39] reveals non-uniformity of the current distribution and the role of disorder close to the superconducting transition temperature.



**Some concluding remarks**

SSM is a powerful and versatile technique that has been used for more than three decades to provide local magnetic measurements of quantum materials. The technique enables sensitive, spatially resolved detection of stray magnetic fields above a sample, at low, but variable temperatures. We suggest that advancements in SSM can extend beyond traditional magnetometry and magnetic response measurements, and can be used to access complementary physical properties with a single scanning probe. We hope that SSM becomes more widely accessible in the future to accelerate the development of new quantum materials.

**Acknowledgement**

T R D and B K were supported by ERC Grant No. ERC-2019-COG-866236 and Israeli Science Foundation Grant No. ISF-228/22. K C N acknowledges support from the Air Force Office of Scientific Research (MURI Grant No. FA9550-21-1-0429).

## 2. NV magnetometry

### 2.1. Imaging electronic transport in materials using colour centres

*J L Webb, U L Andersen and A Huck*

Centre for Macroscopic Quantum States (bigQ), Department of Physics, Technical University of Denmark, 2800 Kongens Lyngby, Denmark

#### Current state of the technique

Optically active point defects (colour centres) in the solid state offer an excellent platform for imaging electrical transport in a wide range of materials, ranging from biomaterials to inorganic 2D vdW heterostructures. Acting as individual atomic scale sensors, colour centres provide high sensitivity to the local environment, high spatial resolution (nano-metre/micro-metre scale for single/ensemble sensors) as well as fast readout and initialisation (both sub- $\mu\text{s}$ ) for time-resolved measurements. The principle of using colour centres to study electronic transport is well established in three material systems: boron vacancies in hexagonal boron nitride (hBN), divacancies in silicon carbide (SiC) and in particular NV centres in diamond. Since the original idea of a spin microscope with optical readout [41], NV-based sensing [42, 43], and the first demonstration of microscopic imaging [44], the field has rapidly developed, and offers strong potential for future materials science studies.

Sensing and imaging with colour centres relies on the well-established technique of optically detected magnetic resonance (ODMR) spectroscopy. In this process, local electric or magnetic fields modulate the defect fluorescence emission when illuminated by pump laser light and a resonant field in the RF to MW range. A more comprehensive discussion on the sensing principle and limitations is extensively covered in [45, 46].

Interest in colour centre imaging has primarily been driven by its ability to image vector magnetic flux density ( $B$ ) or electric field ( $E$ ). For magnetic field imaging, this attribute is particularly appealing for systems without the presence of any material magnetisation, where alternative materials techniques such as Kerr microscopy or MFM cannot be used. Another intriguing advantage is that field imaging with colour centre can be entirely passive, requiring no mechanical contact of the target material. When illumination on target sample is negligible, this allows electronic transport (charge or spin) within a material to be studied by induced electric/magnetic field without risk of disrupting sensitive effects via mechanisms such as heating or inducing photoelectric charge.

State-of-the-art sensing of electrical transport can be divided into ensemble and single defect methods. Ensemble methods utilise a dense ensemble of colour centres within a bulk host, located in close proximity to the material to be imaged, with the defect fluorescence emission captured by a microscope objective and a charge coupled device (CCD)/complementary metal-oxide semiconductor (CMOS) camera. This technique enables capturing transport properties over a wide FOV with diffraction-limited spatial resolution (approximately at the micron level), with the temporal resolution by camera frame rate. In contrast, single NV methods use just a single colour centre, hosted either in a fixed host material or in a scannable tip. Imaging can be achieved by translating either tip or sample in a raster pattern while recording data from the colour centre using a sensitive photodetector. For fast transport processes, the single defect method effectively offers only a point probe measurement due to the slow rate of translating either probe or sample. However, for slow processes, repeating signals that can be averaged or variance detected, the method affords superior spatial resolution, higher proximity to the target material, greater resistance to environmental decoherence and faster readout.

#### Future advances of the technique

High-quality imaging requires a strong defect centre optical response to the target field, strong resistance to decoherence while maintaining sensitivity to environmental changes, and a minimum level of measurement noise. Achieving these requires in turn a stable, contamination-free host material, efficient defect generation, effective defect control, efficient optical readout, and some degree of noise rejection through shielding or filtering.

Future advances in these areas can be divided into materials and process development. This includes both host and defect generation, improvements in coherent control and readout of the defects and methodologies to minimise measurement noise.

Current efforts in host material engineering seek to mitigate factors such as host material strain and contamination and charging, that adversely affect sensing performance, particularly for colour centre ensembles. An important aspect here is isotopic purification, which removes isotopes (e.g.  $^{13}\text{C}$ ) with intrinsic spin that would otherwise cause decoherence and reduce sensitivity. For diamond,  $^{12}\text{C}$  enriched substrates

have begun to be commercialised. A significant recent advance is the achievement of isotopic purification in SiC [47], yielding up to a 50 fold improvement in coherence times. Another promising future route is material growth with preferentially aligned colour centres, ensuring a uniform crystallographic orientation for all colour centres in an ensemble, which optimises the response (in terms of ODMR contrast) [48].

For control and readout, methods derived from NMR/electron spin resonance (ESR) using pulsed control fields are being implemented for imaging experiments to improve the control of the colour centre state. This includes work on optimal control such as MW pulse shaping [49] and advanced pulse sequencing. The aim is to achieve noise rejection and overcome host crystal inhomogeneity, offered by the double quantum four-Ramsey technique [50]. Future work on using dynamical decoupling sequences for AC sensing holds potential for substantial improvement. Similarly, progress in the transmission of MW nearfield antenna signals to the material could yield increased Rabi frequencies and more homogenous control.

Minimisation of measurement and background noise is key to optimal imaging performance. This noise can originate from the system itself (e.g. thermal), background fields or from control/readout. Critical to noise minimisation are the development of common mode rejection methods, use of mu-metal enclosures and/or frequency and time domain filtering methods. Widefield imaging at low frame rates and relying on temporal averaging is especially susceptible to noise sources, which for conventional CCD/CMOS cameras can cause the sensitivity to plateau when time averaging across multiple acquisition frames. Future use of cameras with in-built lock-in amplification [51], balanced detection techniques [52] or using the optically excited states [53] could prove successful in yielding high-SNR imaging.

Ultimately, there remains two fundamental challenges for colour centre imaging. The first is the low efficiency of the colour centre systems, regarding the ratio of the input optical and MW power to usable signal output. This is due to (a) the necessity to keep colour centre density low (to avoid defect-defect interaction), resulting in an overall low light absorption cross section and (b) due to difficulty in efficiently collecting fluorescence emission. The latter is particularly problematic for host crystals with high refractive index such as SiC and diamond. The second substantial problem is the high level of non-signal background fluorescence emission, adding a high level of optical shot noise. It would instead be desirable to have a system emitting light only when activated by an MW field resonant with the colour centre energy levels.

Although these challenges could be tackled by material engineering (e.g. larger substrates) or through optical engineering (e.g. solid immersion lenses), it would be advantageous to explore alternative systems. This may include defects hosted in 2D materials, avoiding the refractive index issue and potentially offering strong, visible-range emission suited to Si-based imaging detectors. Here, ongoing *ab initio* studies utilising automated machine-learning based searches [54] may identify promising optically active defects. This could also include hybrid systems, such as the introduction of additional dopants or overgrowth on diamond in order to increase optical response or suppress decoherence effects [55].

### Current and future use of the technique for material science

To date, the majority of experiments for materials science investigations have relied on widefield imaging using a defect ensemble, as this provides the simplest method to realise high spatial resolution (on the  $\mu\text{m}$ -scale) and wide FOV ( $\mu\text{m}$ -mm scale). A particular focus has been the imaging of electrical current transport in microelectronic devices by induced magnetic field, for both fundamental study and for applications such as circuit inspection.

Imaging can be achieved with sensing bandwidths up to several kHz sub-ms time resolution. Although this may not be fast enough to capture in real-time electronic signal propagation in microdevices, temporal averaging can allow facilitate the imaging of average current propagation patterns. Studies have included inspection of capped and uncapped integrated circuits [56], and in particular study of current flow in novel devices made from 2D materials such as graphene [57, 58] or transition metal dichalcogenides including imaging of photoinduced current [59]. 2D materials are ideally suited for colour centre imaging as transport is confined to a thin layer which can be in high proximity to the colour centre host. Their flatness also makes them highly suitable for study with NV-based SPM systems. Due to the rapid technical development of such SPM systems and their potential for enhanced spatial resolution, it is likely these will play an increasing role in future materials science studies.

Another intriguing application of colour centre imaging is its use in exploring more exotic electronic transport phenomena, such as spin transport or superconductivity. Imaging is particularly suited to superconducting materials with low critical currents and accessible magnetic phase transitions, such as the Fe-based superconductors [60]. NV-based imaging has been successfully used to image spin transport phenomena such as spin waves induced in yttrium iron garnet (YIG) [61]. Looking forward, the development of low temperature NV-SPM systems, capable of probing both electric and magnetic field at

mK temperatures, is of special interest [62]. These systems may prove useful in both the study of both fundamental phenomena (e.g. topological phases) as a complement to traditional transport methods and for applications such as device characterisation, including superconducting qubits for quantum computing.

Finally, colour centre imaging has demonstrated significant potential in recording transport in biomaterials, including electric impulses (e.g. action potentials) in living tissue, enabled by the high degree of biocompatibility of the defect host materials, particularly diamond. Proposals have been put forward to image electrical current transport in living neural networks, for fundamental studies in biology [63]. Compared to existing techniques (e.g. GEVIs), this offers the advantage of requiring no biological modification or use of potentially toxic fluorescent dyes. Although this application has yet to be realised, it could be achieved with modest future improvement in sensitivity. Another relevant development for biomaterials is the use of colour centres in biocompatible nanomaterials (e.g. nanodiamonds) as highly stable fluorescence markers [64] and nanoscale biosensors [65]. Combined with an imaging method such as scanning confocal microscopy using a spinning (Nipkow) disc system, imaging a distribution of colour centre markers could deliver extremely high-quality, sub-wavelength information on both electronic and thermal transport in biomaterials.

### Acknowledgements

We acknowledge financial support from the Novo Nordisk Foundation for the Biomag project (NF21OC0066526) within the challenge programme and the Danish National Research Foundation (DNRF) through the centre for Macroscopic Quantum States (bigQ, Grant No. DNRF0142).

## 2.2. Imaging magnetism using NV magnetometry

David A Broadway<sup>1,2</sup>, Kai Wagner<sup>1</sup> and Patrick Maletinsky<sup>1</sup>

<sup>1</sup> Department of Physics and Swiss Nanoscience Institute, University of Basel, Klingelbergstrasse 82, 4056 Basel, Switzerland

<sup>2</sup> School of Science, RMIT University, Melbourne, Victoria, Australia

### Current state of the technique

The NV centre in diamond is an atomic scale lattice defect that host an electronic spin that allows for optical spin initialisation and readout. Together with coherent spin manipulation using MW-frequency driving fields and operation even under ambient conditions, this renders the NV a highly interesting spin system that finds applications in various quantum technologies and first and foremost quantum sensing and nanoscale magnetometry.

Such nanoscale magnetic imaging is based on placing the NV centre spin in close proximity to a sample of interest, in order to measure the stray magnetic fields generated by the sample. There are various experimental approaches to achieve this goal, each with their own advantages and disadvantages, the most widely used being AFM using diamond tips hosting NV centre spins (NV-AFM, figure 3(a)) [66] and wide-field optical imaging using the ‘quantum diamond microscope’ (QDM, figure 3(b)) [67]. A highly noteworthy variant of the QDM employs NV centres embedded diamond anvil cells (figure 3(c)) [68] for magnetic sensing and imaging under high pressures. A key asset of all these techniques is their applicability over a vast range of conditions, operating from the lowest achievable temperatures up to 1000 K [69] and at pressures from ambient to >10 GPa [68].

In the following, we list and briefly discuss some key technical performance characteristics of NV based magnetic imaging modalities:

### Measurement protocol

All mentioned techniques detect magnetic fields through the Zeeman effect, building on optical spin initialisation and readout, together with MW driving. The most commonly employed approach is ODMR (figure 3(d)), but more refined method using pulsed, coherent NV driving exist and are described in excellent prior reviews [45].

### Reconstruction of the source

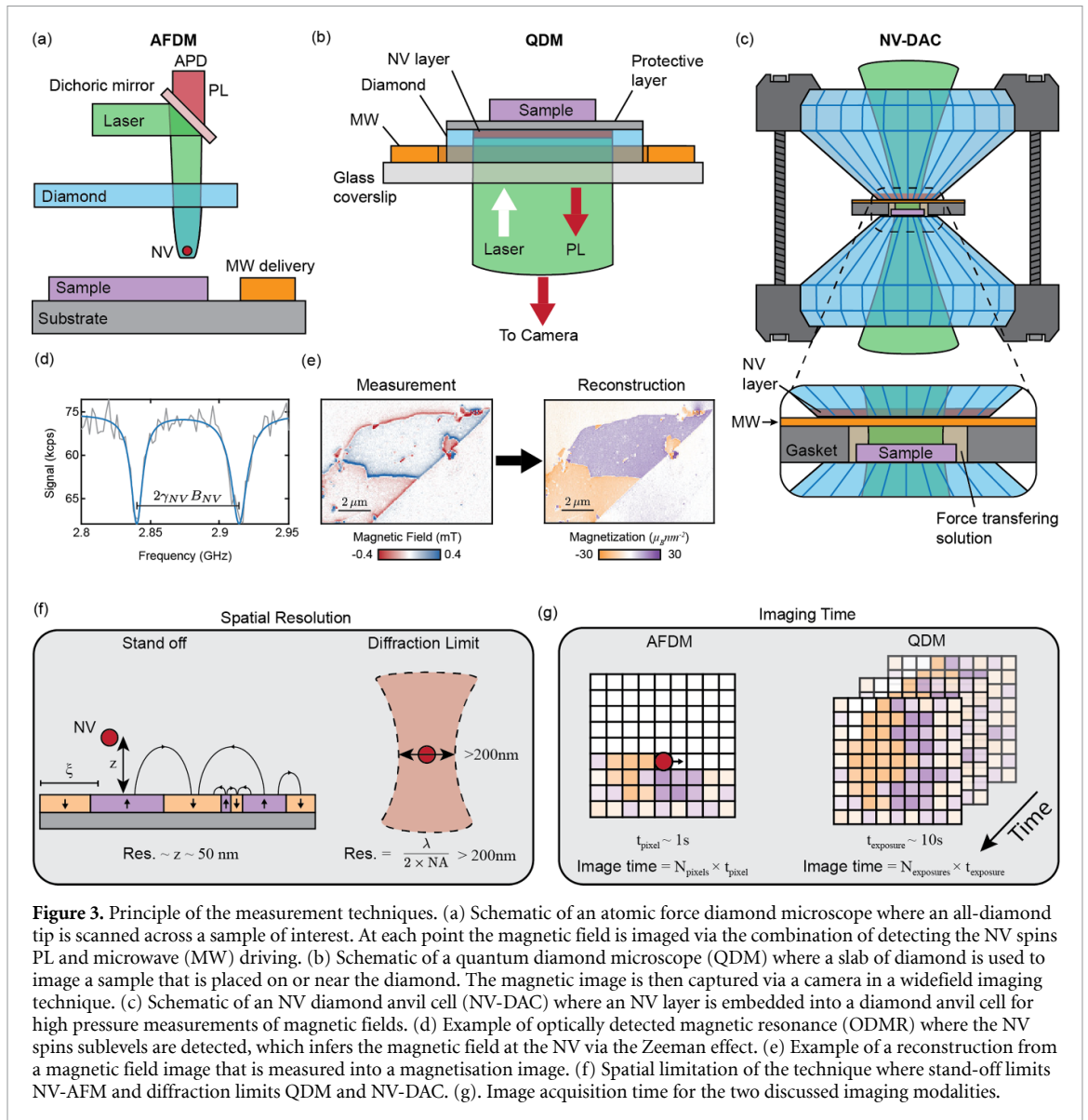
The main goal of NV-based magnetic imaging is to infer information about the magnetic field source, e.g. currents or magnetisation distributions in the sample under study. The required reconstruction from the imaged magnetic field (figure 3(e)) is an ill-posed problem and as such presents major difficulties in the reconstruction process. Various methods have been discussed in recent years for how to best address this difficult inverse problem [70, 71].

### Spatial resolution

Spatial resolution in NV magnetometry is ultimately limited by standoff distance between the NV spin and the sample (figure 3(f)) [70]. For NV-AFM with single NV centres, this limit in spatial resolution is effectively realised and NV-sample distances  $z_{\text{NV}}$  in the range of 20–70 nm have been reported [72]. Conversely, in the QDM, spatial resolution is determined by the diffraction limit, since ensembles of NV centres are addressed using a conventional optical microscopy. While in principle it should be possible to achieve a resolution of  $\approx 200$  nm, in practice this is not achieved due to optical aberrations, leading to a resolution in the range of 400 nm [73].

### Imaging time

Imaging time is determined by the desired SNR, whose scaling with time is determined by the sensitivity. The latter is typically on the order of  $1 \mu\text{T Hz}^{-1/2}$  for single NVs and improves with  $N^{1/2}$ , if a number of  $N$  NVs are employed simultaneously. In NV-AFM, data needs to be acquired pixel-by-pixel, and the imaging time therefore scales linearly with the number of pixels. Conversely, in QDM, all pixels of an image are acquired simultaneously, which leads to a significant speedup in acquisition time (figure 3(g)). In addition, NV-AFM typically employs single NVs for sensing, while in the QDM, several NVs are interrogated in each pixel. As a result, NV-AFM is a comparatively slow technique (acquisition times of several minutes to hours) with high spatial resolution, while QDM has a limited resolution, but allows for rapid data acquisition times on the order of seconds to minutes [73].



**Figure 3.** Principle of the measurement techniques. (a) Schematic of an atomic force diamond microscope where an all-diamond tip is scanned across a sample of interest. At each point the magnetic field is imaged via the combination of detecting the NV spins PL and microwave (MW) driving. (b) Schematic of a quantum diamond microscope (QDM) where a slab of diamond is used to image a sample that is placed on or near the diamond. The magnetic image is then captured via a camera in a widefield imaging technique. (c) Schematic of an NV diamond anvil cell (NV-DAC) where an NV layer is embedded into a diamond anvil cell for high pressure measurements of magnetic fields. (d) Example of optically detected magnetic resonance (ODMR) where the NV spins sublevels are detected, which infers the magnetic field at the NV via the Zeeman effect. (e) Example of a reconstruction from a measured magnetic field image. (f) Spatial limitation of the technique where stand-off limits NV-AFM and diffraction limits QDM and NV-DAC. (g). Image acquisition time for the two discussed imaging modalities.

### Future advances of the technique

There are several key improvements that can be made to further enhance the sensitivity, operating range and utility of NV scanning magnetometry in the future.

#### Spatial resolution

The main fundamental limit to NV spatial resolution is the NV-sample distance  $z_{NV}$ , mentioned above. While current implementations of the NV-AFM realise values of  $z_{NV}$  in the range of tens of nm, efforts are underway to achieve single-digit nanometre values of  $z_{NV}$  and therefore in spatial resolution. For this, two challenges need to be addressed. First, NVs need to be stabilised close to the surface, without significantly degrading their optical, spin, and charge-properties—a formidable and still open challenge in diamond materials engineering. Second, the geometry of all-diamond scanning probes should be adapted to allow for such small values of  $z_{NV}$ . Current realisations of diamond tips consist of cylindrical tips of  $\sim 200$  nm diameter, which pose a limitation on its own with regard to optimising  $z_{NV}$ . However, interesting new approaches are on the horizon [74] that may allow for the realisation of ultrasharp diamond tips for the realisation of single-digit-nm resolution in NV-AFM.

#### Sensitivity

Sensitivity in NV magnetometry is primarily determined by the NV spin coherence time,  $T_{coh}$ , and the efficiency of optical addressing of the NV spin, as characterised, e.g. by the photon count-rate  $I_{PL}$  of NV photoluminescence. NV magnetic field sensitivity scales with  $1/(T_{coh} * I_{PL})^{1/2}$ , which motivates improvements in NV optical addressing (e.g. through photonic engineering) and NV coherence times (e.g. by



materials engineering or advanced, pulsed quantum sensing protocols). Both aspects have been discussed in detail in past reviews [45].

An additional, sometimes overlooked aspect is the control over the NV charge state. Indeed, only the negatively charged NV centre,  $NV^-$ , is directly applicable to magnetometry, while other charge states (i.e.  $NV^+$  and  $NV^0$ ) exist but should be avoided for magnetometry. Charge state stability directly translates to enhanced magnetic field sensing capabilities, as only the single negatively charged state is utilised for sensing. Approaches for stabilising the NV charge state have been demonstrated (e.g. through diamond surface modifications [75] or multi-colour laser illumination) and are subject to ongoing research in particular with respect to the required nano-fabrication for scanning probes [76]. The different possible NV charge states, however, also offer interesting opportunities for NV spin-to-charge conversion and thereby efficient (even single-shot) NV spin readout [77, 78], or photoelectric readout of the NV spin state [79, 80].

### Real-time spin and charge control

Advanced electronic components for NV spin control, such as field programmable gate arrays offer another ongoing avenue for future improvements in NV sensing capabilities. For example, using such hardware, it was recently shown that the NV charge state can be deterministically initialised in real-time [81]. Furthermore, NV spin-based measurement sequences can be adapted on the fly providing efficient, real-time tracking of the instantaneous NV spin transition frequency [82]. Real-time computation can also aid in targeting measurement points of expected higher information gravity via Bayesian updating/sampling and speed-up towards live quantum optimal control [83].

All the above-mentioned aspects closely relate to improving the spatial resolution, as stabilising NV centres close to the surface has the potential of reaching the single-digit nanometre scale. This will also increase all detected signals from a target sample. Therefore, the aforementioned surface modifications and control schemes will take a key factor with particular importance for fundamental physical questions on the nanometre-scale.

### Multimodal sensing operation

Next to the excellent sensitivities in measuring static and AC magnetic fields, the NV centre grants access to several other quantities, such as electric fields or temperature. This is complemented by spectroscopy of magnetic [84] or electrical noise [85].

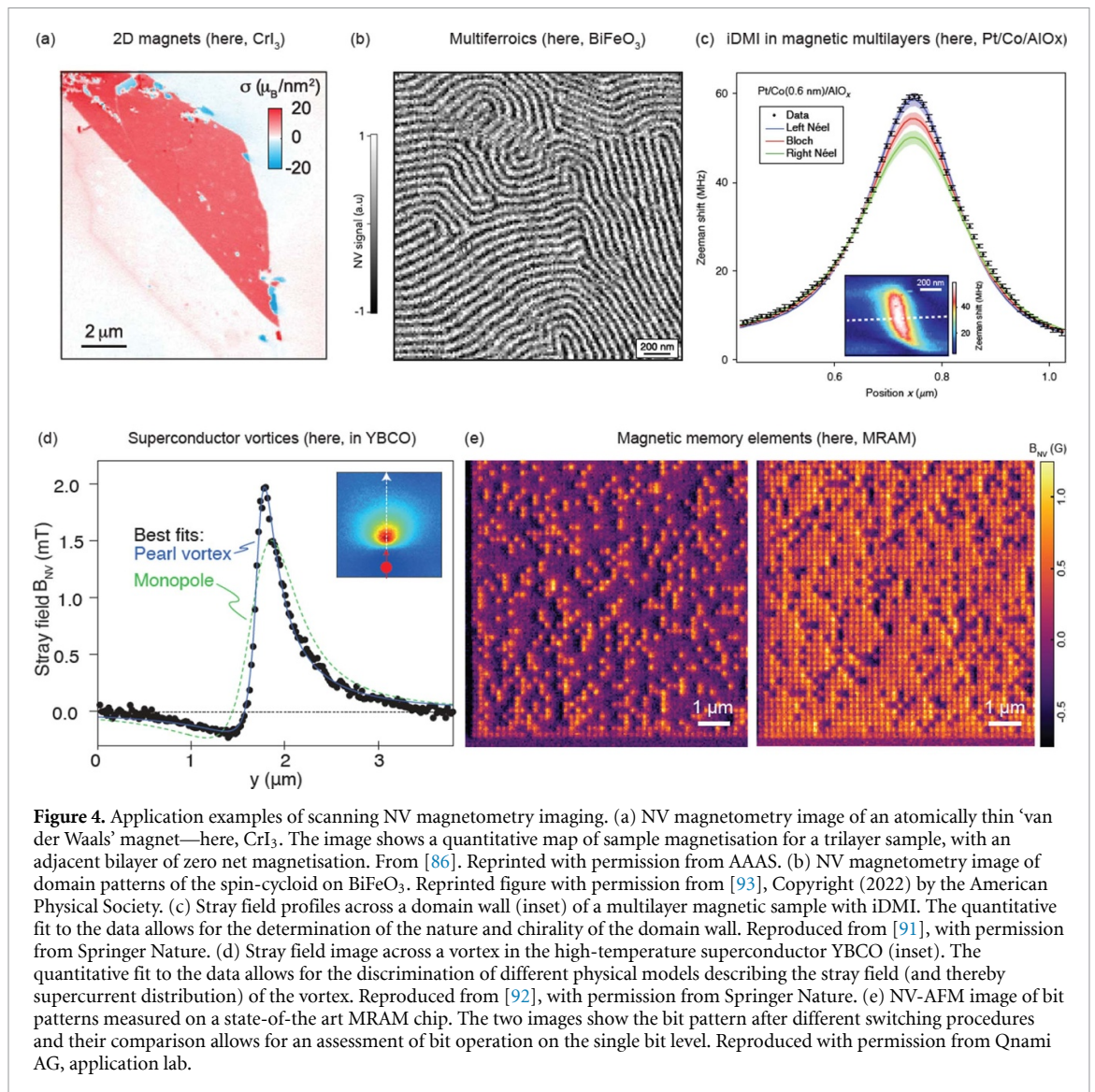
The possibility to sense several different physical quantities using the same sensing platform and conditions is another key asset of NV based sensing that can help develop a better microscopic understanding of physical processes in samples under study—a prime example is the simultaneous addressing of ferroelectric and magnetic ordering in multiferroic materials (see below). While such multimodal imaging typically increases measurement time and protocol complexity, research is underway to further advance and simplify such multimodal NV imaging.

### Current and future use of NV magnetometry for material science

Scanning NV magnetometry has the potential to significantly impact the field of material science in the coming years. This was already shown with several use-cases and key publications over the last ten years, and the resulting impact is poised to grow further in the future. The power of NV magnetometry in this regard lies not only in its high sensitivity and nanoscale spatial resolution, but for many applications also in the quantitative and minimally invasive nature of the technique. These combined properties open up new possibilities for the study and development of advanced magnetic materials and devices. In this section, we will discuss how NV magnetometry impacts several areas of application, including 2D magnetic vdW materials and their heterostructures, antiferromagnets and multiferroics, skyrmion hosting magnetic materials, materials with engineered interfacial Dzyaloshinskii–Moriya interaction (DMI), superconductors, and industry applications such as failure analysis and novel memory devices.

2D magnetic vdW materials have emerged as a promising platform for spintronics and magnonics, with potential applications in data storage, processing, and communication. Scanning NV magnetometry offers a unique opportunity to probe these materials' magnetic properties. In fact, scanning NV magnetometry is so far the only technique that allowed for nanoscale imaging of spin-textures in atomic monolayers of a vdW magnet [86] (figure 4(a)). Furthermore, the technique allows for addressing magnetism in vdW heterostructures (such as moiré stacks or gated magnets) and thereby extract, e.g. information about interlayer exchange couplings. This deeper understanding could lead to the development of new vdW magnetic materials and devices with tailored properties, such as tuneable magnetism, spin transport, and magnon-mediated information transfer.

Antiferromagnets and multiferroics are another class of materials, where NV magnetometry already showed impact and will continue to do so. Antiferromagnets exhibit magnetic ordering but (near) zero net



magnetisation and as such are hard to address by conventional experimental techniques. However, weak uncompensated magnetic moments are known to occur when bulk symmetries of an antiferromagnet are broken, i.e. on surfaces or domain walls. The resulting weak magnetic stray fields enable access to nanoscale properties of antiferromagnets by exploiting the sensitivity offered by NV magnetometry. Multiferroics are materials that simultaneously exhibit ferroelectric and magnetic ordering, offering the possibility of controlling magnetism with electric fields and vice versa. Often-times, such as for the prototypical multiferroic BiFeO<sub>3</sub>, the magnetic order in multiferroics is of AFM nature. Scanning NV magnetometry then enables the observation and control of complex magnetic textures in these materials [87–89] (figure 4(b)), ultimately leading to the development of novel AFM and multiferroic devices with improved energy efficiency and new functionalities.

Skyrmions, topologically protected magnetic quasiparticles, have garnered significant interest due to their small size, stability, and potential for low-power spintronic applications. Both bulk and thin-film skyrmion hosting magnetic materials can be studied and manipulated using scanning NV magnetometry [90] which can provide crucial insights into the formation, stability, and manipulation of skyrmions. This knowledge could pave the way for the development of skyrmion-based devices, such as racetrack memories and logic gates, that can operate at ultralow power consumption and high speed.

Magnetic multilayer stacks that induce interfacial DMI (iDMI), a chiral exchange interaction that stabilises non-collinear magnetic textures, form another area where NV centre magnetometry can make a significant impact. By studying materials with engineered iDMI, researchers can not only quantify the strength of iDMI [91] (figure 4(c)), but also gain insights into the complex interplay between magnetic moments and crystal lattice structures. This understanding could lead to the development of advanced materials with tailored DMI properties for use in spintronic and magnonic applications.



Superconductors and superconducting devices also stand to benefit from advances in scanning NV centre magnetometry [92] (figure 4(d)). By probing the local magnetic fields in superconducting materials, researchers can gain a better understanding of the mechanisms behind superconductivity, vortices and vortex dynamics, ultimately leading to the development of improved superconducting devices, such as SQUIDs.

Finally, NV centre magnetometry's nanoscale sensitivity and non-destructive nature make it an ideal tool for various industry applications, including failure analysis and the development novel memory devices, such as magnetic random-access memories (MRAMs) (figure 4(e)). In failure analysis, the ability to identify and localise defects in magnetic materials and devices can lead to improved quality control, reliability, and lifetime. In the realm of novel memory devices, NV magnetometry could enable the development of high-density, energy-efficient storage solutions based on new memory architectures.

## Acknowledgements

We acknowledge financial support from the Swiss Nanoscience Institute, the National Centre of Competence in Research Quantum Science and Technology, a competence centre funded by the Swiss National Science Foundation (SNF), by SNF Project No. 188521, and by the ERC consolidator Grant Project QS2DM. DAB acknowledges support through an Australian Research Council DECRA Fellowship (Grant No. DE230100192)

## 2.3. Imaging magnetic dynamics using spins in diamond

Toeno van der Sar<sup>1</sup>, Chunhui Rita Du<sup>2</sup> and Amir Yacoby<sup>3</sup>

<sup>1</sup> Department of Quantum Nanoscience, Kavli Institute of Nanoscience, Delft University of Technology, Lorentzweg 1, 2628 CJ Delft, The Netherlands

<sup>2</sup> School of Physics, Georgia Institute of Technology, Atlanta, GA 30332, United States of America

<sup>3</sup> Department of Physics, Harvard University, 17 Oxford Street, Cambridge, MA 02138, United States of America

### Current state of the technique

Condensed matter physics thrives on the interplay between discoveries and the development of new measurement capabilities that are enabled by such discoveries. While in some cases this development involves measuring new physical quantities, often it simply involves using novel physics to report new quantities of interest. For example, global positioning systems report position accurately through the measurement of time and uses general relativity to connect the measured quantity to the reported one. Therefore, by using established physics one can report quantities that are very hard to measure.

Magnetometers based on NV spins in diamond (figures 5(a) and (b)) measure magnetic fields in a rather broad range of frequencies (DC to several GHz) and with high spatial resolution (down to  $\sim 50$  nm) [94]. Recent developments in NV centre magnetometers have made considerable impact in our ability to explore quantum matter through the use of novel physics that enables reporting quantities that are otherwise hard to measure. One such example is the use of the fluctuation dissipation theorem to connect the NV spin relaxation to the spin chemical potential in a magnetic system [95, 96]—an important quantity governing spin transport. The spin excitations in a magnet result in magnetic field fluctuations that leads to relaxation of a nearby NV sensor and thereby enables quantitative determination of the local chemical potential [95].

Another capability enabled by NV magnetometers is the use of spin waves in scattering experiments. Such experiments consist of launching coherent waves or particles with well-defined energy and momentum at a target. Through the measurement of the amplitude and phase of the scattered signal one can reconstruct the target's structure factor [97] and probe its magnetic susceptibility [61]. NV spins can measure the amplitude [98] and phase [97, 99] of propagating spin waves through the magnetic field they generate in time and in space (figures 5(d)–(g)). A key advantage of the technique is that it allows probing the coherent interactions of the spin waves with target materials and the spin textures and electric currents they host.

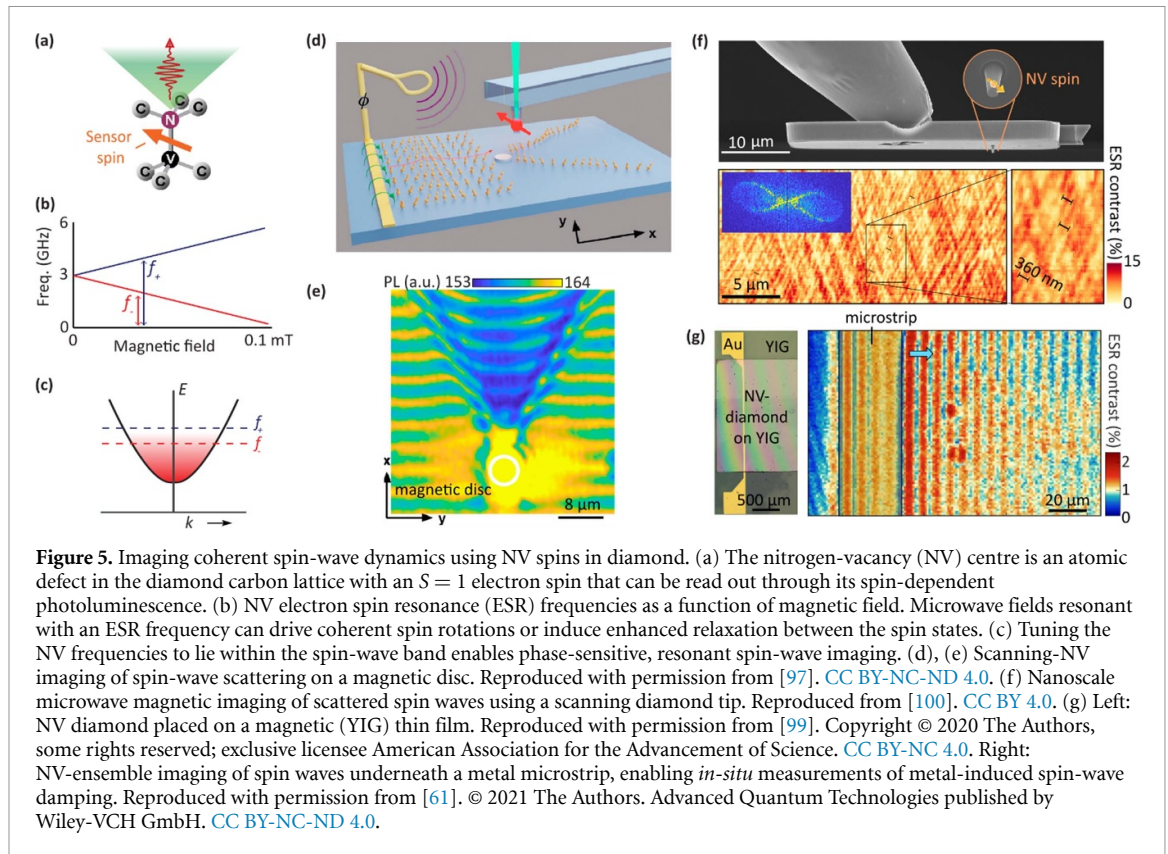
### Future advances of the technique

We expect the future advances of the technique to focus on studying new magnetic materials, on the interaction of spin waves with spin textures and electric currents, and on expanding the detectable frequency and wavelength range. In current NV-based spin-wave detection schemes, two situations can be distinguished. In the first, the NV spin resonance frequencies lie within the spin-wave band (figure 5(c)), such that spin waves resonant with the NV can directly drive NV spin transitions. In the second, the NV spin resonance frequencies lie below the band (next section), such that the NV is sensitive to longitudinal magnetisation dynamics generated by spin-wave mixing [96]. Both situations provide excellent new opportunities for probing magnetic materials and spin-wave—material interactions. Here we describe the current focus of efforts to advance these techniques.

With phase-sensitive imaging of NV-resonant spin waves established (e.g. figures 5(d)–(g)), a key goal for realising broad applicability of the technique is to develop protocols that will enable imaging of spin waves far detuned from the NV resonance frequency. Off-resonant spin-wave detection using NV magnetometers would enable high-resolution probing of spin-wave physics in the wide range of magnetic materials with strong saturation magnetisations or magnetic anisotropy, such as ultrathin vdW magnets [86]. We discuss two approaches to meet this challenge.

The first approach focuses on developing NV spin-control techniques that render the NV spin sensitive to MW signals that are far detuned from the ESR frequency. The development of such protocols is an active field of research, with state-of-the-art methods having already demonstrated detection capabilities at  $\sim 100$  MHz detuning from the NV ESR frequency [101]. However, detecting signals that are detuned by multiple gigahertz, as required for many magnetic materials of interest, remains a challenge as the sensing sensitivity typically decreases with increased detuning. As such, there remains a need for new NV-protocols enabling broadband MW detection.

The second approach focuses on converting the intra-band spin-wave excitations to magnetic dynamics at frequencies that are detectable with standard NV sensing protocols. On the one hand, mixing between *thermal spin waves* leads to NV-resonant longitudinal spin *noise* that provides access to the magnetisation of a broad range of magnetic materials (next section). On the other hand, mixing between *MW-driven* spin waves

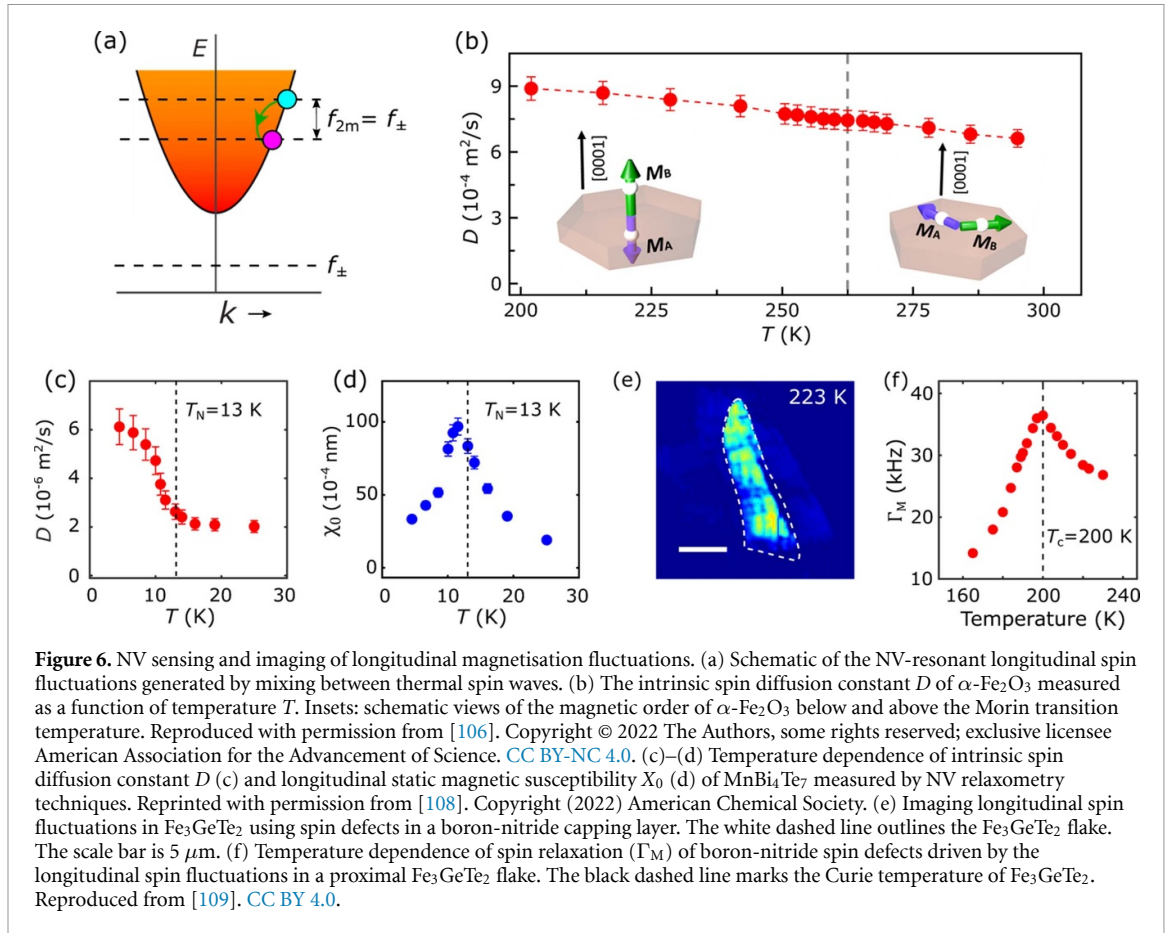


results in *coherent* dynamics of the longitudinal magnetisation [102]. Broadband detection of the difference-frequency signals is possible by flipping the NV spin at a rate that matches the difference frequency using AC magnetometry/dynamical decoupling protocols [94]. We anticipate that the driven spin-wave mixing processes will enable mapping spin-wave dispersions even at large detuning with the NV frequency, as the frequency—and wavelength—difference of the driven primary spin waves directly determines the longitudinal magnetisation dynamics.

A key drive in the field spintronics is the prospect of developing nanoscale spin-wave logic devices. This would benefit from spin-wave imaging techniques with nanoscale resolution, motivating a push of the spatial resolution of NV-imaging of magnetisation dynamics down to the few-nanometre scale. The challenge is that the spin-wave frequency quickly increases to the  $>10$  GHz regime for wavelengths below  $\sim 100$  nm because of the exchange interaction. As such, the required matching of NV- and spin-wave frequencies (figure 5(c)) limits the smallest detectable wavelengths. Moreover, the exponential decay of the spin-wave stray fields puts a strenuous requirement on the NV-sample distance: for instance, the stray field of a 50 nm spin wave decays exponentially on an 8 nm scale. While scanning-NV magnetometers readily achieve NV-sample distances of  $\sim 50$  nm (see, e.g. [86]), few-nm NV-sample distances have until now only reliably been obtained for samples directly deposited onto diamonds with shallow NV centres. Such shallow NV centres enable the detection of few-nm spin waves with a diffraction limited spatial resolution, and as such we anticipate their future use for probing nanoscale spin-wave physics.

### Current and future use of the technique for material science

After the pioneering study on YIG [103], NV magnetometry has been successfully applied to detect spin-fluctuations in other magnetic materials such as permalloy [104], the ferromagnetic insulator nickel–zinc–aluminium–ferrite [105], the AFM insulator  $\alpha$ - $\text{Fe}_2\text{O}_3$  [106], a synthetic cobalt-based antiferromagnet [107], the topological magnet  $\text{MnBi}_2\text{Te}_4$  ( $\text{Bi}_2\text{Te}_3$ )<sub>*n*</sub> [108], and the low-dimensional magnet  $\text{Fe}_3\text{GeTe}_2$  [109]. For these and many other materials, the magnon band lies above the NV ESR frequencies (figure 6(a)) such that the resonant detection methods of figure 5 are not applicable. Instead, the NV sensing of magnon dynamics relies on detecting the fluctuations of the longitudinal spin density [96]. The longitudinal spin noise is characterised by the two-magnon scattering processes, where a magnon with a frequency  $f + f_{2m}$  can undergo a transition to another magnon with a frequency  $f$ , emitting magnetic noise at frequency  $f_{2m}$  as illustrated in figure 6(a). In contrast to the transverse spin noise with a minimum



frequency determined by the magnon band gap, the frequency bandwidth of its longitudinal counterpart starts from zero, which can be comfortably addressed by NV centres in the low frequency (field) regime.

The longitudinal spin noise is fundamentally related to diffusive spin transport behaviours via Brownian motion [106], which is further connected to the underlying magnetic parameters of the material systems studied. Building on this new measurement platform, one of us used NV relaxometry to show that the intrinsic spin diffusion constant and longitudinal dynamic spin susceptibility of AFM insulator  $\alpha$ - $\text{Fe}_2\text{O}_3$  and vdW magnet  $\text{MnBi}_4\text{Te}_7$  could be directly measured in the absence of external spin biases (figures 6(b)–(d)). The measurement of longitudinal spin noise provides an appealing way to explore local magnetic phase transitions and domain wall motions in magnetic materials. In addition to NV centres, it is also worth mentioning that spin defects hosted by 2D materials, such as hBN, can also be utilised to perform spin relaxometry measurements to detect longitudinal spin noise as demonstrated in  $\text{Fe}_3\text{GeTe}_2/\text{hBN}$  based vdW heterostructures very recently (figures 6(e) and (f)).

Moving forward, promising directions include extending the presented spin-relaxometry sensing platform to magnetic quantum states of matter, such as spin liquids with frustrated magnetic ordering [110], Moiré magnetism with stacking engineered atomic registry and lattice interactions [111], and quantum anomalous Hall insulators with unconventional current flow patterns. We share the optimism that quantum magnetometry techniques could provide previously inaccessible information to explore unconventional spin dynamics at the forefront of condensed matter physics.

## Acknowledgements

C R D acknowledges the support from the Air Force Office of Scientific Research under Award No. FA9550-21-1-0125. T S acknowledges support from the Dutch Research Council (NWO) under Awards VI.Vidi.193.077, NWA.1160.18.208, and OCENW.XL21.XL21.058. A Y Acknowledges support from the ARO Grant W911NF-22-1-0248, by the Quantum Science Center, a National Quantum Information Science Research Center of the US Department of Energy (DOE), Gordon and Betty Moore Foundations EPIQS Initiative through Grant No GBMF9468, the STC Center for Integrated Quantum Materials NSF Grant No. DMR-1231319, and the Aspen Center of Physics supported by NSF Grant PHY-1607611.

### 3. Scanning Hall microscopy

David Collomb<sup>1</sup>, Simon Bending<sup>2</sup> and Ahmet Oral<sup>3</sup>

<sup>1</sup> Quantum Brilliance GmbH, Tullastraße 72, Freiburg im Breisgau, Germany

<sup>2</sup> Department of Physics, University of Bath, Bath BA2 7AY, United Kingdom

<sup>3</sup> NanoMagnetics Instruments Ltd, Suite 290, 266 Banbury Road, Oxford OX2 7DL, United Kingdom

#### Current state of the technique

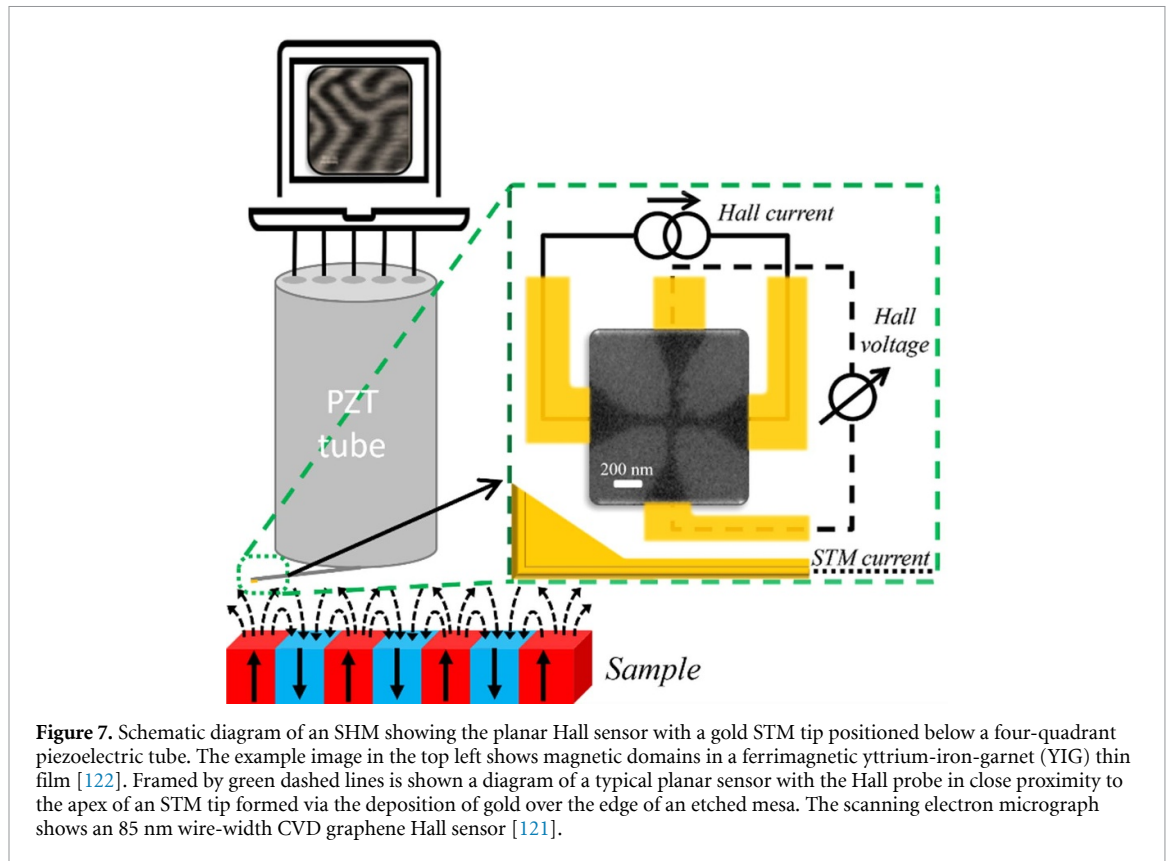
There are many complementary approaches to magnetic imaging and scanning Hall-effect sensors are often chosen because of their high magnetic field sensitivities and wide field range, quantitative linear response, non-invasive performance and fabrication versatility. This allows them to be employed in applications where other semi-quantitative and potentially invasive sensor types with higher spatial resolution, such as MFM cantilevers, may not meet requirements. They are also much more compact and easier to use than the recently developed diamond NV centre magnetic microscope, which requires the precise fabrication of a single crystal diamond with an NV centre at the apex of an AFM tip as well as additional laser and MW excitation. Although the spatial resolution of NV microscopy can in principle approach the atomic scale, imaging is relatively slow due to the need to average the signal at each pixel for typically for a few seconds, and its application is generally limited to low magnetic fields up to a few milli-teslas [112], while SHM is capable of operation up to several teslas.

Magnetic imaging with Hall sensors dates back over 60 years, driven by their quantitative outputs and non-invasive operation. State-of-the-art SHM generally places a Hall sensor within a few micrometres of a STM tip patterned over a deep etched mesa, forming a planar device as shown in figure 7. This is fixed to a fine positioning piezoelectric tube at a shallow angle with the tip closest to the sample surface. The SHM sensor can be brought into sub-nanometre tunnel contact with a sample and then scanned over the surface, collecting the Hall voltage at each pixel. By knowing the Hall coefficient of the sensor, calibrated magnetic field maps at the surface can be constructed. Early examples used Hall sensors a few micrometres in size patterned in thin films of Bi, InAs or epitaxial GaAs. The invention of modulation-doped semiconductor heterostructures [113] led to a renaissance of this field [114, 115] due to their low 2D carrier densities and very high carrier mobilities at low temperatures giving rise to minimum detectable fields,  $B_{\min}$ , of  $< 0.1 \mu\text{T Hz}^{-0.5}$  for  $1 \mu\text{m}$  sensors at  $T \sim 4 \text{ K}$ . GaAs heterostructures remain the choice for cryogenic imaging but their mobilities deteriorate dramatically at room temperature and edge wall depletion limits their minimum dimensions to  $\geq 100 \text{ nm}$  [116]. To reduce the spatial resolution  $50 \text{ nm}$  Bi Hall probes with  $B_{\min} \sim 0.08 \text{ mT Hz}^{-0.5}$  have been patterned in a polycrystalline film using a  $\text{Ga}^+$  focused ion beam [117], however these suffer from poor mechanical and chemical stability. Narrow gap semiconductor sensors have also been widely explored due to their much higher 300 K mobilities. For example, a  $0.5 \mu\text{m}$  InSb probe has been demonstrated with  $B_{\min} \sim 0.72 \mu\text{T Hz}^{-0.5}$  [118]. Although challenging, significant progress has been made in the growth of epitaxial InSb (and other III–V alloy) heterostructures in recent years allowing higher spatial resolutions to be achieved. The isolation of graphene in 2004 initiated a second renaissance in this field. The very low, tuneable carrier density of this atomically thin material and the extremely high 300 K mobilities arising from its unique band structure, make it ideal for nanoscale Hall sensor fabrication [119]. Imaging with  $1.5 \mu\text{m}$  chemical vapour deposition (CVD) graphene sensors has been demonstrated with  $B_{\min} \sim 20 \mu\text{T Hz}^{-0.5}$  [120] while  $85 \text{ nm}$  Hall cross devices with  $B_{\min} \sim 59 \mu\text{T Hz}^{-0.5}$  were subsequently reported [121], both operated at 300 K.

#### Future advances of the technique

The two foremost challenges in the field of Hall microscopy relate to the search for new materials to enable high performance operation at room temperature and the fabrication of sensors with deep sub-micron spatial resolution without severe degradation of minimum detectable fields. High magnetic field resolution requires a large Hall coefficient (low carrier density) and low Johnson and  $1/f$  noise. These criteria are not satisfied in GaAs/AlGaAs heterostructures at 300 K where the carrier mobility is low and  $1/f$  noise dominates the output at low drive currents. Moreover, surface state pinning of the chemical potential leads to sidewall depletion, limiting the minimum practical size of sensors to  $\sim 100 \text{ nm}$  [123]. Although narrow gap III–V semiconductors exhibit much higher 300 K carrier mobilities and have surface states that generally lie in the conduction band and do not lead to carrier depletion, they are typically patterned in films with thicknesses in the range in 200–500 nm, leaving little scope for improved spatial resolution. The growth of epitaxial InSb heterostructures is challenging but in recent years modulation-doped InSb/ $\text{Al}_x\text{In}_{1-x}\text{Sb}$  quantum well structures have been realised with 300 K mobilities up to  $51\,000 \text{ cm}^2 (\text{Vs})^{-1}$  for carrier concentrations as low as  $5.8 \times 10^{11} \text{ cm}^{-2}$  [124]. These structures are typically based on  $\sim 30 \text{ nm}$  wide InSb quantum wells that lie  $\sim 50 \text{ nm}$  below the surface of the wafer and promise excellent performance for Hall microscopy applications



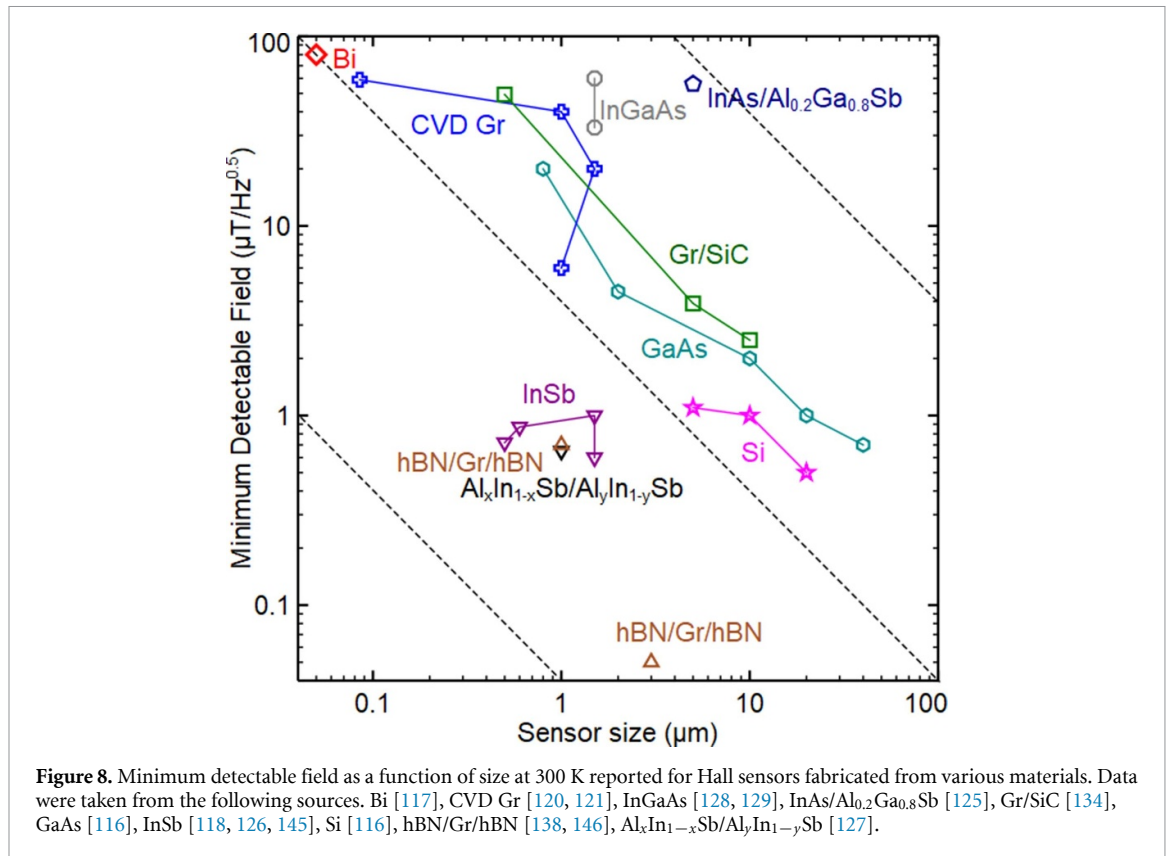


**Figure 7.** Schematic diagram of an SHM showing the planar Hall sensor with a gold STM tip positioned below a four-quadrant piezoelectric tube. The example image in the top left shows magnetic domains in a ferrimagnetic yttrium-iron-garnet (YIG) thin film [122]. Framed by green dashed lines is shown a diagram of a typical planar sensor with the Hall probe in close proximity to the apex of an STM tip formed via the deposition of gold over the edge of an etched mesa. The scanning electron micrograph shows an 85 nm wire-width CVD graphene Hall sensor [121].

with resolutions down to  $\sim 100$  nm. There exist many III–V material options and epitaxial heterostructure sensors with other alloy combinations have also been explored including InAs/Al<sub>x</sub>Ga<sub>1-x</sub>Sb [125, 126], Al<sub>x</sub>In<sub>1-x</sub>Sb/Al<sub>y</sub>In<sub>1-y</sub>Sb [127], as well as AlGaAs/InGaAs/GaAs [128, 129].

Graphene's unique band structure yields very high 300 K carrier mobilities while its single atomic layer form places few limits on nanoscale patterning. The different sources of graphene include CVD graphene [130–133], epitaxially grown graphene on SiC [134] and mechanically exfoliated graphene [135]. CVD graphene is cheap and readily available in very large areas and despite its relatively defective granular structure still exhibits quite high mobilities. Nanoscale CVD graphene Hall sensors have been demonstrated with sizes as small as 85 nm with excellent  $B_{\min} \sim 59 \mu\text{T Hz}^{-0.5}$  at 300 K, patterned on Si/SiO<sub>2</sub> substrates by electron beam lithography and dry etching [121]. These devices could be tuned with a back gate allowing  $B_{\min}$  to be optimised close to the charge neutrality point with scope for further reduction in sensor size down to  $\sim 50$  nm when a rapid decrease in mobility due to edge scattering is expected to set in. CVD bilayer graphene Hall sensors have also been shown to exhibit higher sensitivity than comparable monolayer devices in the absence of substrate bias [136]. Epitaxial growth of graphene on SiC can achieve higher mobilities than CVD material, albeit at a higher cost and lacking the ability to tune the carrier concentration with a back gate. The highest carrier mobilities have been demonstrated in mechanically exfoliated graphene devices. The mobility of structures deposited on Si/SiO<sub>2</sub> substrates is limited by scattering at charge centres in the adjacent oxide layer. This can be suppressed by the encapsulation of graphene in two layers of hBN [137] when 300 K mobilities as high as  $10 \text{ m}^2 (\text{Vs})^{-1}$  have been demonstrated. This represents a very exciting area for future developments and  $B_{\min}$  as low as  $0.7 \mu\text{T Hz}^{-0.5}$  has been estimated in  $1 \mu\text{m}$  encapsulated sensors at 300 K [138]. It remains to be seen if high mobilities can be retained in nanoscale devices, but from a fabrication perspective sizes comparable to the 10–20 nm thickness of hBN capping layers should be achievable. Moreover, breakdown current densities for graphene nanoribbons in excess of  $10^{12} \text{ A m}^{-2}$  have also been demonstrated [139] and just working at 10% of this value would represent  $30 \mu\text{A}$  through a  $1 \mu\text{m}$  wide graphene strip with potentially achievable noise levels in the range  $B_{\min} \sim 0.01\text{--}0.1 \mu\text{T Hz}^{-0.5}$  at 300 K. The very high carrier group velocity associated with the linear electronic dispersion of graphene also enables very high frequency operation and graphene FETs working at up to 46 GHz have been demonstrated [140]. Figure 8 summarises key milestones made in Hall sensor technology, with the recent developments in graphene and its encapsulation showing promise to lead the technique to higher spatial and field resolutions.

SHM requires the integration of a second sensor to control the scan height during imaging and/or map the surface topography. The combination of the two complementary data sets yields new information about



the sample. The most common secondary sensor is an STM tip microfabricated as close to the active Hall element as possible. However, these cannot be used with insulating samples, driving the development of sensors with atomic force tracking. Hall sensors have been integrated onto micromachined cantilevers with piezoresistive deflection detection [141]. Simpler approaches to force tracking have been demonstrated whereby a nanoscale GaAs/AlGaAs Hall sensor was either glued to a GaAs cantilever sandwiched between two piezo plates [142] or to one of the tines of a quartz tuning fork [143]. There have also been several attempts to pattern Hall sensors directly on the end of AFM tips, although this has proved to be very challenging [144]. However, one can anticipate new developments in these areas in order to optimise the imaging spatial resolution.

### Current and future use of the technique for material science

Unconventional type II cuprate high temperature superconductors (HTSs) are notable for having a superconducting transition temperature,  $T_c$ , above 77 K without the need to exceed atmospheric pressure. These have generated significant interest not just for their fascinating physics, but also for utilisation in energy-based applications such as lossless transfer of electricity from remote sources to their point of consumption. A major challenge in this development involves maximising the critical current density of manufactured HTS tapes through optimising the flux pinning without degrading  $T_c$ . SHM is well positioned to study these materials due to its quantitative nature, as well as its variable temperature operation from milli-Kelvin temperatures to well above the highest cuprate  $T_c$ . By mapping the pinning of vortices to defect sites as well as quantitatively measuring the temperature dependent vortex profiles, SHM provides critical feedback for the optimisation of flux pinning [122, 147–149]. The miniaturisation of Hall probes down to the nanoscale using graphene will enable an expansion of this research by improving the correlation between vortices and nanoscale pinning sites. This is in addition to being able to access higher magnetic field and applied current regimes where higher resolution is required to resolve the densely packed vortices.

Of particular interest for materials science today are magnetic skyrmions. These topologically protected magnetic features, which can be as small as 1 nm in diameter, are predicted to represent discrete energetically favourable magnetic states. This has led skyrmion-based logic and memory devices such as nanoscale racetrack-type technology or ferroelectrically tuneable skyrmions in oxide heterostructures [150, 151] to be envisioned. Which materials skyrmions manifest in and how they behave in their host material are questions not yet fully understood. To tackle these skyrmions need to be visualised via real-space imaging, where the ability to resolve skyrmions with sizes up to 100 nm is required. The recent improvements in this resolution

range using graphene sensors demonstrates SHM's ability to play an important role in this field in the future, especially in materials that exhibit skyrmions at room temperature.

By inverting the conventional set-up, whereby the Hall probe remains stationary and the magnetic object now rasters above the probe, one can use SHM to characterise an MFM tip's stray field, which allows for quantitative information to be extracted out of the otherwise qualitative MFM images [152]. Here, the tip is scanned over the active Hall probe and at each scan pixel a corresponding Hall voltage is measured. While the measured Hall voltage is dependent on the tip–sensor displacement, the magnetic moment of the tip, and the tip oscillation amplitude, an additional unwanted contribution arises from the electrostatic interaction between tip and probe. This unwanted interaction can be removed by also operating in a force modulated Kelvin probe force mode. With the recent advances in graphene-based Hall probes, this technique could be readily rolled out for calibration of commercially available AFM systems, allowing a user to unlock with ease quantitative MFM with high spatial resolution.

The growing complexity and scale of industrial manufacturing processes and infrastructure from metal bridges to pipework has placed a critical demand for a new arsenal of non-destructive testing (NDT) tools. By exploiting the continued miniaturisation of sensors, NDT tools can become capable of mapping defects and structural changes down to the nanoscale. Scanning Hall probes have also been used as an NDT tool through scanning Hall Susceptometry. For example, a 200  $\mu\text{m}$  spatial resolution SHM system was combined with a small electromagnet for the monitoring of heat-resistant steels as they age, where the microstructural changes correlated with a change in magnetic response [153]. However, such designs have relied upon micrometre resolution Hall probes, as well as using an excitation coil surrounding the sample. To advance the technique and its applications, nanometre spatial resolutions and the move to 'local' field excitations are desirable. By combining the recent developments in graphene-based probes with previous work integrating a field excitation coil and a Hall probe on the apex of an AFM tip [144], SHM could become a routine inspection tool in various manufacturing processes such as for the control of magnetic storage media or structural monitoring of metallic artefacts.

## Acknowledgements

D C and S J B acknowledge financial support from EPSRC in the UK under Grant Number EP/R007160/1 and the Superqumap COST Action CA-21144. D C was also supported by a PhD studentship from Lloyds Register Foundation ICON (Award Nos. G0086).



## 4. Magnetic force microscopy

Hans J Hug<sup>1,2</sup>, Andrada-Oana Mandru<sup>1</sup>, Volker Neu<sup>3</sup>, Hans Werner Schumacher<sup>4</sup>, Sibylle Sievers<sup>4</sup> and Hitoshi Saito<sup>5</sup>

<sup>1</sup> Empa, Materials Science and Technology, Ueberlandstrasse 129, 8600 Duebendorf, Switzerland

<sup>2</sup> Department of Physics and Swiss Nanoscience Institute, University of Basel, Klingelbergstrasse 82, 4056 Basel, Switzerland

<sup>3</sup> Leibniz Institute for Solid State and Materials Research, Helmholtzstrasse 20, D-01069 Dresden, Germany

<sup>4</sup> Physikalisch-Technische Bundesanstalt, 38116 Braunschweig, Germany

<sup>5</sup> Department of Mathematical Science and Electrical-Electronic-Computer Engineering, Graduate School of Engineering Science, Akita University, Akita, Japan

### Current state of the technique

MFM is a laboratory-based scanning probe technique that uses a magnetic tip on an oscillating cantilever to image magnetic stray fields emanating from a sample with a spatial resolution down to 10 nm [154]. MFM is a versatile tool for characterising materials as it can simultaneously map the topography and local Kelvin potential [155] apart from the stray field, requires minimal sample preparation and can image magnetic layers covered by oxidation protection layers or thinner top contacts. Additionally, MFM can be performed under various conditions, including high magnetic fields and variable and cryogenic temperatures [156], and in liquid environments [157].

Typically, MFM is performed under ambient conditions using a line-by-line two-pass measurement method (figure 9). During the first pass, the tip intermittently contacts the sample while utilising a feedback to trace the topography. In the second pass, the tip is lifted off the surface to predominantly measure the longer-ranged tip-sample interaction forces. These include the desired magnetic force on the tip that can be expressed as the convolution between the tip magnetisation  $\mathbf{M}_t$  and the derivative of the sample stray field  $\mathbf{H}_s$  [154]:

$$\mathbf{F}_t(\mathbf{r}, z) = \mu_0 \int_{V_t} [\mathbf{M}_t(\mathbf{r}', z') \cdot \nabla] \mathbf{H}_s(\mathbf{r} + \mathbf{r}', z + z') d\mathbf{r}' dz' \quad (1)$$

Magnetically-hard tips, where the magnetisation  $\mathbf{M}_t$  can be considered as fixed, are generally used. The force acting on the tip is detected via the phase shift signal of the oscillating cantilever,  $\Delta\phi \propto dF_n/d\mathbf{n}|_{\text{eff}}$ , where  $\mathbf{n}$  represents the canted cantilever normal and also the canted direction of the oscillating cantilever. The index '|<sub>eff</sub>' indicates that the derivative  $dF_n/d\mathbf{n}|_{\text{eff}}$  is a weighted average of the local derivative over the oscillation amplitude of the tip [154].

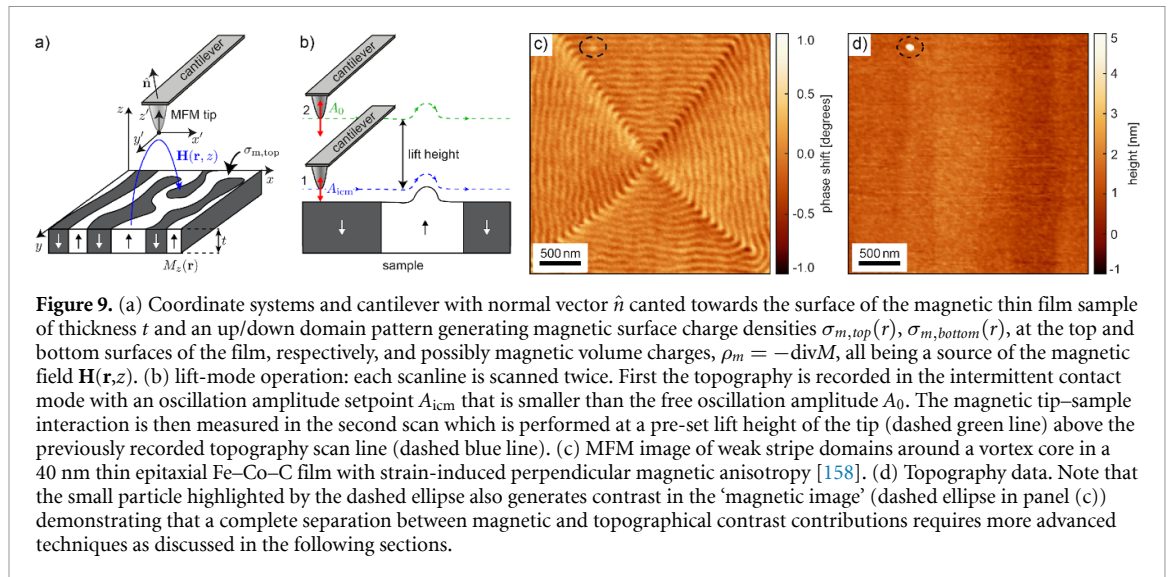
Note that since  $\mathbf{F}_t$  arises from a sample field-tip magnetisation convolution, it depends on the spatial wavelength of the stray field pattern. Without calibration, the interpretation of results remains qualitative. It is also worth noting that an increased magnetic tip volume (e.g. thicker coating) increases the signal amplitude as per equation (1). However, this comes at the cost of spatial resolution, and an increased risk of altering the sample's micromagnetic state by the stray field of the tip.

Therefore, further advancements in MFM require improvements in SNR, measurement functionality, data analysis, tip calibration procedures, and the development of more advanced magnetic probes.

### Future advances of the technique

The following paragraphs detail proven methods to unlock the complete potential of MFM. Presently, these methods are adopted by only a few research groups focused on advancing MFM. To make these advancements accessible to a broader research community, there is a need for instruments that are commercially available, along with standardised metrology procedures, software, and MFM tips that incorporate these advanced techniques.

*Improved sensitivity and advanced operation modes:* working at minimum tip-sample distance during scanning is key to a large SNR. Furthermore, sensitivities improve by up to 100× by operation in vacuum, using cantilevers coated only at the tip-end and by using softer silicon cantilevers. Fields down to a few tens of  $\frac{\mu\text{T}}{\sqrt{\text{Hz}}}$  [159] can then be measured even with low magnetic moment tips (figure 10(a)). However, more complex instrumentation, such as a phase-locked loop to drive the cantilever on resonance, along with non-contact operation and tip-sample distance control modes, are needed. An operational mode that provides extremely robust non-contact tip-sample distance control is based on the modulation of the tip-sample electric potential [160]. This leads to sidebands at  $f_0 \pm f_{\text{mod}}$  and  $f_0 \pm 2f_{\text{mod}}$ , which can be used to map the local Kelvin potential by Kelvin Probe Force Microscopy (KPFM) and for a (slow) control of the



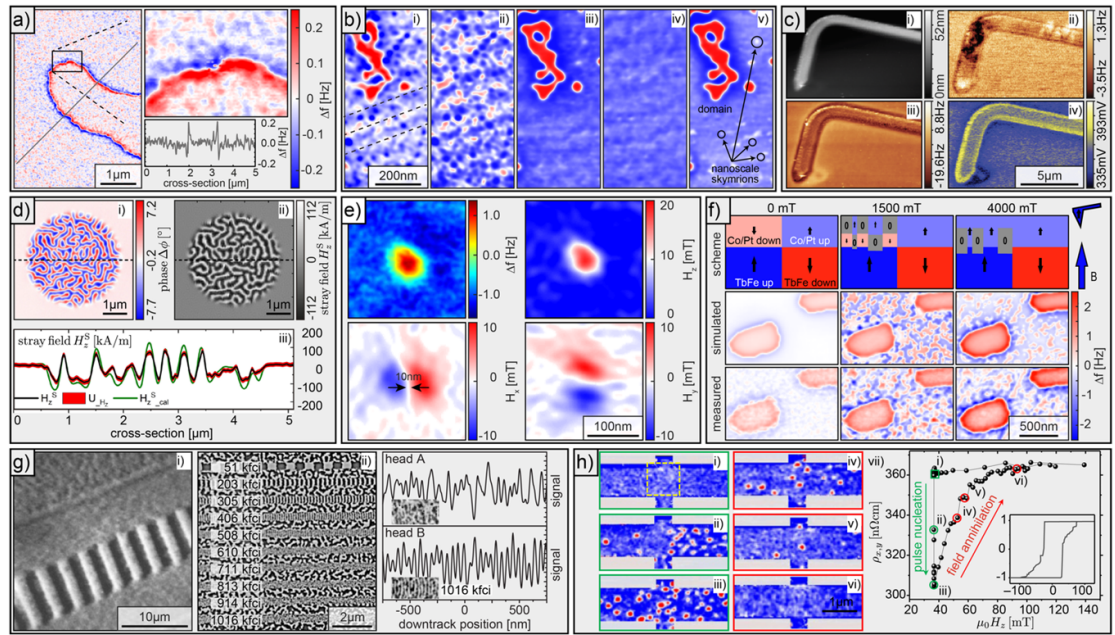
tip-sample distance. This enables quantitative comparison of data acquired on different samples, as needed for quantitative MFM as an example. Differential imaging (figure 10(b)) [156] and KPFM techniques can be used to disentangle various contributions to the measured contrast (figure 10(c)). In addition, a further miniaturisation of the cantilever or the use of nanowire cantilevers is a promising approach to optimise the stiffness-to-resonance frequency ratio and thus the force sensitivity.

*Tip-response calibration and quantitative measurements:* a calibration of the tip’s response to a range of spatial stray field wavelengths [161, 162] (figure 10(d)) allows obtaining the vector magnetic field (figure 10(e)) from the measured MFM signal [161] or testing the fidelity of candidate micromagnetic structures (figure 10(f)) [156]. Such a tip calibration can be performed by measuring a reference sample [161, 162] or by directly measuring the tip’s stray field distribution [163]. Together with the disentanglement of different contrast contributions, quantitative MFM is a powerful method to resolve nanoscale magnetic textures (figures 10(b) and (f)) [156]. Future developments should include calibration in an extended range of spatial frequencies to address multiscale problems, and field-dependent calibrations. However, note that for advanced tip calibration and deconvolution procedures [162] high SNR data are beneficial. This is achievable with high-quality factor cantilevers, which are best operated under vacuum conditions [159], and requires the use of advanced MFM instrumentation.

*Advanced magnetic probes:* high-resolution imaging is typically performed with high-aspect ratio tips with tip radii below 5 nm, coated with a few nanometres of ferromagnetic material on one side. However, the magnetic stray field of the tip can still be too large, leading to changes in the sample’s micromagnetic state. In such a case, the magnetic moment of the tip must be reduced, which can however not be achieved by further lowering the tip’s coating thickness. Instead, alternative approaches, such as the use of low-moment magnetic alloys or the design of more sophisticated tip coatings, consisting of two antiferromagnetically-coupled ferromagnetic layers can be explored to further advance the technique.

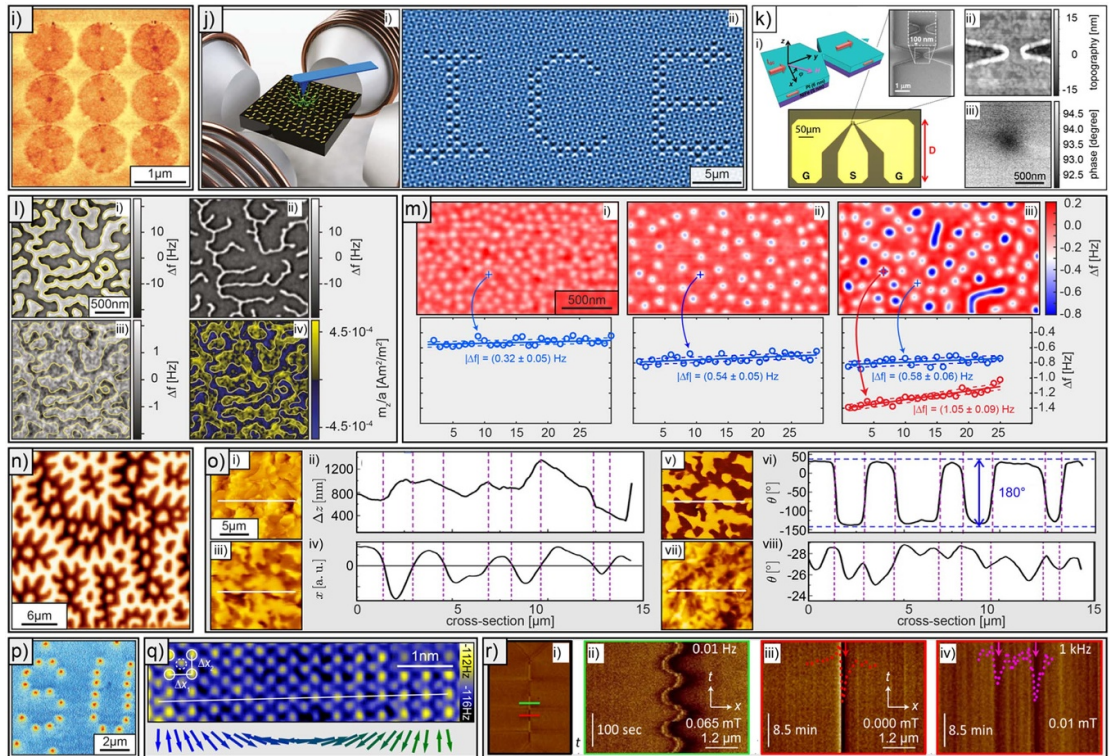
Generally, tips permitting an adjustment of their magnetic state are beneficial for MFM [159]. Although the magnetic state of high-aspect ratio magnetic tips can be reliably flipped, it would be advantageous to change the tip’s magnetic state at higher frequencies, potentially even in the presence of the sample. This would aid in separating magnetic contrast contributions from those arising from the sample topography.

One promising approach is the use of tips coated with a superparamagnetic granular alloy (SP tips), which allows for the adjustment of the size and direction of the tip’s magnetic moment with an external field. The tip moment is proportional to the external field and follows it without hysteresis [164]. This technique, known as alternating MFM (A-MFM) (figure 10(o)), significantly improves spatial resolution even at ambient conditions, provided that the sample magnetisation is stable in the applied oscillatory field. In addition, A-MFM permits mapping the gradient of the sample vector field along the direction defined by the applied ac-field driving the tip with high precision [164]. Moreover, concepts to measure the magnetic field value including uniform field component are presently explored.



**Figure 10.** (a) Advanced tip and cantilever design: overview image acquired with a vacuum MFM using an ultra-high quality factor cantilever shown together a zoomed image (shown right) acquired at the location of the rectangular box, and a cross-section. The red/blue frequency shift pattern arises from a local variation of the areal density of the magnetic moment corresponding to a variation in magnetic layer thickness by  $\pm 1$  atomic monolayer [159]. (b) Disentangling topography and magnetism: MFM data nanoscale skyrmions (panel (v)) in  $\text{SrIrO}_3/\text{SrRuO}_3$  bilayers observed by MFM after a careful subtraction of the topographic (panel (ii)) and magnetic (panel (iv)) backgrounds from the raw MFM data (panel (i)) [156]. (c) Disentangling topography, magnetism, and Kelvin contrast: Co nanowire topography, Kelvin potential, and frequency shift image. The magnetic domain pattern could be revealed by compensating electrostatic force. Reproduced from [155]. CC BY 2.0. (d) MFM tip calibration: MFM image of a patterned Co/Pt multilayered thin film (i) and thereof calculated stray field distribution at  $z = 60$  nm (ii). Simulated and calibrated stray field data with expanded uncertainties are compared in the cross-section image (iii) [161, 162]. (e) Quantitative MFM—from  $\Delta f$  to the vector components of the stray field: MFM data of a skyrmion in a  $[\text{Ir}(1)/\text{Co}(0.6)/\text{Pt}(1)]_{x5}$ -multilayer measured  $12.0 \pm 0.5$  nm above the surface (i). Vector components,  $H_z$ ,  $H_x$ , and  $H_y$ , of the magnetic field obtained from the deconvolution of a using the tip calibration function [192]. (f) Quantitative MFM—simulation of MFM contrast from candidate micromagnetic structures: simulated (middle row) and measured (bottom row) MFM data of exchange-coupled double layers; the simulated MFM data was obtained by using the candidate magnetisation structures shown in the top row [193]. (g) Data storage applications: (i)  $8 \mu\text{m}$  wide track on a longitudinal magnetic recording medium (CoPtCr) [169]. (ii) Overview of tracks written by high-density magnetic recording head B. In contrast to head A, head B permits the writing of a 1016 kfcI track (see track cross-sections right to the image) [194]. (g) Reprinted from [169], with the permission of AIP Publishing. (h) Spintronic applications: varying number of individual skyrmions in a Hall cross and corresponding transverse resistance  $\rho_{xy}$ . Reproduced from [195], with permission from Springer Nature.





**Figure 10. (Continued.)** (i) Magnetic vortices in nanostructures: MFM data obtained on vortices in circular permalloy islands with  $1\ \mu\text{m}$  diameter and  $50\ \text{nm}$  thickness with different core polarities. From [174]. Reprinted with permission from AAAS. (j) Spin-ice: experimental setup (i) and result of the local magnetic state manipulation (ii) of a magnetic charge ice. Reproduced from [175]. [CC BY 4.0](#). (k) MFM on an operational spin Hall nano-oscillator: (i) schematic (left top), SEM image of a nanoconstriction (right top), and optical microscopy image of the final device in a conventional design with top contact coplanar waveguides (with length of  $D$ ) providing electrical access to the SHNO (bottom). Topography (ii) and MFM image (iii) of an SHNO with a constriction width of  $300\ \text{nm}$ , operated at  $H = 800\ \text{Oe}$  and  $I_{\text{dc}} = 6\ \text{mA}$ . Reproduced from [176]. [CC BY 4.0](#). (l) Mapping pinned uncompensated spins in antiferromagnet/ferromagnet exchange-biased samples: MFM data obtained at  $8.3\ \text{K}$  on a  $\text{Si}/\text{Pt}(20)/\text{Pt}(2)[\text{Co}(0.4)/\text{Pt}(0.7)]_{20}/\text{Co}(0.6)/\text{Co}(1)$  exchange biased sample (layer thicknesses in nm). Panel (i) after zero-field cooling showing the ferromagnetic domain pattern highlighted by the yellow lines. (ii) Domains in an applied field of  $200\ \text{mT}$ , and (iii) saturated state. The yellow lines (outlines of the initial domain pattern) reveal that below a white domain there is a granular pattern with a dark contrast. (iv) Sub-monolayer areal density of the pinned uncompensated spins at the antiferromagnet/ferromagnet interface obtained from a quantitative analysis of the MFM frequency shift data depicted in (iii) [177]. (m) Reproducible tip-sample distance control for quantitative MFM: determining the location of different skyrmion types (strong and weak  $|\Delta f|$  contrast—panel (iii)) within ferromagnetic/ferrimagnetic/ferromagnetic trilayers by performing MFM at the same distance and with the same (unaltered) tip on a subset of layers (panels (i) and (ii)) from the original trilayer [179]. (n) MFM on permanent magnet materials: domain contrast on bulk  $\text{Sm}_2\text{Co}_{17}$  grain imaged with a high coercive ( $\mu_0 H_c > 0.8\ \text{T}$ )  $\text{SmCo}_5$  tip (DOI: 10.1039/c8nr03997f). (o) A-MFM imaging of the topography and magnetic fields on rough surfaces using superparamagnetic tips: magnetic domain imaging of a very rough, fractured surface of Sr ferrite magnet without topographic crosstalk by A-MFM imaging of the topography and magnetic fields on rough surface using superparamagnetic tips (i), (ii) topographic image and its line profile of grain boundary fractured surface of Sr ferrite magnet. (iii), (iv) Magnetic field gradient image and its line profile with fixed measuring direction. (v), (vi) Magnetic phase image and its line profile.  $180^\circ$  phase difference corresponds to the up and down magnetic field direction. (vii), (viii) Conventional MFM image and its line profile, with uncertain measuring direction [164]. (p) Manipulation of the micromagnetic state: manipulation of vortices in Nb thin film and observation at  $5.5\ \text{K}$ . Reprinted from [196] with the permission of AIP Publishing. (q) Magnetic exchange force microscopy: atomic resolution exchange force frequency shift image of a Mn monolayer on  $\text{W}(110)$  measured at constant height relative to a reference tip-sample distance obtained by STM operated at  $U_{\text{bias}} = -10\ \text{mV}$  and  $I = -2\ \text{nA}$  reduced by  $0.29\ \text{nm}$ . As a cantilever a tuning-fork sensor is used oscillated with an amplitude of  $50\ \text{pm}$ . The bottom part shows a sketch of one half of the cycloidal spin spiral along the  $[154-163]$  direction. Reproduced from [189]. [CC BY 4.0](#). (r) Imaging dynamic phenomena: domain wall oscillation in a permalloy Landau structure. (i) Demagnetized state, (ii) real-time imaging of the quasi-static domain wall oscillating with  $0.01\ \text{Hz}$  (iii) static DW in the  $x-t$  plane, (iv) time-averaged imaging of the DW oscillating with a frequency of  $1\ \text{kHz}$  ( $\mu_0 H_{\text{ac}} = 0.10\ \text{mT}$ ) in the  $x-t$  plane [168].

Freitag *et al* [165] have developed a new cantilever design that uses ultra-soft bending modes and higher flexural modes to simultaneously measure stray fields and their derivatives at the sample surface. This approach can improve surface analysis in materials science, nanotechnology, and biophysics.

*Imaging dynamic magnetic processes:* as a mechanical scanning technique and because of the small signal detected, MFM imaging is limited to line rates of a few Hertz. Imaging dynamic magnetic phenomena thus appears challenging. However, side-band or A-MFM enables the imaging of temporally oscillating processes, such as the magnetic excitation of superparamagnetic particles [166] or high-frequency fields from a recording head [167]. Recently, time-averaged MFM has also been used to image spatially varying processes, including nanoscale domain wall movements oscillating at frequencies up to 1 kHz, with the potential for measuring at hundreds of kHz in the future [168].

### Current and future use of the technique for material science

Each of the following subsections describes typical current application fields and then points out future applications that become accessible through a wider availability of more advanced MFM instrumentation or a further development of the MFM technique as outlined in the previous section.

*Data storage applications:* the initial MFM development was boosted by the need to characterise magnetic recording media (figure 10(g) (i)) [169], to investigate the impact of specific head designs on the bit structure and storage density, and study magnetisation reversal processes in longitudinal and perpendicular hard disk media (figure 10(g) (ii)). MFM was essential for the development of advanced data storage methods such as heat or MW assisted recording [170] and domain-wall-based racetrack devices [171]. More recent research topics address topologically protected non-collinear magnetic structures (figure 10(h)) [172].

*Characterising the micromagnetic state of thin film multilayers:* MFM has become the standard method for characterising magnetic thin films and multilayers [173], unravelling the physics governing their micromagnetic state and its dependence on applied magnetic fields and temperature (figures 10(b) and (f)). Patterning film materials introduces new anisotropies, functionalities, and applications. MFM is applied to studying closure domain structures and magnetic vortices (figure 10(i) [174], figure 9(c)), and could reveal monopole-like excitations in artificial spin ice (frustrated magnets) (figure 10(j)) [175] and the spatial profile of a spin-nano-oscillator (figure 10(k)) [176].

*Materials and system design based on the quantitative MFM analysis of the micromagnetic state:* through quantitative MFM, all vector components of the sample's stray field can be measured [154, 162] and the comparison to micromagnetic models enables quantitative materials research [156]. An example is the correlation of the areal density of pinned uncompensated spins at the interface between ferromagnetic and AFM layers with the exchange-bias field (figure 10(l)) [177]. More recent quantitative MFM work in combination with micromagnetic simulations was performed to reveal the various reversal mechanisms of 3D metallic multilayers exchanged-coupled to rare-earth-transition metal alloys for spintronic applications (figure 10(f)) [178]. Regarding topological chiral structures, MFM could determine the different soliton types present in these systems (figure 10(m)) [179]. Very recent MFM work also addressed 2D materials, where different magnetic skyrmion types were investigated at low temperature in vdW heterostructures [180].

Permanent magnetic materials (figure 10(n)) have been studied by A-MFM which is suitable for distinguishing topographical and magnetic contrast contributions relevant of these sintered materials with magnetic degradation-free grain boundary fractured rough surface (figure 10(o)) [164]. Moreover, to measure the absolute local field which can distinguish magnetic and non-magnetic grain boundary phases will become available in the future.

Combining MFM with further techniques like Kelvin probe microscopy to eliminate electrostatic contributions enables a quantitative characterisation of current distributions via the detection of Oersted fields to characterise leakage currents in integrated circuits for device failure analysis [181] and the local sensitivity of sub-micron Hall-crosses [182]. MFM was also combined with *in-situ* modulating electric fields, to image both magnetic and ferroelectric domains in multiferroic systems [183]. More recently, combining MFM with piezo-response force microscopy and STEM lead to the realisation of magneto-electrical-optical coupling in BiFeO<sub>3</sub>-based thin films [184]. MFM is therefore expected to play a key role in novel low power consumption ME memories research.

*Manipulation of sample states by the tip field:* from the beginning of MFM, the tip stray field has also been used to locally manipulate nanoscale magnetic objects or magnetisation structures, for example to nucleate

and manipulate vortices in superconductors (figure 10(p)). Recent work addressed the nucleation and manipulation of skyrmions in chiral materials [185]. With advanced scanning control, these activities are expected to have significant impact in future studies of magnetic nanoparticles [186], artificial spin ice [175], and reconfigurable magnonic crystals [187].

*MFM and magnetic exchange force microscopy (MExFM) in ultra-high vacuum (UHV)* (see also chapter 5): MFM performed under UHV conditions permits the study of reactive surfaces or magnetic system which must be prepared on single crystalline surfaces, which may gain importance particularly if combined with atomic-resolution AFM or Kelvin probe microscopy techniques. Furthermore, in UHV, MExFM can characterise magnetism on the atomic scale to image, for example, the AFM unit cell in NiO [188] or, more recently, the cycloidal spin spiral in one monolayer of Mn on W(110) [189] (figure 10(q)).

*Dynamic phenomena studied by MFM*: existing time-averaged and time-resolved MFM techniques allow studying fundamental aspects of nanoscale domain wall oscillations (figure 10(r)) [168]. These are not only relevant in novel magnetic nanodevices such as domain wall resonators [190] and memory devices, but also in sensor materials based on a domain wall-based magneto impedance effect [191] and in classical soft magnetic materials for transformer applications. Further advances in imaging dynamic magnetic processes would arise if one could rapidly switch the magnetic state of the tip on a time scale down to a few nanoseconds using optical triggering, independent from the state of the sample. This will allow for the implementation of highly resolving stroboscopic imaging techniques, thus advancing the development of modern spintronic devices.

## Acknowledgements

The A-MFM measurement system was developed by the support of Japan Science and Technology Agency (JST)/SENTAN Program in 2008–2015.

## 5. Spin sensing on the atomic scale with STM and non-contact AFM

A A Khajetoorians<sup>1</sup>, N Hauptmann<sup>1</sup>, S Baumann<sup>2</sup>

<sup>1</sup> Scanning Probe Microscopy Department, Institute for Molecules and Materials, Radboud University, Nijmegen, The Netherlands

<sup>2</sup> University of Stuttgart, Institute for Functional Matter and Quantum Technologies, Stuttgart, Germany

### Current state of the technique

Spin sensing on surfaces with atomic-scale resolution has been achieved using two different, but compatible methods: STM, and non-contact AFM (nc-AFM) [203]. STM-based spin sensing methods rely on local detection of a spin-component of the tunnelling current, whereas nc-AFM-based spin sensing methods rely on the detection of local forces related to magnetic exchange. SP-STM is the magnetic sensing technique based on STM/spectroscopy that uses a probe with a non-zero magnetisation. This leads to an additional contribution to the measured tunnelling current that is sensitive to the projection of the tip and sample magnetisation, analogous to a magnetic tunnel junction. The method usually requires an applied external magnetic field to measure changes of this projection to infer the change in magnetisation. SP-STM has been used to image and manipulate magnetic order on a variety of metallic surfaces [204, 205]. Most recently, it has been used at sub-Kelvin to mK temperature to measure the magnetisation of individual and coupled atoms [199, 206] as well as to measure spin dynamics [201, 202, 207].

MExFM is a variant of nc-AFM which is based on measuring the short-range magnetic exchange force between a magnetic probe and sample [188]. This is in contrast to MFM, which is sensitive to long-range magnetic dipole forces. MExFM detects the local exchange force between the probe and sample, and therefore is a short-range method that can atomically resolve the magnetic ground states of materials, independent of current, but also be combined with SP-STM detection (figure 11(a)) [197, 208]. MExFM can be also used to quantify the distant-dependent exchange forces between a tip and sample [189, 208, 209].

While both SP-STM and MExFM provide spatial resolution down to individual atoms, historically they suffered from poor time resolution (>microseconds). However, over the last decade, new capabilities have been developed to extend the spatial and temporal resolution, currently reaching all the way to the picosecond-range [210]. These include all-electrical pump-probe spectroscopy, and stochastic resonance spectroscopy [200, 207, 211]. Additionally, ESR techniques have been developed in order to coherently excite atomic spins on surfaces (figure 11(f)) [201, 202]. In terms of magnetic sensing, the all-electrical based methods have been used to extract weak magnetic interactions, as well as extract the intraatomic hyperfine interaction of an individual atom [212, 213].

### Future advances of the technique

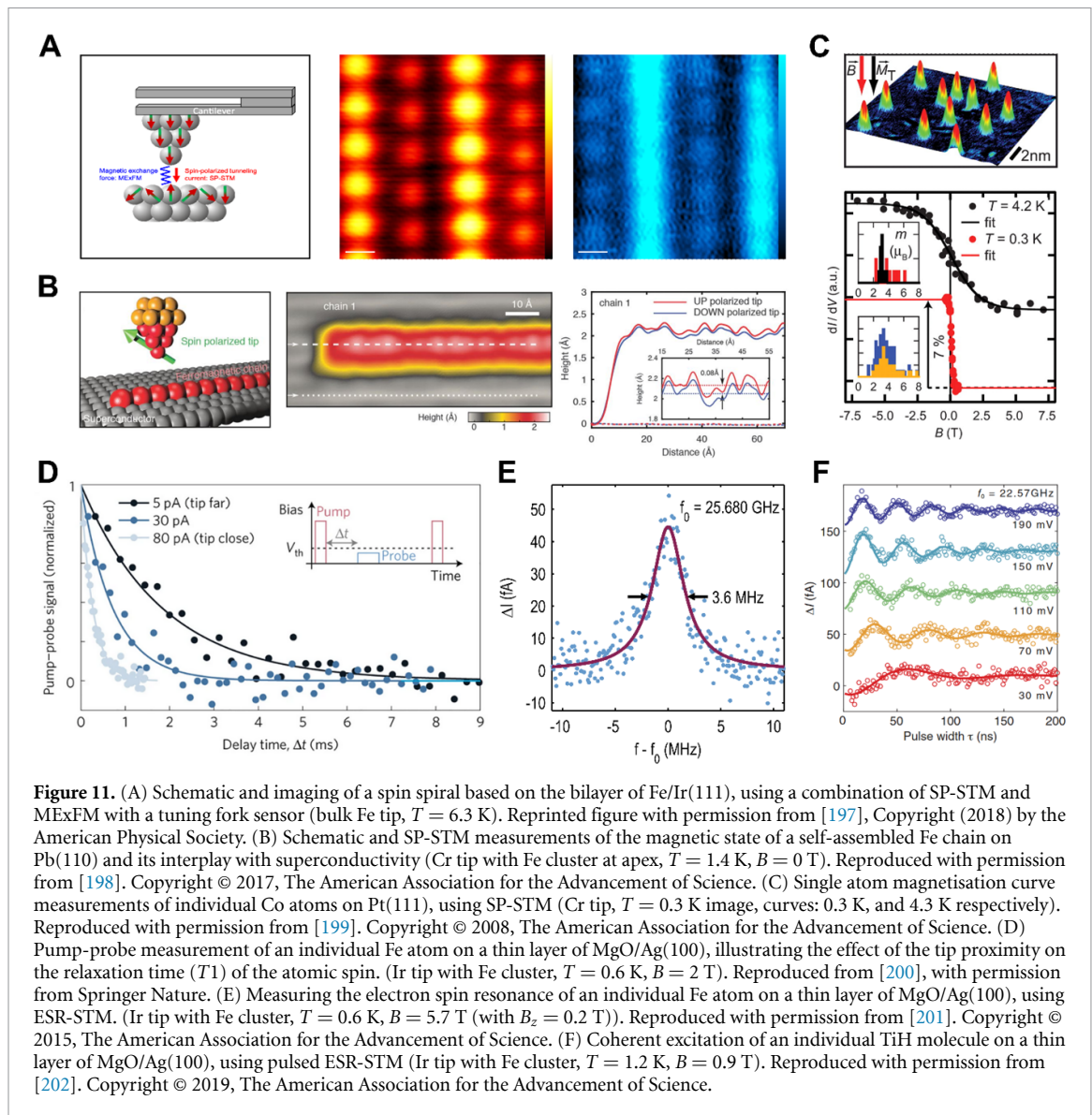
The future of magnetic sensing at the ultimate smallest length-scales involves pushing beyond simple imaging for example by adding new physical understanding by changing the detection schemes for MExFM or by extending the temporal resolution for SP-STM. Furthermore, integration with other methods will be a key advancement for either technique. Here we focus specifically on: (1) instrumental development of MExFM, (2) integration of spin sensing with optical time-resolved techniques, (3) linking local spin sensing with other methods in single measurement setups.

### MExFM

An advancement of the MExFM method would be its implementation to <sup>3</sup>He temperature and below to access fragile quantum states of matter that emerge at low temperature. However, there are multiple hurdles towards this end based on signal detection. MExFM is based on a tuning fork design in order to facilitate simultaneous current detection, and requires signal amplification near the cold junction. But conventional electronics require higher temperature to properly operate, and typically cannot withstand high magnetic fields.

In addition, further developments need to be made on the magnetic probe used in the method. Up to date, primarily ferromagnetic probes have been used [188, 208, 214], for example 3d transition metals [188]. The use of ferromagnetic probes strongly limits the ability to detect robust antiferromagnets on surfaces, as well as to investigate the magnetic structure without additional perturbations driven by large stray fields. Thus, using AFM probes for exchange force detection, with minimal stray fields and a much more robust stable magnetic structure, will be key. Moreover, further developments to measure the distance-dependent exchange force are necessary and depend on controllable functionalisation of the tip. Here, tip functionalisation, as done with CO molecules to enhance spatial resolution [215], will allow to tune the





**Figure 11.** (A) Schematic and imaging of a spin spiral based on the bilayer of Fe/Ir(111), using a combination of SP-STM and MEXFM with a tuning fork sensor (bulk Fe tip,  $T = 6.3$  K). Reprinted figure with permission from [197], Copyright (2018) by the American Physical Society. (B) Schematic and SP-STM measurements of the magnetic state of a self-assembled Fe chain on Pb(110) and its interplay with superconductivity (Cr tip with Fe cluster at apex,  $T = 1.4$  K,  $B = 0$  T). Reproduced with permission from [198]. Copyright © 2017, The American Association for the Advancement of Science. (C) Single atom magnetisation curve measurements of individual Co atoms on Pt(111), using SP-STM (Cr tip,  $T = 0.3$  K image, curves: 0.3 K, and 4.3 K respectively). Reproduced with permission from [199]. Copyright © 2008, The American Association for the Advancement of Science. (D) Pump-probe measurement of an individual Fe atom on a thin layer of MgO/Ag(100), illustrating the effect of the tip proximity on the relaxation time ( $T_1$ ) of the atomic spin. (Ir tip with Fe cluster,  $T = 0.6$  K,  $B = 2$  T). Reproduced from [200], with permission from Springer Nature. (E) Measuring the electron spin resonance of an individual Fe atom on a thin layer of MgO/Ag(100), using ESR-STM. (Ir tip with Fe cluster,  $T = 0.6$  K,  $B = 5.7$  T (with  $B_z = 0.2$  T)). Reproduced with permission from [201]. Copyright © 2015, The American Association for the Advancement of Science. (F) Coherent excitation of an individual TiH molecule on a thin layer of MgO/Ag(100), using pulsed ESR-STM (Ir tip with Fe cluster,  $T = 1.2$  K,  $B = 0.9$  T). Reproduced with permission from [202]. Copyright © 2019, The American Association for the Advancement of Science.

exchange mechanisms responsible for the measurement and to quantify the magnetic exchange energy landscape between different chemical terminations of the probe and the surface.

### Combinations with time resolution

STM/AFM methods have been able to sense fast magnetic switching times. However, currently spin sensing is still limited to electrical detection, and thus in the range of 100 ps. Yet, there are possibilities to improve these magnetic sensing capabilities to extend the time resolution by orders of magnitude. In particular, methods using ultra-fast THz [210] or other fs-optical pulses as excitation pulses, to excite optical transitions, vibrational modes or phonons, combined with the magnetic sensitivity and high spatial resolution in STM promise new insight into ultra-fast processes in materials and can even give access to unknown, short-lived transient states. Similarly, experiments that involve light detection, rather than excitation, allow for studying new material properties once magnetic sensitivity is added to the existing capabilities. The combination of electrical and optical excitation/detection will create unique possibilities for example to study the dynamics of coupled modes, such as magnon–phonon coupling. Furthermore, extending the use of frequency-resolved methods, such as stochastic resonance spectroscopy or ESR-SPM (figure 11(a–c)) [201, 202, 211, 216], will allow for new ways to map out driven dynamics in complex materials with sub-Angstrom resolution. Thereby short-lived excited states of matter, that are not easily detectable with established methods, can be probed and utilised.

### Hybrid measurements with spin sensing

While spin sensing with STM/nc-AFM based methods boasts fantastic spatial and time resolution, there are extreme demands on the measurement conditions. Due to the sensitivity of the probe, samples are typically



studied in UHV environments and require atomically clean surfaces to probe, often restricted to single crystalline surfaces (figure 11(d–e)). These combinations make it extremely challenging to apply more advanced sample fabrication, such as systems that are lithographically defined. Therefore, one of the developments in the field includes being able to explore a vaster amount of samples, such as heterostructures synthesised *ex situ* and *in situ*. Moreover, there is a continuing trend in combining STM and nc-AFM based methods, with others, such as transport-based measurements. Future spin sensing developments in this realm will also involve computability with such methods, as well as integration with other methods reviewed in this roadmap.

### **Current and future use of the technique for material science**

The current state-of-the-art as well as many of the cutting-edge developments outlined here related to spin sensing with STM/nc-AFM-based methods will be extremely valuable for future fundamental characterisation in a variety of material systems. A key advantage of these scanning probe methods over most other magnetic sensing techniques is their unprecedented spatial resolution which adds local information to the processes present in complex materials that are otherwise impossible to resolve. This local information can bring understanding to the influence of defects and other impurities in materials, for example in their role as pinning centres or photon emitters, or it can be used to resolve local structural rearrangements, e.g. due to correlated electronic states. Thus, this information can guide future developments of materials with specific functionality driven by their local composition and/or their local response to time-dependent excitation.

### **Lower dimensional materials**

Understanding the local electronic structure of correlated states of matter and its interplay with spin is currently a vast field of research, examples of which are charge/spin density waves or Mott insulator phases. This particularly includes looking at the role of dimensionality in layered, vdW materials [217]. Currently, the local interplay of spin and charge carriers in such materials remains elusive, in both the absence as well as in the presence of a magnetic field. In this realm, local spin sensing based on combinations of STM/nc-AFM will be vital in understanding the role of the spin degree of freedom. Similarly, there are many open questions linked to the role of defects in these materials, and their influence on static spin ordering as well as magnetic excitations, such as magnons. Likewise, probing spin textures in novel lower dimensional superconductors, and imaging new phases of matter, will be of key interest. Furthermore, the role of dimensionality is becoming more and more relevant in 2D and 1D magnets, in understanding the role of various exchange mechanisms, as well as in less frequent types of ordering, such as ferrimagnetism.

### **Complex magnetic states**

It is imperative to understand how complex magnetism emerges based on the interplay between structure and interactions, in order to create new generations of magnetic materials. For example,  $\pi$ -orbital magnetism, as derived from carbon-based materials like graphene, is extremely difficult to detect without local probe measurements [218]. Understanding the potential of magnetic ordering in these materials, and manipulating the spin degree of freedom is possible with the development of STM/nc-AFM based methods. On top of that, STM/nc-AFM based methods can allow for on-surface synthesis of such systems which opens up new possibilities for future design directions. In addition to  $\pi$ -orbital based magnetism, other types of magnetic order are emerging which necessitate a local characterisation of the spin structure. This includes looking at new types of chiral ordering, e.g. skyrmion-based materials, as well as at magnetic states observed in frustrated magnets, such as self-induced spin glasses [204, 205]. Spin based imaging will also be relevant in looking at material systems with complex phase diagrams, for example in the case of 1D ordered systems which exhibit spin-charge separation [219], as well as systems described by multiple order parameters (such as multi-ferroics). Finally, spin sensing based on STM/nc-AFM methods will be vital in understanding spin in topological materials [198], especially in edge states [220].

### **Atomic-scale spin dynamics**

The ability to resolve spin behaviour at the atomic scale and with high temporal resolution opens up the pathway to understand atomically resolved spin relaxation dynamics as well as to quantify coherence times and ultimately coherent control of spins [201, 202, 207, 211]. The state-of-the-art plus future development of spin sensitive time resolved methods based on STM/nc-AFM may address multiple different material systems. With the further development of THz based STM/AFM interrogation [210], it may be possible to resolve magnon dynamics, or coupled magnon–phonon modes, at the atomic scale in a variety of ordered magnetic structures, with spatial resolution. With the advancement of stochastic resonance spectroscopy and electrical pump-probe spectroscopy [207, 211], it is also possible to probe the relaxation dynamics of

coupled spin structures, for example those that exhibit complex relaxation pathways due to various Markov processes [221]. Using similar techniques, it may be possible to look at non-equilibrium dynamics in 1D and 2D spin structures on surfaces, for example those that are chiral, which self-assemble or are artificially constructed. Using a combination of relaxation dynamics and coherent dynamic techniques, based on continuous wave excitation, it may also be possible to understand local ordering in fragile states of matter, like quantum spin liquids as well as pursue coherent manipulation of quantum nanostructures that rely on spin-based read-out [222].

### **Acknowledgements**

This project has received funding from the ERC under the European Union's Horizon 2020 Research and Innovation Programme (Grant Agreement Nos. 818399 and 947717), NWO-VICI programme (VI.C.212.007), and was financed by the Baden Wurttemberg Foundation Program on Quantum Technologies (Project AModiQuS).

## 6. Nanoscale magnetic resonance imaging

Alexander Eichler and Christian L Degen

Laboratory for Solid State Physics, ETH Zürich, 8093 Zürich, Switzerland

### Current state of the technique

Nanoscale magnetic resonance imaging (NanoMRI) was suggested by theoretical physicist John A. Sidles in the early 1990s as a means to resolve biomolecular structures with three-dimensional atomic resolution [223, 224]. At that time, Sidles worked at the Department of Orthopedics at the University of Seattle, and understood how useful it would be to image individual complex molecules, such as single proteins, on the atomic scale. His proposal, known as magnetic resonance force microscopy (MRFM), came shortly after the invention of the AFM and MFM (see section 4). The vision of MRFM relies on mechanical force sensors to detect the weak magnetic signal of (ideally) single nuclear spins. Technically, its function is similar to that of an MFM (see section 4), while using the pulsed RF techniques known from MRI to manipulate the spins and encode information.

Over the course of the past 30 years, MRFM has enabled impressive advances in NanoMRI. Experimental milestones include the detection of a single electron spin in 2004 [225], the 3D imaging of tobacco mosaic viruses at a resolution of 5–10 nm [226], and the demonstration of magnetic resonance diffraction with sub-Angstrom precision [227]. By comparison, conventional MRI used in clinical medicine has a detection limit of order  $10^{18}$  nuclei, limiting the technique to the millimetre length scale.

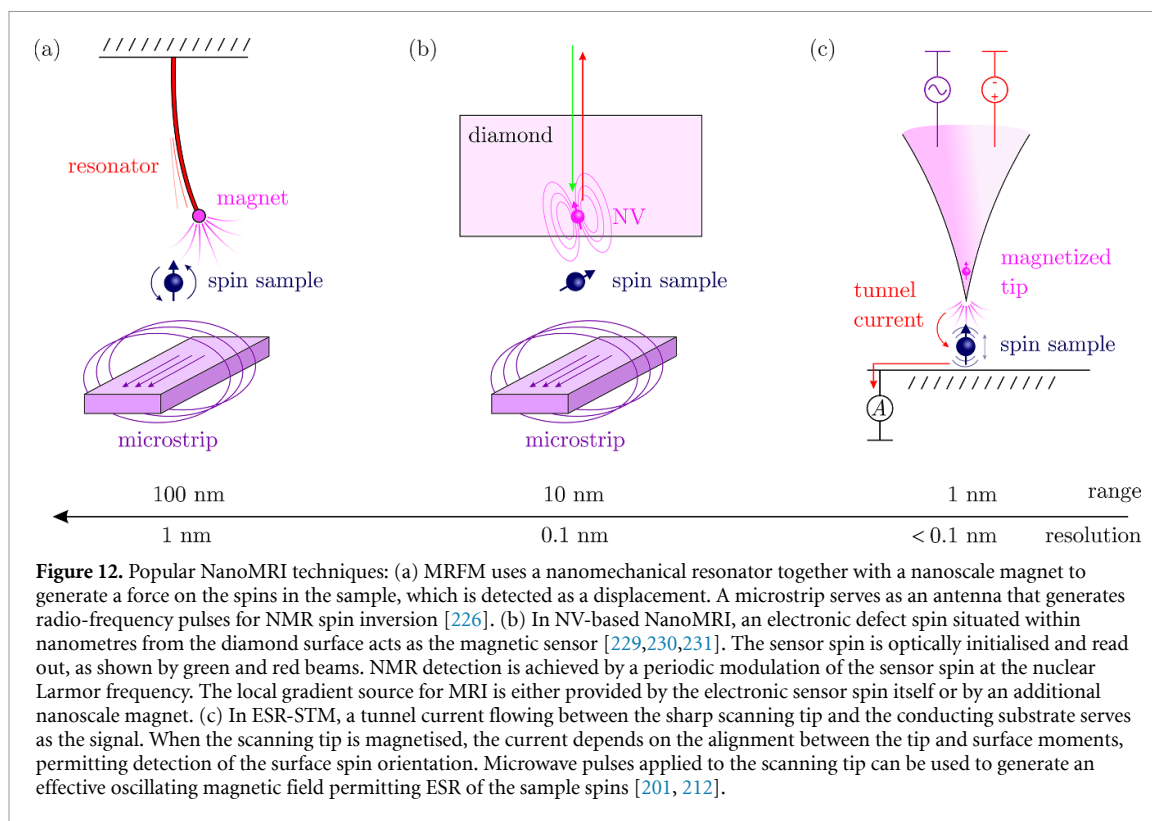
MRFM has not remained the only approach to NanoMRI. Over the past decade, other types of magnetic sensors have emerged that allow for the detection and mapping of magnetic fields at the nanometre scale [228]. Most prominently, these include NMR sensors based on single NV centres in diamond, as well as the technique of ESR-STM, illustrated in figure 12 (see also sections 2 and 5). The various NanoMRI techniques not only differ in their experimental approach and sensitivity, but also in their FOV and scope of use. The FOV is roughly determined by the physical size of the field gradient source. It ranges from  $>100$  nm for a nanomagnet employed in MRFM to  $<10$  nm for the atomic-sized sensors associated with an NV centre or STM tip. It is therefore likely that the different techniques will find application for purposes that require either large imaging depth or optimal spatial resolution [228]. In addition, some techniques are compatible with room-temperature operation while others require cryogenic vacuum environment.

### Future advances of the technique

In order to study the structure of complex molecules, NanoMRI will need to image magnetic signals with sub-nanometre resolution over distances from 10 to 100 nm. Approaching this goal requires advances in many directions including sensitivity, spatial selectivity, long-term stability of the apparatus, and sample preparation. The various NanoMRI subcommunities are currently pursuing different trajectories to reach this goal. To give an impression of the diverse aspects entering this process, we highlight below a few selected efforts:

*Surface functionalisation*—as the technique operates on the nanometre scale, control of surface properties is paramount. For example, the coherence time and charge stability of NV centres is strongly impacted by imperfections on the diamond surface [232]. Likewise, the exact surface termination plays a crucial role in the intrinsic dissipation of nanoscale mechanical oscillators [233] and in non-contact friction effects as the oscillators are approached to the sample surface [234]. Much work is therefore invested into mitigating the harmful effects of surface imperfections on the sensor performance. Obviously, the exquisite atomic surface control in ESR-STM is a great advantage of that technique.

*SNR*—the central bottleneck of NanoMRI is detection sensitivity. This sensitivity can be improved both through optimised sensor geometry and by a reduction of sensor noise. Nanomechanical sensors are typically limited by thermomechanical force noise, which is why the most sensitive experiments are performed at cryogenic temperatures. Furthermore, mechanical sensors are designed with the aim of minimising the damping coefficient  $\gamma = \omega_0 m / Q$  where  $\omega_0$ ,  $m$ , and  $Q$  are the sensor's angular resonance frequency, mass, and quality factor, respectively. There are two main strategies to achieve low  $\gamma$ : on the one hand, long and thin cantilevers or nanowires have low  $m$  and  $\omega_0$  [228]. On the other hand, strings or membranes made from e.g. strained nitride have higher  $m$  and  $\omega_0$  but can possess enormously high  $Q$  [228, 235]. NV-based NanoMRI, by contrast, is limited by optical shot noise as well as the spin coherence time;



here, the avenues towards improvement are a photonic structuring of the sensor, schemes to enhance the optical contrast and readout fidelity, as well as dynamical decoupling detection protocols [45].

*Compressed sensing and image reconstruction*—in most NanoMRI applications, the averaging time required per voxel is significant and leads to timescales of hours, days, or even weeks per image. Many strategies have been devised to reduce this time. A universal approach to speeding up the data taking is through ‘compressed sensing’ [236]. In this method, the measurement apparatus samples only a random subset of the points on the sampling grid in real space. Since NanoMRI data is rarely acquired in real space, such undersampling can result in almost the same image quality as a fully sampled data set after the reconstruction step. Beyond data acquisition, efficient and robust routines for reconstructing the 3D spin density from datasets will be further developed [237, 238].

*Hyperpolarisation*—conventional MRI relies on the bulk magnetisation that is generated in the sample by the large polarising field of a superconducting magnet. Although the thermal polarisation is very small (ca. 1 in  $10^5$  nuclei), this is compensated by the large number of nuclear spins in a voxel ( $>10^{18}$ ). By contrast, in NanoMRI, voxels only contain  $1\text{--}10^6$  nuclear spins, and the thermal polarisation becomes minute. Consequently, many NanoMRI experiments rely on statistical polarisation, which, however, has many drawbacks in terms of signal averaging and spectroscopy. A promising approach to increasing the sensitivity while maintaining the advantage of bulk polarisation are hyperpolarisation techniques, where nuclei are polarised via electronic spins or optically. Nuclear spin alignment can either be achieved by an external polarising agent [239] or, for NV-NMR, by the optically-polarised NV centre itself [238]. Future implementations of NanoMRI are expected to take advantage of nuclear hyperpolarisation.

### Current and future use of the technique for material science

A grand goal of NanoMRI is the imaging of single molecules with true 3D resolution, thus extending clinical MRI to the realms of structural biology and nanoscale materials research. Ideally, this imaging will have near-atomic resolution and the full chemical specificity known to NMR and MRI. One key advantage of NanoMRI over competing methods such as x-ray crystallography or cryo-electron microscopy is the fact that it is non-destructive, allowing to scan a single specimen various times without damaging it and obviating the need for crystallisation or ensemble averaging. The results of such a scan will be complementary to computer-based methods aided by artificial intelligence (AI), which recently have made encouraging progress. As in many other fields, simulations and direct measurements will contribute different information to provide a mature and powerful platform for structure determination.

Beyond structural biology, NanoMRI is expected to feature many interesting niche applications in various areas of materials science, spin physics, and quantum engineering. In materials science, the method can be used to study paramagnetic defects and spin diffusion in surface layers [240]. For example, magnetic defects contribute to  $1/f$  noise in superconducting qubits, limiting their coherence time. Such undesirable surface defects are difficult to investigate with other methods, who lack the sensitivity and elemental resolution to clearly identify them at the nanoscale. Further, NanoMRI will enable the study of spin diffraction in crystalline and amorphous solids at the scale of lattice constants and inter-atomic distances [227]. This capability will be useful for studying spin interactions in periodic materials, investigate spin transport, and identifying mechanisms of decoherence and dissipation. Another interesting application is the imaging of multi-qubit spin registers in solid materials platforms [237]. The capability to directly image and characterise individual spin registers is valuable to build a solid understanding of the interaction in the host material, and to thereby improve the platform quality. Such imaging and control showcase the potential of NanoMRI for designing functional materials for quantum applications.

Finally, a most exciting and fruitful aspect of NanoMRI are the many technical improvements made to sensor platforms, such as pristine mechanical resonators, ultrapure diamond, or spin control and image reconstruction methods. These improvements are by themselves a driving force towards better materials systems and instruments. The exciting technologies that are being envisioned in the quantum sensing and nanotechnology communities depend on such advances to turn our improved understanding of nature at the atomic scale into concrete and useful applications.

## Acknowledgements

The authors thank Fabian Natterer, Konstantin Herb, Nils Prumbaum, and John Abendroth for discussions.



## 7. Magneto-optical imaging

### 7.1. Kerr microscopy

Jeffrey McCord<sup>1,2</sup> and Michael Vogel<sup>1,2</sup>

<sup>1</sup> Department for Materials Science, Kiel University, Kiel, Germany

<sup>2</sup> Kiel Nano, Surface and Interface Science (KiNSIS), Kiel University, Kiel, Germany

#### Current state of the technique

MO microscopy [241, 242] is a polarised light microscopy technique for the characterisation of magnetic materials and the imaging of magnetic domains [243]. It is suitable for general characterising of materials with magnetic phases. Due to its straightforward nature, this traditional technique has become a commonly employed laboratory method for examining magnetic textures down to the nanoscale.

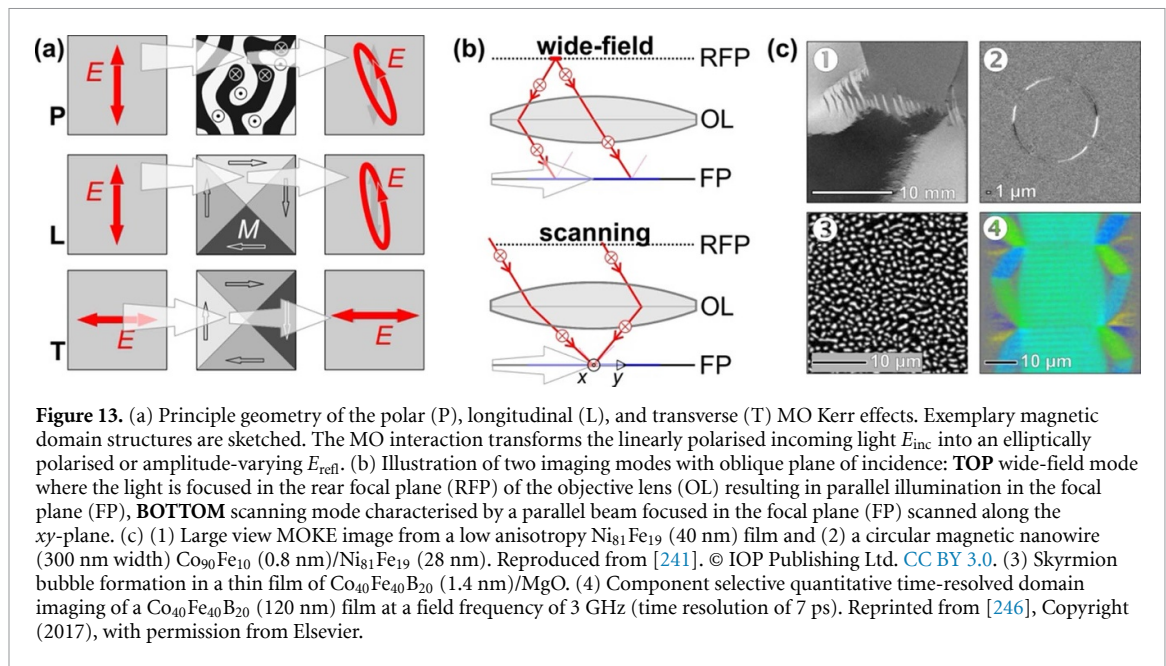
There are several MO effects used to image magnetic domain activity. In transmission and reflection, the MO Faraday and Kerr effect (MOKE) are two linear MO effects. The MO sensitivities are defined by the relative alignment of the magnetisation to the plane of incidence and the orientation of the polarisation of the light, giving rise to three fundamental linear effects: the polar, longitudinal, and transverse MO effects (figure 13(a)). The quadratic MO effect is the Voigt effect. Depending on the geometry, the various effects contribute simultaneously, the separation of which is done by advanced illumination schemes for wide-field microscopy [244, 245, 246]. Confocal scanning MOKE allows for selective sensitivity [247, 248]. By this, magnetic vector information can also be extracted.

MO microscopy uses modified optical polarisation microscopes or similar setups [241]. It allows wide-field imaging or scanning confocal configurations (figure 13(b)) with flexible scales from centimetres to hundreds of nanometres. The Abbe diffraction limit is the fundamental limit to spatial resolution. A basic MO setup uses a separate illumination and observation path, particularly suitable for large-scale sample MO imaging with in-plane magnetisation alignment. Other setups are based on modified polarised light microscopes using a single objective to illuminate and image the magnetisation state. This is suitable for high spatial resolution and high magnification imaging. Where wide-field schemes allow direct imaging of magnetic objects regardless of magnification or image size, scanning methods allow MO signal extraction with a higher MO SNR. Due to the surface sensitivity of the MO Kerr effect, probing magnetic states of sub-nm thick films is standard. By relying on optics, flexible temporal resolution down to the sub-ps regime can be achieved with appropriate pulsed light sources, typically through stroboscopic imaging schemes. Example images are shown in figure 13(c).

#### Future advances of the technique

Future developments to improve SNR and spatial resolution include several needed adaptations of the MO imaging technique. It is expected that novel imaging modes will be required to advance. New functionalities are projected to become available.

- *MO contrast recovery*: while the SNR in scanning Kerr microscopy can be enhanced by lock-in amplification, the SNR for wide-field imaging is significantly lower. With the advantage of wide-field imaging in a single exposure, related image detection schemes with improved SNR are an obvious extension to wide-field Kerr microscopy. This will allow changing from the most commonly used differential imaging scheme with a magnetically homogeneous magnetic background image. Such imaging schemes require significant but worthwhile changes in microscope illumination and image detection. In-line polarisation variations based on real-time GPU image processing have already been demonstrated in this direction [249]. The fundamental goal is to obtain a low-noise wide-field magnetic micrograph without the need for additional images of a different magnetic state in a most efficient way. The addition of plasmon filter structures for MO contrast enhancement is another possible development avenue. AI and machine learning based methods for image noise reduction and local contrast enhancement can be implemented.
- *MO effects*: more intelligent illumination schemes will enable direct component selective MO imaging. Multiple simultaneously applied illumination sources with different polarisation states should allow switchable and single exposure imaging of e.g. pure and continuously configurable in-plane MO contrast independent of the actual planes of incidence. Furthermore, the possible use of full illumination aperture is expected to reduce unwanted effects from topographic structures. Such an implementation will also lead to a significantly better spatial resolution.
- *Improvement in spatial resolution*: dielectric microelements, such as microspheres, can focus light in their vicinity. Due to a near-field interaction a so-called photonic nanojet (PNJ) forms. This produces a virtual image with enhanced resolution and magnification [250]. Incorporating them into MO imaging systems can



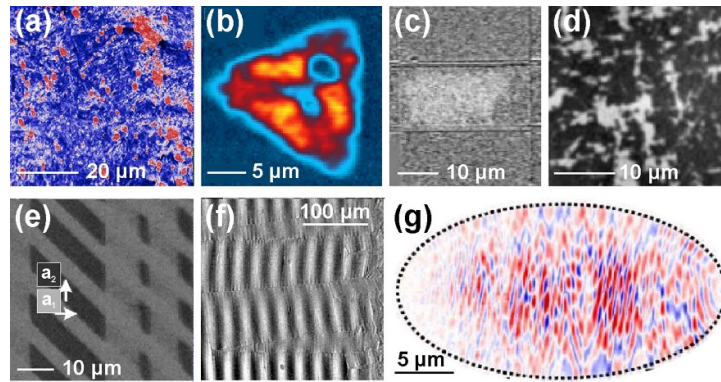
improve spatial resolution significantly. A challenge for MO microscopy is to maintain the polarisation of the light. Alternatively, moving to shorter imaging wavelengths is another way to improve spatial resolution in MO microscopy. This requires a change to quartz optics and is more easily implemented in scanning Kerr microscopy setups in conjunction with CW or pulsed UV or soft-x-ray laser sources. Yet, imaging by PNJs appears to be the most promising technique for pushing the spatial resolution limit of wide-field MO imaging well below 100 nm. Direct magnetic scanning near-field optical microscopy may play a role in the future, with the first practical realisation for time-resolved MO imaging [251].

- **Time-resolved MO microscopy:** the ability to image fast magnetisation dynamics using stroboscopic imaging techniques is standard in scanning and adapted for wide-field MO microscopy. Yet, extending fast wide-field imaging modes to single-shot imaging of non-repetitive events has yet to be explored for timescales beyond the microsecond regime. Providing single-shot imaging for the full range of magnetisation dynamics down to at least the nanosecond regime will open a whole new field for imaging the magnetisation behaviour of magnetic materials and devices.
- **Conoscopic MO imaging:** observation of the objective's back focal plane, or Fourier imaging, allows new domain imaging modes. In conoscopic mode, a range of magnetic texture information (e.g. lateral dimensions of magnetic domains and orientation) can be extracted from Fourier space images. Temporal resolution for imaging magnetisation dynamics using stroboscopic imaging techniques is standard in MO microscopy. Yet, fast single-shot imaging is limited by the brilliance of the illumination source used within lab-based microscopy setups. For time-resolved imaging at timescales below the SNR limit, averaging in Fourier space would allow for extracting otherwise hidden domain structure parameters.
- **Integration with other magnetic microscopy techniques:** the possible high temporal resolution and the ability to extract magnetometric data with MO imaging offer substantial benefits in combination with different techniques, especially scanning microscopy techniques. Therefore, MO microscopy is expected to be further integrated with complementary techniques of higher spatial resolution for the simultaneous extraction of magnetic domain information by different techniques. Adapted optical setups will be developed for this purpose.

### Current and future use of the technique for material science

MO microscopy is becoming increasingly important in materials science. MO microscopy provides a straightforward method of identifying magnetic phases in ferrous alloys, complementing the research and development of advanced materials [252]. (See figure 14 for examples from this section). The adaptation of MO microscopy to materials engineering is just beginning. With the growing availability of commercial MOKE microscopes, MO imaging is expected to be increasingly applied in the future. Research on magnetic materials for existing and future energy applications with the possibility of time-resolved operando device characterisation is only possible with MO microscopy, often assisted by MO indicator films [241].

Furthermore, the MO surface sensitivity enables studies of 2D crystallographic materials, with numerous new material systems being introduced. 2D magnets within the family of vdW bonded materials have been



**Figure 14.** Example MO images from various materials and devices. (a) MO microscopy image of a deep cryogenic heat-treated steel surface in an applied magnetic field, where the (colour-coded) blue regions correspond to martensite. Reprinted from [252], Copyright (2021), with permission from Elsevier. (b) Scanning MOKE microscopy image of a  $\text{WS}_2$  island, displaying a spatial map of spin density. A central core of low spin/valley density surrounded by regions of higher spin/valley density become visible. Reproduced from [253]. © IOP Publishing Ltd. CC BY 3.0. (c) Spin-orbit-torque driven magnetisation switching measurements in a  $\text{WTe}_2$  (80 nm)/NiFe heterostructure. Reproduced from [254], with permission from Springer Nature. (d) Magnetic domains in an antiferromagnetic NiO (20 nm) film obtained from an MO microscope. Reprinted figure with permission from [257], Copyright (2019) by the American Physical Society. (e) Remanent ferromagnetic domain distribution in a ferroelectric/ferromagnetic  $\text{BaTiO}_3/\text{CoFe}$  structure. [258] John Wiley & Sons. [Copyright © 2011 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim]. (f) Time-resolved wide-field image of a moving SAW, interacting with the magnetic domain walls in a quartz/ $\text{FeCoSiB}$  structure. [259] John Wiley & Sons. © 2022 The Authors. Advanced Electronic Materials published by Wiley-VCH GmbH. (g) Spatially-resolved time-resolved measurements of non-linear spin waves in a  $\text{Ni}_{80}\text{Fe}_{20}$  thin film element. Reproduced from [260]. CC BY 4.0.

obtained experimentally, with properties markedly different from their bulk phase. The need of cryogenic temperatures and high magnetic fields to control the magnetism in such materials makes advanced high-resolution MO microscopy an ideal technique for investigation. Time-resolved scanning Kerr was already used to investigate spin/valley lifetimes in monolayer  $\text{WS}_2$  [253]. MO microscopy will further contribute to the investigations of magnetic topological semimetals like magnetic Weyl semimetals, which are of interest for novel spintronic devices. MO microscopy has so far proven to be effective for the investigation of spatial effects of Weyl semimetal heterostructures. Non-trivial ways of purely spin-orbit torque-driven magnetisation switching, aided by interfacial chirality effects, can be easily studied and quantified by low-temperature MO microscopy [254]. The integration of high temporal resolution MO microscopy to study the dynamics and lifetimes of processes in semimetals prepares MO microscopy for future challenges in this area. It is anticipated that with improved MO sensitivity and spatial resolution, MO microscopy will continue to adapt to the various challenges of studying new classes of materials. With the move towards small device integration, advanced MO microscopy will become even more indispensable.

Magnetic materials provide a versatile platform for the formation of real-space non-trivial topological winding spin solitons like skyrmions. Improved MO microscopy will continue to be used to investigate skyrmionic bubbles [255]. The previously discussed advances in spatial resolution will be vital in studying the magnetisation dynamics of single skyrmions and skyrmion strings. Together with the rise of sufficiently low damping materials, this will lead to insights into the resonance modes of skyrmions. The observation of which will be within the reach of future MO microscopy. Also, AFM materials are candidates for future spintronic devices because of their scaling advantages due to minimised stray magnetic fields and ultrafast switching speeds. However, only non-collinear AFMs exhibit MOKE contrast [256]. The MO Voigt contrast and the linear dichroic material contrast allow spatial detection of the AF spin distribution [257]. As the symmetry of both effects is equivalent, a clear distinction between the effects is not straightforward. Current and future AFM materials studies will benefit from the proposed improvement in contrast sensitivity and advanced time-resolved imaging schemes. Second harmonic MO signal generation (section 7.2) offers a direct way of imaging fully compensated AFMs.

The spatiotemporal characterisation of functional magnetic material devices will remain at the heart of MO microscopy. In this respect, multiferroic or ME materials will be studied, including piezoelectric/ferromagnetic composites with large ME coupling coefficients. MO microscopy can provide unprecedented insight into the local interactions in the multiferroic materials [258]. For ME sensing devices, the magnetic domain behaviour is of relevance due to domain wall mediated noise. As has been demonstrated for surface acoustic wave devices, high-resolution time-resolved MO microscopy can also provide understanding of electrical response characteristics [259]. Because MO microscopy allows

quasi-static and time-resolved imaging of functional devices, even when covered in transparent layers. It is often the only choice for obtaining magnetic micrographs of incorporated materials.

In this context, time-resolved MO microscopy provides unique insight into the dynamic properties of a wide range of magnetic materials in bench-top setups, such as the study of collective magnetic excitations, e.g. various spin-wave-related phenomena for envisaged future devices. Of particular interest for future information processing are non-linear processes, in which degenerated phase-states were identified by advanced time-resolved microscopy schemes [260]. This demonstrates the adaptivity of MO microscopy to new scientific challenges. Another glimpse on future advances is investigating the ultra-fast behaviour of magnetic domains [261] by MO diffraction pattern analysis.

MO microscopy, especially with improved MO sensitivity and spatial resolution, will continue to adapt to the various challenges of studying new classes of materials and devices. Although it is a traditional technique, the extreme flexibility of MO microscopy means that it will continue to evolve.

## Acknowledgements

J M acknowledges funding through the German Research Foundation (DFG) through the DFG Heisenberg programme (Mc9/9-1, Mc9/9-2), and the Collaborative Research Centre CRC 1261 'Magnetoelectric Sensors: From Composite Materials to Biomagnetic Diagnostics'. M V acknowledges funding by the US-DOE, Office of Science, MSED and by the US-DOE, Office of Science, BES, under Contract No. DE-AC02-06CH11357.

## 7.2. Second-harmonic-generation microscopy

Manfred Fiebig

Department of Materials, ETH Zurich, Zurich, Switzerland

### Current state of the technique

Optical SHG denominates the frequency doubling of light in a material. In the electric-dipole description of the light field, it is expressed by  $P_i(2\omega) = \epsilon_0 \chi_{ijk}^{(2)} E_j(\omega) E_k(\omega)$ , where  $\mathbf{E}(\omega)$  and  $\mathbf{P}(2\omega)$  describe the incident light at frequency  $\omega$  and the frequency-doubled polarisation induced in the material, respectively [262]. As a multi-photon process, SHG requires intense light fields to be observed, so its first detection occurred in 1961 [263], right after invention of the laser. SHG is a highly symmetry-sensitive technique [264] and as such perfectly suited to study systems exhibiting ferroic order [265] because these break point-group symmetry by definition [266]. For example, electric-dipole SHG is only permitted in non-centrosymmetric systems and is therefore often used for the investigation of ferroelectric states, which inherently break inversion symmetry. The first SHG contributions from magnetic order were reported in 1991 [267], following predictions and precursor experiments from the 1960s [265]. Since SHG, in contrast to linear-optical techniques, does not couple to the magnetisation but to the symmetry of a magnetically ordered material, SHG is particularly efficient in the study of antiferromagnetic materials, where the magnetisation is zero [268].

The investigation of magnetically ordered materials with SHG is particularly advantageous when the specific degrees of freedom of optical techniques come into play [265]. (i) Polarisation-dependent SHG spectroscopy can distinguish between differently ordered sublattices of a material. In particular, the coexistence between magnetic and electric orders in a so-called multiferroic is accessible with the same SHG experiment and in a single experimental run, which gives direct access to the magnetoelectric (ME) coupling effects between the coexisting ordered states. (ii) Spatial resolution up to the limit of the wavelength of about 1  $\mu\text{m}$  provides access to domain patterns. This permits to image the distribution of antiferromagnetic spin-reversal domains within seconds, which is unattainable with other techniques. Domain structures are mostly imaged in the near-surface bulk region, but interference techniques [269] and Cerenkov SHG [270] are also capable of 3D imaging. Despite the wavelength-limited resolution, the ordered state of nanoparticles can be addressed. On the other hand, spatial resolution may be increased to, in principle, <100 nm by SHG near-field imaging [271]. (iii) Finally, with a temporal resolution up to the laser pulse length of a few fs, magnetisation-dynamical processes beyond the equilibrium of electron, spin, and lattice temperatures are resolved, which offers prospects in the direction of ultra-fast magnetic reading or writing.

On the theoretical side, mechanisms coupling the magnetic order to the nonlinear photon field were identified [265], and a general tool (#SHAARP) calculating the polarisation selection rules for SHG experiments on ordered materials was developed [272, 273].

### Future advances of the technique

The most important future development in SHG microscopy is probably its evolution from a basic-research tool towards practical or even commercial applications. For example, in conceptual experiments SHG was used as an *in-situ* characterisation tool that tracks the emergence of functional properties like a spontaneous polarisation or magnetisation during the growth of epitaxial heterostructures [274]. It is thus possible to observe the threshold thickness at which the ordered state emerges, the interaction between the sequentially deposited constituents, and the formation of interim phases that, albeit disappearing again with the ongoing growth, determine the structure and properties of the finalised heterostructure. Thus far, such *in-situ* SHG (ISHG) probing was mostly applied to ferroelectric materials, but the application to magnetically ordered systems is imminent. Since it is possible to track the evolution of the magnetic order at any time during the growth process, the formation of undesired magnetic states can be avoided and targeted magnetic functionalities can be enhanced by feedback-controlled adaption of the growth conditions. In heterostructures combining magnetically and electrically active materials, their cross-correlation can be optimised in the same way and help to obtain novel multiferroics or enhanced ME coupling effects.

It is foreseeable that the conceptual demonstration of ISHG will soon lead to the development of a prototype that can be flanged onto the growth chamber for sputtering deposition, just like the electron gun for *in-situ* electron diffraction. In combination with a software like #SHAARP for signal analysis [273], this will open up ISHG probing to non-specialists in lasers and nonlinear optics. A commercial setup for ISHG maybe expected two to five years from now.

In parallel, the concept of probing a magnetic material during its composition (*in situ*) will be expanded towards probing a magnetic device during its operation (*operando*) by SHG. In magnetic materials, functionality is usually synonymous to modification, mostly switching, of a magnetic state, and because of



the non-invasive, remote nature of SHG probing, this can be tracked with spatial, temporal, and sub-lattice resolution without destroying the working component.

In moving towards the technological-commercial implementation of SHG microscopy, the methods of nonlinear-optical probing of magnetic order will be developed further. The present focus on the leading-order electric-dipole contribution to SHG will be expanded towards processes including magnetic-dipole or electric-quadrupole contributions to the light field. This permits more selective addressing of the electronic transitions of a material and in particular the investigation of systems with centrosymmetric magnetic order [265, 275, 276]. The associated processes are significantly weaker than the leading electric-dipole contribution, so ways to counteract this inherent attenuation will be explored. This may involve resonant tuning of one or more of the light fields involved in the nonlinear excitation process to the electronic transitions of the probed material. Along with this, nonlinear processes other than SHG, like sum-frequency generation or third harmonic generation, may be involved [265]. Finally, the SHG process may be transferred from (near-) visible frequencies towards other photon-energy ranges. THz SHG appears to be specifically promising as it allows to probe the elementary excitations in materials resonantly and thus very directly and effectively. Hence, significant extensions of SHG and other nonlinear optical methods are still foreseeable in the field of spectroscopy. In contrast, the limits in terms of temporal and spatial resolution may have been reached already. A laser pulse length of 100 fs seems to remain the optimal compromise between temporal and spectral resolution. While in terms of thickness, the complex ordered states of a single molecular layer of a material can already be detected [277], the sub-wavelength spatial resolution of optical-near-field experiments is diminished by the extreme complexity of near-field SHG techniques. Here, a compromise may be an improvement of the statistical tools for analysing the far-field SHG response from structures ordering on sub-resolution length scales [265].

SHG microscopy may be expanded by combining the previously mentioned degrees of freedom of optical techniques (i) to (iii) in a synergetic way. For example, ultrafast magnetisation-dynamical experiments with SHG probing may be performed with optical resolution with the help of specially enhanced digital cameras ('EM-CCD') or improved statistical image analysis. This will permit to observe the time-dependent as well as the spatial dissipation of a non-equilibrium state of the spin system at the same time. In pump-probe experiments, intense excitation in the x-ray, (near-) visible, or THz range will be synchronised to SHG probing. The aforementioned combinations of techniques will especially give a boost to the study of spin dynamics in systems with compensated magnetisation, where SHG probing is the method of choice.

### **Current and future use of SHG microscopy for materials science**

One of the greatest strengths of SHG microscopy is spatially resolved imaging of spin-related structures with 'hidden' order. These are foremost forms of magnetic order exhibiting zero net magnetisation like antiferromagnetism. Because of a lack of magnetisation to couple to, established linear-optical techniques, mostly the MOKE, are not applicable, and symmetry-sensitive SHG microscopy comes into play for imaging AFM domains [268]. With the growing interest in antiferromagnetism for spintronics applications [278] or as neighbouring phase to other correlated states, convenient experimental access to AFM domains will become ever more important because essentially the functionality of any ferroic material roots in the distribution and manipulability of its domains.

Currently, there is also a growing interest in novel forms of magnetically compensated ferroic order. Since their physical origin is still under investigation, domain imaging by SHG microscopy is all the more important since it addresses the ordered state via macroscopic symmetry properties with no need to know the microscopy. One of the new types of compensated magnetic order is ferrotoroidicity, the spontaneous alignment of magnetic whirls, which is interesting because of its inherent ME effect. Other materials may exhibit spin-driven ferroaxial order, where magnetic screws break mirror symmetries, or magnetic-multipolar types of order. Broken spatial and/or temporal inversion symmetries in all these systems can act as a source for new ME optical or transport properties [265].

The study of ME correlations in multiferroics will remain a central theme of SHG microscopy. ME correlations are ultimately based on the coupling between magnetic and electric domains, and the only technique that can image both within the same experiment is SHG. There are also a number of spin states, however, whose correlations are not captured by the symmetry formalism used for ferroic systems and to which SHG microscopy may nevertheless couple. This includes dynamical order in the form of loop currents [279] or supercurrents [280] as they occur in superconducting materials. The conducting surface state of topological insulators may be accessed by SHG because of its capability to distinguish between surface and bulk contributions. Other systems with topological properties are skyrmions and Weyl semimetals, which are interesting for new electronics and computing technologies. First very preliminary characterisation experiments by SHG are presently reported for all these systems [275], so a rapid expansion of SHG from ferroic towards topological and spin-correlated systems at large is likely.

There is a good chance that oxide electronics will be revolutionised by the development and commercialisation of ISHG in the growth of epitaxial heterostructures with ME functionalities. The current, mostly methodical experiments show that ISHG measurement can be used to control the growth process in many ways [265]. In feedback loops during the deposition, multi-domain formation can be promoted or suppressed. The deposition may be stopped at the point of emergence of magnetic or electric order to obtain particularly large susceptibilities. The influence of and coupling between the interfaces separating neighbouring constituents can be identified and used as an additional parameter to design functionality. Interim phases occurring during growth are detected, which helps to understand the relation between the deposition process and the final product. In addition, the impinging laser light itself may be used to modify the formation processes locally through the thermal and electronic excitations it introduces.

The connection of spatial and ultrafast temporal resolution in pump-probe imaging experiments will help to connect the relaxation processes between electrons, spins, and the lattice to spatial dissipation processes and gradient effects. In systems with multiple competing interactions, this can lead to better understanding and control of the subtle balance between states of similar energy. This may lead to an improved concept for ultrafast magnetisation processes and, ultimately, towards all-optical magnetic phase control.

An inherent disadvantage of optical experiments around the visible range is that the associated photon energies of  $\sim 1$  eV may be far away from the characteristic energy of the states that are actually probed. Magnetic excitations, for example, are quantified in magnons with energies of around 10–100 meV. Here the extension of SHG experiments into the THz range with either THz SHG or a combination of THz stimulation and visible-SHG probing may offer significantly more direct access to magnetic correlation dynamics. This may be applied in studies of the pseudogap of (especially unconventional) superconductors [279], in investigations of the competition between magnetic order and disorder near quantum-critical points, or in the generation of radiation-driven multiferroic-like states with dynamical ME correlations [281].

## 8. X-ray techniques

### 8.1. Magnetic x-ray microscopies in real space

Peter Fischer<sup>1,2</sup>, Aurelio Hierro-Rodriguez<sup>3,4</sup>, Simone Finizio<sup>5</sup> and Sarnjeet Dhesi<sup>6</sup>

<sup>1</sup> Materials Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States of America

<sup>2</sup> Department of Physics, University of California Santa Cruz, Santa Cruz, CA 95064, United States of America

<sup>3</sup> Departamento de Física, Universidad de Oviedo, 33007 Oviedo, Spain

<sup>4</sup> CINN (CSIC-Universidad de Oviedo), 33940 El Entrego, Spain

<sup>5</sup> Swiss Light Source, Paul Scherrer Institute, 5232 Villigen PSI, Switzerland

<sup>6</sup> Diamond Light Source, Chilton, Didcot, Oxfordshire OX11 0DE, United Kingdom

#### Current state of the technique

Magnetic x-ray microscopy in real space can be done with XPEEM [282] and with x-ray Fresnel zone plate based systems, realised in magnetic STXM and full field magnetic transmission x-ray microscopy (MTXM) [283–285] (figure 15).

There is complementarity as XPEEM provides surface sensitivity, whereas STXM and MTXM allow for bulk studies up to the absorption lengths of x-rays ( $\sim 150$  nm).

Prerequisites for both are polarised x-rays to use x-ray magnetic circular or linear dichroism (XMCD/XMLD) as element-specific and quantifiable magnetic contrast for FM/FiM/AF/AM materials, high resolution and efficiency x-ray optics for STXM/MTXM, various x-ray detectors, i.e. 2D x-ray detectors for MTXM and XPEEM and x-ray point detectors, e.g. avalanche photo diode (APD) or fluorescence detectors for STXM, or sample current detection in STXM. Various detection schemes are adopted, i.e. x-ray reflection, x-ray fluorescence, secondary electron (SE) detection. Contrast enhancement and background reduction is achieved by modulating the x-ray polarisation (left/right for XMCD and various angles for XMLD), comparing contributions from various spin-orbit coupled x-ray absorption edges, e.g.  $L_3/L_2$ ,  $M_5/M_4$ , or normalisation to magnetically saturated (flat) images.

A variety of substrates can be used, e.g. x-ray transparent membranes (mostly  $Si_3N_4$ ) for absorption detection in MTXM/STXM, solids and conductors for XPEEM or STXM in sample current detection mode, luminescence or phosphorescence for specific STXM applications.

Stroboscopic pump-probe schemes using the inherent time structure of synchrotron-radiation (SR) sources allow for time-resolved magnetic microscopy of fully repeatable magnetisation dynamics [286–288].

3D imaging of FM materials is performed through tomographic [289] (or laminographic [290]—for extended surface samples) imaging, by the acquisition of several projections at different orientation angles (ideally around more than one rotation axis) of the sample with respect to the probing x-ray beam considering that the XMCD contrast measures the projection of the local magnetisation onto the photon propagation direction. To acquire the single projections transmission techniques such as STXM and MTXM are used but also XPEEM in transmission (shadow) mode [290].

External parameters can be applied during the image recording: temperature, external magnetic fields (limited with XPEEM), excitations ( $B$ -field, current,  $E$ -field).

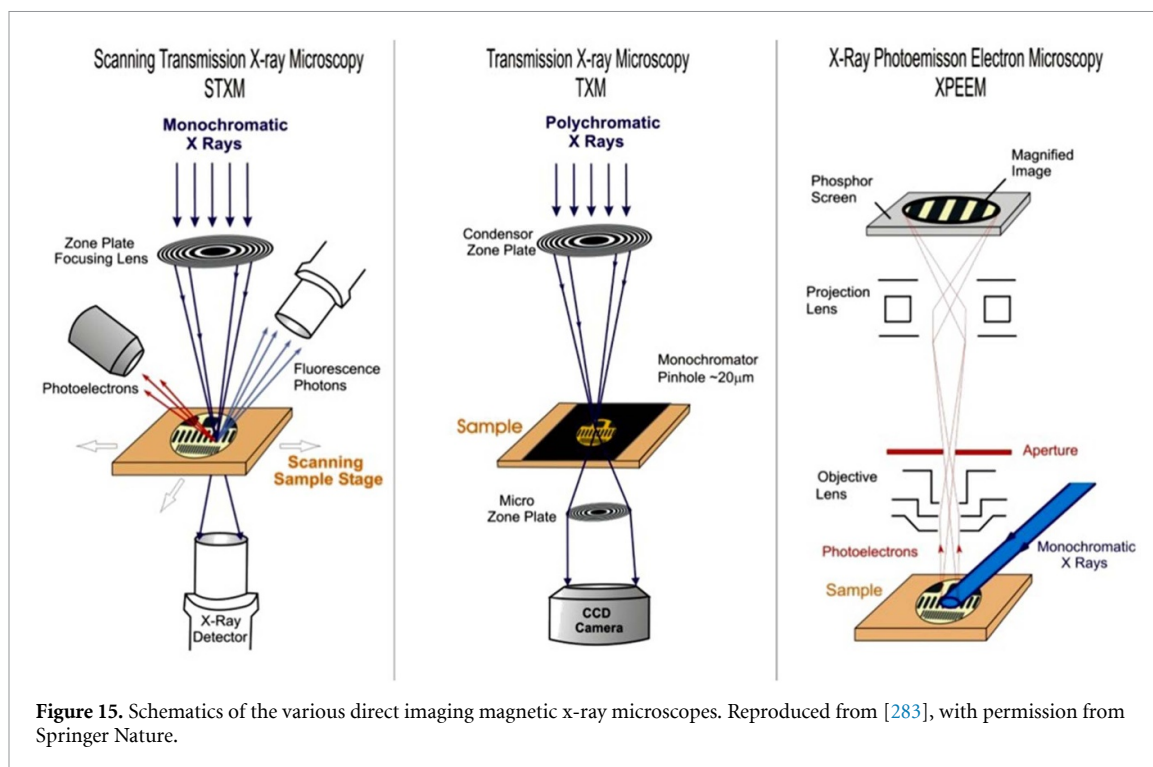
In general, radiation damage is not an issue for magnetic systems of interest in real space x-ray microscopy, such as metallic films, multilayers or nanopatterned magnetic structures.

#### Future advances of the technique

Several large-scale facilities worldwide are currently in the process of upgrading their synchrotron sources mostly to enhance the coherent flux, which opens new possibilities for magnetic x-ray microscopies. A few examples, which are not meant to be conclusive, showcase current trends.

- Future STXM at SLS

The Swiss Light Source (SLS) will undergo an upgrade to a diffraction limited source during 2023–2026. SLS 2.0 will increase the photon flux at both the soft and hard x-ray energy ranges by several orders of magnitude, allowing for a reduction of measurement time, and to image samples with weaker XAS/XMCD signals. However, the increased photon flux yields an increased probability of multi-photon events at the point detector used in STXM imaging, leading to erroneous measurements [291]. In order to overcome this



**Figure 15.** Schematics of the various direct imaging magnetic x-ray microscopes. Reproduced from [283], with permission from Springer Nature.

issue, the APD installed at the PolLux STXM will be replaced with a multi-sector APD, each sector with its separate amplifier and discriminator, thus reducing the probability that two photons interact with the same sector, producing an erroneous count, if compared with a single-sector APD.

- Future PEEM at DIAMOND

The Diamond-II upgrade will provide improved capabilities for magnetic imaging using coherent probes, enable an increase in the number of beamlines using a new lattice as well as increasing the photon flux at higher photon energies by increasing the energy from 3 GeV to 3.5 GeV. A new flagship coherent soft x-ray imaging and diffraction beamline will focus activity on coherent 3D tomography from ferroic materials and host a comprehensive portfolio of sample environments for imaging under magnetic and electric fields and at cryogenic temperatures. The beamline is designed to take full advantage of the increase in coherent flux and aims to provide state-of-the-art 3D magnetic vector tomographic ptychographic reconstructions in a fraction of the time that is currently possible. The upgrade will also replace the existing insertion devices on the XPEEM facility with fast polarisation switching devices enabling considerably accelerated magnetic imaging of thin films and interfaces.

- Future developments with MTXM

There are no current plans worldwide for significant upgrades in MTXM. The MISTRAL beamline at ALBA is currently the only MTXM with a dedicated setup for performing magnetic tomography in novel 3D nanomagnetic structures, which is one of the frontiers in current nanomagnetism research [292, 293]. Developing novel optical designs based on reflective rather than diffractive optics seems to be an interesting alternative as it would increase the sample-optics distance and would allow to image buried interfaces. Operating at low  $T$  and at high magnetic fields would be essential for *in-situ* and *in-operando* studies of quantum effects in magnetic materials. Utilising the increased coherence at future SR facilities, which so far has been considered counter-productive to MTXM might require significant dedicated efforts to advance conventional diffractive x-ray optical elements [294].

Other noteworthy developments with complementary advanced magnetic spectromicroscopy techniques are:

- SPLEEM

Spin-polarised low-energy electron microscopy (SPLEEM) is used for the study of surfaces and interfaces with the ability to detect the spin-polarisation of the electron beam. Recent developments include sample stages to vary the sample's temperature down  $\sim 4$  K. The strength of SPLEEM are investigations during *in-situ* deposition of metallic and oxide films or multilayers. Important contributions by SPLEEM include investigations of chiral spin textures [295].

- Momentum microscopes

Momentum microscopy (MM) [296] allows to investigate the electronic structure of solids by utilising a combined PEEM/ARPES spectrometer to couple lateral information of about 50–100 nm with the spectroscopic analysis. Sample information is obtained either in real space mode (PEEM) or in reciprocal space mode (ARPES and MM) resulting in a constant energy map showing a real space image or a constant energy surface of the samples band structure (MM), or an energy dispersive image of a real or reciprocal space coordinate (ARPES). Coupling ARPES instruments with spin sensitive devices, e.g. Mott detectors allows to extract also the three components of the spin polarisation vector with high energy and momentum resolutions.

### Current and future use of the technique for material science

The specific features of magnetic x-ray microscopies will continue to provide unique information to understand the underlying microscopic properties and behaviour in magnetic materials. The elemental specificity represented in the spectroscopic fingerprint is probably the single most important feature and will be key to disentangle spatial and temporal distributions of spin and orbital magnetic moments and their dynamics in heterogeneous magnetic materials.

The sensitivity to the projection of magnetisation onto the photon helicity in ferro(i)-magnetic systems and the angular dependence of linear dichroism in antiferromagnetic systems will allow to address the full 3D character of spin textures including topologically modified systems, such as curved surfaces or core–shell nanoparticles. Revealing the depth dependence of spin textures, e.g. via reflection mode imaging will be critical to understand the role of interfaces, and modifications of magnetic interfaces in the presence of spin–orbit or asymmetric exchange interactions [297].

The increase of coherence at next generation light source, which will tremendously benefit x-ray scattering based techniques, is probably less relevant to direct real space imaging magnetic x-ray microscopies, although using phase sensitive x-ray optics might be an interesting avenue to pursue, harnessing a combination of amplitude and phase magnetic contrast methods.

The recently discovered altermagnetic systems [298] will most likely exhibit interesting topological spin textures, albeit initial magnetic imaging experiment will most likely address the anisotropic spin textures in *k*-space, however exploring a combination of XMLD and XMCD might be worth exploring.

The relatively large cross sections of dichroism effects and the quantifiability (MO sum-rules [299, 300]) will provide microscopic information and complement macroscopic high sensitive magnetisation information, e.g. through SQUID magnetometry.

The inherent temporal structure of x-ray sources, which can range from the 10–100 ps range at SR storage rings down to the fs regime at XFELs can provide unprecedented insight into the dynamics of spin textures. In addition to fully repeatable dynamics, including standing and propagating spin waves [301], which can be also be readily simulated in micromagnetic simulations, it will be the uncorrelated spin fluctuations on the nanoscale [302] which can be anticipated to become a focus of future research. However, real space x-ray microscopies will be most likely not feasible to experimentally address those topics due to the high intensity of the free-electron laser (FEL) pulses, and their strong variability due to the SASE emission. Among the challenges would be enormous space charges in XPEEM, radiation damage to the illumination optics in MTXM and STXM and enormous photon exposures to the sample with STXM.

Topics of increased interest will be *in-operando* studies of future real magnetic devices, including investigations of spin currents and spin accumulation in nanoscale magnetic materials [303].

Although the spatial resolution of x-ray microscope will most likely not surpass other techniques, such as SP-STM or EM-based microscopies, correlating magnetic x-ray microscopies with those techniques will enable a coherent understanding of magnetic phenomena across length and time scales and in multiple dimensions.

Combining tomography with time resolution seems to be a viable path to explore complex systems [304], notably artificial spin ice systems and other artificially created 3D magnetic nanostructures.

The implementation of variable temperature sample environment and much larger magnetic fields will be of high interest to investigate the microscopic origin of magnetic phase transitions.

Magnetic x-ray microscopies, notably when operating in x-ray transmission, which is the simplest detection mode with the least perturbation of the recorded signal requires x-ray transparent substrates similar to TEMs. Until recently, those membranes used mostly non-epitaxial materials. However, recent developments in focused-ion-beam lithography enabled the fabrication of thin (<100 nm) membranes from single-crystalline substrates such as SrTiO<sub>3</sub> and Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub>, opening the possibility to study epitaxial materials, e.g. multiferroic perovskites, magnetic garnets) with transmission x-ray microscopy (TXM) techniques [305].



## Acknowledgements

P F acknowledges support by the US DOE, Office of Science, Office of Basic Energy Sciences, Materials Sciences and Engineering Division under Contract No. DE-AC02-05-CH11231 (NEMM program MSMAG). A H-R acknowledges support by Spanish MICIN under Grant PID2019-104604RB/AEI/10.13039/501100011033 and by Asturias FICYT under Grant AYUD/2021/51s185 with the support of FEDER funds.

## 8.2. Coherent x-ray imaging of magnetic systems

Claire Donnelly<sup>1,2</sup>, F Büttner<sup>3,4</sup>, Ofer Kfir<sup>5</sup>, Wen Hu<sup>6</sup> and Sergey Zayko<sup>7</sup>

<sup>1</sup> Max Planck Institute for Chemical Physics of Solids, Dresden, Germany

<sup>2</sup> International Institute for Sustainability with Knotted Chiral Meta Matter (WPI-SKCM2), Hiroshima University, Hiroshima 739-8526, Japan

<sup>3</sup> Helmholtz-Zentrum für Materialien und Energie, Berlin, Germany

<sup>4</sup> Experimental Physics V, Center for Electronic Correlations and Magnetism, University of Augsburg, Augsburg, Germany

<sup>5</sup> School of Electrical Engineering, Iby and Aladar Fleischman Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel

<sup>6</sup> National Synchrotron Light Source II, Brookhaven National Laboratory, Upton, NY 11973, United States of America

<sup>7</sup> Max Planck Institute for Multidisciplinary Sciences, Göttingen, Germany

### Current state of the technique

Coherent x-ray magnetic imaging is a lensless approach to magnetic microscopy. Highlights of this approach are the simultaneous access to the element-specific phase and absorption contrasts and the promise of wavelength-limited spatial resolution in the x-ray regime. Its development sparked a new era of low-dose, high-contrast imaging with unique capabilities ranging from element specific phase contrast imaging [312–315] to spatial resolutions of 5 nm in 2D [316] and 10 nm in 3D [317].

The procedure for coherent magnetic imaging is shown in figure 16. First, an (ideally brilliant) coherent source of polarised x-rays (synchrotron, FEL, or high-order harmonic generation (HHG)) is used to illuminate the magnetic sample. This illumination is micro-focused rather than nano-focused, which simplifies the need for advanced x-ray optics significantly. Spatial information is encoded in the coherent scattering distribution (high angles correspond to high resolution) collected on a 2D pixelated camera in the far field. Finally, the complex magnetic-refraction image is decoded from the scattering intensity algorithmically.

There are three main types of coherent imaging:

*Fourier transform holography (FTH)*: the first experimentally implemented approach to coherent magnetic imaging was FTH [306], where the interference between the sample's exit wave and a reference wave (for example scattered from a point-like feature) physically encodes the phase information on the camera. Thus, a single Fourier transformation provides an image of the sample's exit wave, cross-correlated with the reference. Experimentally, apertures in an opaque (typically gold) layer form the sample's FOV and the reference hole.

*CDI*: magnetic imaging by CDI captures the diffraction pattern in the far field and retrieves the sample's exit wave algorithmically. Here, the resolution is limited by the wavelength, the collected scattering angle, the signal above the noise, and the fidelity of the reconstructed phase. Oversampling (recording the diffraction pattern above the Nyquist frequency) and a well-defined transmission region ('support') are required to recover the phase of the complex-valued diffraction signal.

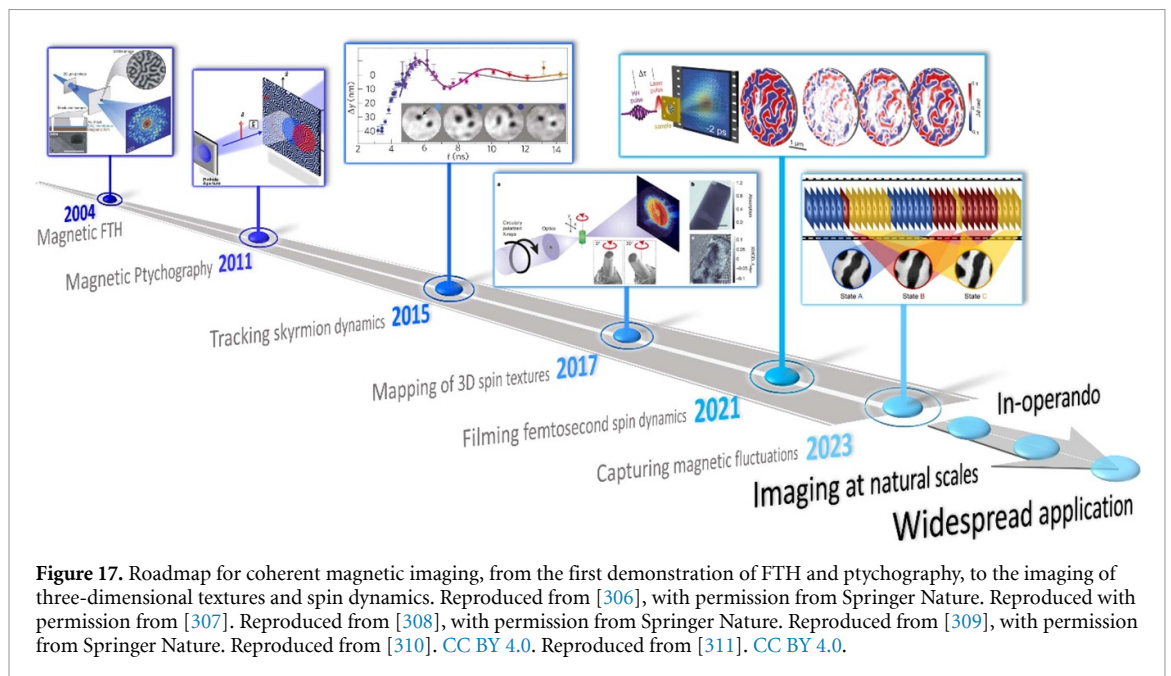
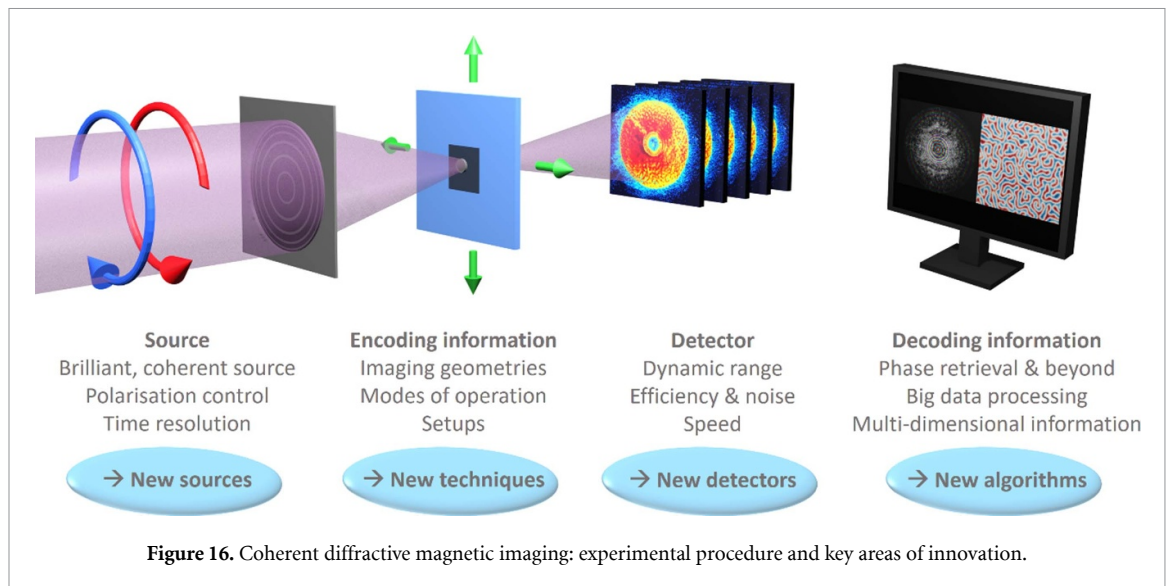
*Ptychography (PTY)*: in contrast to CDI and FTH, PTY [307, 313] is a scanning technique in which an extended object can be imaged. An overlap between adjacent illumination positions relaxes the requirement for a well-defined FOV and facilitates algorithmic retrieval of the amplitudes and phases of both the object and of the incident beam ('probe').

### Future advances of the technique

Technical developments of coherent sources and detectors as well as new reconstruction methods are driving advances in coherent magnetic imaging, leading to new science opportunities. Here we emphasise the diffraction-limited resolution, that provides a perspective for wavelength-limited resolution imaging. To date, both soft-x-ray PTY [318, 319] and CDI [316] have demonstrated magnetic imaging down to sub-10 nm, but the potential is far from being fully exploited. In the extreme-UV spectral range the combination of holography and CDI has pushed the spatial resolution to below the illumination wavelengths, reaching 16 nm spatial resolution at 21 nm wavelength [310]. Thus, coherent magnetic imaging with soft- and hard-x-rays [313] has the potential to access a few- or sub-nm resolutions in both 2D and 3D.

### Sources and detectors

Magnetic coherent imaging relies on synchrotron, FEL, and HHG radiation sources that deliver stable, bright beams with high flux, coherence and monochromaticity. Despite significant developments, the coherent flux



remains a limiting factor. There is a massive global effort to address this issue, both at large-scale facilities (e.g. 4th generation synchrotrons and seeded FELs) and with the development of sources based on HHG and laser–plasma accelerators. We anticipate that the resulting gain in coherent flux, together with a stable, quantifiable incident wavefront, will soon enable widespread, high-throughput use of coherent x-ray magnetic imaging, even at sub-wavelength resolution and in higher dimension (tomographic) experiments [309].

Although x-ray large-scale facilities are leading in their potential spatial resolution, extreme-UV scattering from HHG beams [320] demonstrate the promise of coherent magnetic imaging by reaching comparable performances at a ten-fold longer wavelength. Moreover, these laser-based sources are becoming ubiquitous in the field of ultrafast laser science due their stable fs-level accuracy and the method developments that a lab-scale setting with unrestricted beam-time allows for.

Detector improvements propel the prospects of coherent magnetic imaging through four main parameters: dynamic range, readout speed, quantum efficiency, and number of pixels. Quantitative high-resolution imaging requires capturing the dominant unscattered beam alongside the weaker ( $10^{-6}$  to  $10^{-9}$ ) high-angle magnetic scattering that is comparable with the shot- and instrument noise. The quantum efficiency should be high and the instrumental noise low to discern the scattering. Higher numbers of pixels (e.g. 16 Megapixel) allow to retain the oversampling necessary for CDI while increasing the number of resolution elements within the FOV.

On top of the above, the frame rates of conventional cameras (CCD, CMOS) determine the speed of the data collection, benefitting dynamic measurements, and shortening the acquisition ofptychographic images. However, we believe that the rise of event-based detectors (e.g. from the MediPix collaboration) would change the fundamental function of the camera: first, time-tagged photon hits would add correlative, i.e. quantum information to the signal, in addition to the evaluation of classical-field scattering. Second, the figure of merit in event-based acquisition is the hit rate (Mhits  $s^{-1}$ ), making the concept of ‘frames’ or ‘frame-rate’ an arbitrary post-processing feature in such detectors.

### **Algorithm developments**

A key aspect of coherent imaging is the decoding of information from the measured data with reconstruction algorithms. So far, the vast majority of magnetic coherent imaging has involved adaptations of non-magnetic approaches, because circular dichroism maintains the coherent imaging assumption [313, 321]. Future advances could open up possibilities to disentangle multiple polarisation modes [322], extract noisy or weak signals from a non-magnetic background, and resolve multi-dimensional information from single-shot images [323]. Indeed, algorithms have allowed to discern correlative dynamics from holographic data [311]. Extending this to higher dimensions and alternative datasets will allow material case studies at a previously unimaginable level.

One of the most exciting directions involves the ongoing revolution of machine learning and AI, which is impacting our entire society. First demonstrations of this for coherent imaging have shown potential towards more rapid, robust reconstructions of multidimensional data via neural networks [324], and PTY [325], allowing both significant increases in processing times, as well as the introduction of physics-guided reconstructions [326]. If these can be applied to magnetic imaging, and if the incident beam is sufficiently under control, there could be the possibility for quantitative extraction of material properties, transforming coherent imaging from an observation to a quantitative magnetometry technique.

### **Current and future use of the technique for material science**

Even though coherent x-ray magnetic imaging is relatively young, it has already contributed uniquely to several milestone discoveries in the field of magnetism, as shown in figure 17. For example, in 2015, time-resolved x-ray holography was used to confirm the quasi-particle nature of magnetic skyrmions in [Pt/CoB]<sub>x30</sub> up to the GHz regime [308] via tracking the skyrmion with sub-3 nm precision and thereby discovering that the equation of motion of a skyrmion includes a mass term, while static holography has been used to image the 3D shape of skyrmion tubes [327] and identify skyrmionics cocoons [328]. Later, PTY has allowed to resolve the actual profile of skyrmions in a similar [Ta/CoFeB/MgO]<sub>x16</sub> magnetic multilayer [329], to map the 3D configuration of Bloch point singularities [309, 317] as well as providing the first observation of magnetic vortex rings and torons [330], to name just a few. Moreover, in the ultrafast temporal range, coherent magnetic imaging for the first time unambiguously resolved the fs dynamics of the domain walls during laser-induced demagnetisation [310].

Indeed, the absence of lenses sets coherent imaging apart and removes associated challenges. For example, the intrinsic achromaticity of such purely far-field techniques makes them suitable for spectroscopic studies. With spatial resolutions below 5 nm, the imaging of magnetic domain wall profiles and spin textures at material defects are accessible [316], and even imaging of magnetic textures below the exchange length (e.g. Bloch points) is in reach. In combination with higher dimensional—spatial or temporal—imaging [304, 309], this advance promises access to magnetism beyond current micromagnetic models.

Moreover, the phase contrast capability of coherent magnetic imaging offers reduced dose damage and higher penetration depth compared to traditional absorption-based imaging techniques [312], thereby extending soft x-ray imaging to dose-sensitive and micrometre-thick magnetic samples [315].

### **New mechanisms targeted**

Coherent imaging currently addresses FMs, but it has a potential to impact other materials systems. For example, linear dichroism offers information on the spin orientation perpendicular to the incident beam, not only for FMs but also for antiferromagnets. Alternative geometries, such as coherent imaging in reflection, allow phase contrast imaging of AFM domain walls [331] and their fluctuations [332] in natural as well as in artificial antiferromagnets [333]. Combining with coherent x-ray imaging such as reflection PTY [334] or CDI would allow for resolving the details of the fluctuations of nanoscale textures. In PTY, structured probes (e.g. a vortex beam [335]) may be extended to study spin-orbit effects in, for example, topological spin textures, polar vortices in ferroelectrics, and other materials. With the significant increase in

coherent flux expected from the next generation of synchrotron sources, we envisage that the use of coherent approaches will expand significantly to include a wide variety of samples and physical phenomena.

### ***Mapping dynamic behaviours***

Coherent x-ray magnetic imaging has also opened access to new territories in the temporal domain by providing a means to perform imaging at the fastest x-ray sources (both large-scale FEL, and table-top HHG). For an in-depth discussion of x-ray dynamic imaging, see section 8.3. Although at an early stage, the first discoveries highlight the potential of these developments for materials science research. Examples include the observation of spatial variations of ultrafast optical demagnetisation [310, 336] and of local materials inhomogeneities which mediate fluctuations between energetically degenerate stripe domain states [311]. Similar dynamic studies have been demonstrated for antiferromagnets [332]. We expect that coherent x-ray imaging of antiferromagnets, either by x-ray magnetic linear dichroism or phase contrast, will soon be extended to the ultrafast regime. This will allow imaging of switching processes in AFM materials at their fundamental length and time scales and provide a previously unattainable understanding of the physics at play in this class of materials.

With the global rise of coherent x-ray facilities and the associated availability of instrumentation and data analysis support, we expect that coherent x-ray magnetic imaging will shift from singular highlight experiments to an important tool for the materials science community. Quantity and ease-of-use will enable quality by resolving standing challenges. Examples include the dynamical vector spin orientation profile of travelling magnetic domain walls and their correspondence with the 1D domain-wall model from the 1970s [337]; observing the generation dynamics of DMI- and frustration-stabilised magnetic skyrmions and their adherence with the Thiele equation that was developed for micrometre-scale, stray-field-stabilised bubble domains [338]; the imaging of irreversible processes, which are the dominant processes in real materials, either due to intrinsic (e.g. turbulences [339]) or extrinsic (e.g. inhomogeneities [340]) effects; and ultimately, also the imaging of magnetic materials at or near phase transitions.

### **Acknowledgements**

C D acknowledges funding from the Max Planck Society Lise Meitner Excellence Program and funding from the European Research Council 618 (ERC) under the ERC Starting Grant 3DNANOQUANT 619 101116043. F B acknowledges funding from the Helmholtz Young Investigator Group Program. This research used resources of the National Synchrotron Light Source II, a US DOE Office of Science User Facility operated for the DOE Office of Science by Brookhaven National Laboratory under Contract No. DE-SC0012704. O K gratefully acknowledges the Young Faculty Award from the National Quantum Science and Technology program of the Israeli Planning and Budgeting Committee and the Israel Science Foundation (Grant 1021/22).



### 8.3. Time-resolved x-ray imaging of magnetisation dynamics on the nanometre scale

Stefan Eisebitt<sup>1,2</sup>, Simone Finizio<sup>3</sup> and Bastian Pfau<sup>1</sup>

<sup>1</sup> Max Born Institute, Max-Born-Str. 2a, 12489 Berlin, Germany

<sup>2</sup> Institute of Optics and Atomic Physics, Technische Universität Berlin, Straße des 17. Juni 135, 10623 Berlin, Germany

<sup>3</sup> Swiss Light Source, Paul Scherrer Institute, Forschungsstrasse 111, 5232 Villigen PSI, Switzerland

#### Current state of the technique

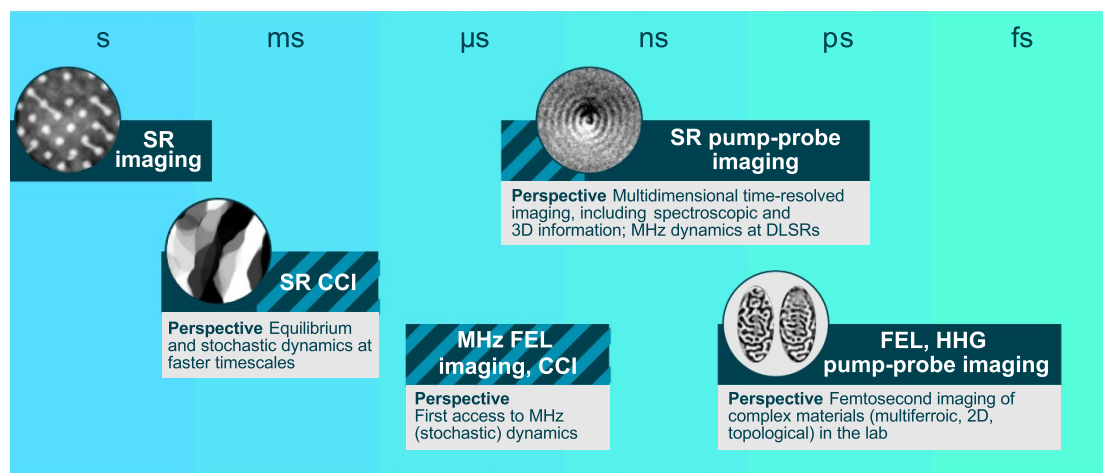
Imaging of lateral magnetisation distributions with sub-30 nm spatial resolution and the ability to add temporal resolution is provided by several x-ray imaging techniques. In general, sensitivity to the projection of the magnetisation is provided via XMCD/XMLD, at suitable resonant transitions which typically are located in the soft x-ray energy regime. Detection with 2D detectors allows for full-field imaging, typically in transmission from thin film samples, of amplitude or phase contrast. The full-field imaging approaches can be distinguished by the mode of image formation into lens-based TXM [341], interference-based FTH [308] and CDI based on iterative phase retrieval. See sections 8.1 and 8.2 for more details on these methods. Note that photoemission electron microscopy (PEEM) is not discussed in this chapter.

Given the readout times of current 2D detectors and signal-strength limitations, image acquisition happens on a timescale of seconds in a quasistatic regime [342]. Temporal resolution in the nano-, pico- and fs regime is derived from pump-probe schemes; see figure 18 for an overview. Most commonly, current pulses or AC currents are used for the study of magnetisation dynamics at SR sources, either to apply spin torques [343–345] or to be transduced into oscillating or pulsed magnetic fields [308, 346–348]. The temporal resolution at SR sources is fundamentally limited by the electron-bunch lengths to at least several tens of picoseconds. At FELs, fs temporal resolution can be achieved, typically in combination with laser-pulse excitation [336]. Here, single-shot pump-probe ‘snapshot’ experiments are in principle possible given that a single x-ray pulse contains sufficient photons to create a high-resolution image after a pump event [349]. Typically, however, such an approach will severely perturb or completely destroy the sample in the soft x-ray regime [349–351]. Multi-shot imaging of a static ferromagnetic domain pattern by FTH at an FEL without sample damage has been reported at 2.25 MHz [352].

Raster imaging can be implemented with fast point detectors in STXM. Pump-probe imaging in STXM is mostly performed in an asynchronous fashion, with the actual pump-probe delay for each recorded photon count individually sorted according to their phase relation with respect to the pump [353]. By measuring the time-of-arrival of each x-ray photon, temporal resolutions beyond the width of the x-ray pulses can be achieved, but those are still ultimately limited to the tens of ps scale [354]. An additional advantage of measuring the x-ray photon time-of-arrival is that it becomes possible to image non-locked dynamics by post-processing the acquired time-of-arrivals [355]. Imaging of laser-induced magnetisation dynamics after impulsive excitation by 2 ps (fwhm) laser pulses has recently been demonstrated in a STXM at 50 MHz repetition rate without observing sample damage at a laser fluence of  $3.1 \text{ mJ cm}^{-2}$  [356]. PTY as an extension of CDI is an alternative scanning-based imaging modality. Relying on slower 2D detectors, only time-resolved experiments where the excitation is directly locked to the master clock frequency of the SR source have been performed so far [357].

#### Future advances of the technique

Many of the foreseeable advances of the techniques are driven by the increasing availability of x-ray sources with significantly increased brilliance and hence coherent photon flux, allowing to focus to smaller spot sizes, and—together with the improvements in computational power—allowing to harness the information content from interference patterns of x-rays transmitted through the sample. In this fashion, the increased availability of x-rays with polarisation control at diffraction-limited storage rings (DLSRs) and FELs will push the spatial resolution in (static) imaging more routinely below the 10 nm limit, if sample damage is not an issue. At SR sources, the push to DLSRs is typically associated with bunch lengths typically exceeding 100 ps and hence a loss of temporal resolution over what is achievable today. On the other hand, the exploitation of sub-100 fs pulses at FELs for imaging of ultrafast magnetisation dynamics has so far not really started, although the spatial resolution achievable should match well with the temporal resolution. Here, sample damage of solid samples, often in the form of thin films on membranes, is a severe issue for systematic studies with a focus on materials science. Studies with hard instead of soft x-rays and hence less absorption and radiation damage may be able to mitigate this, possibly paving the way for video-like imaging of dynamic magnetisation phenomena with MHz repetition rates and non-perturbative pump-probe experiments accessing timescales down to fs.

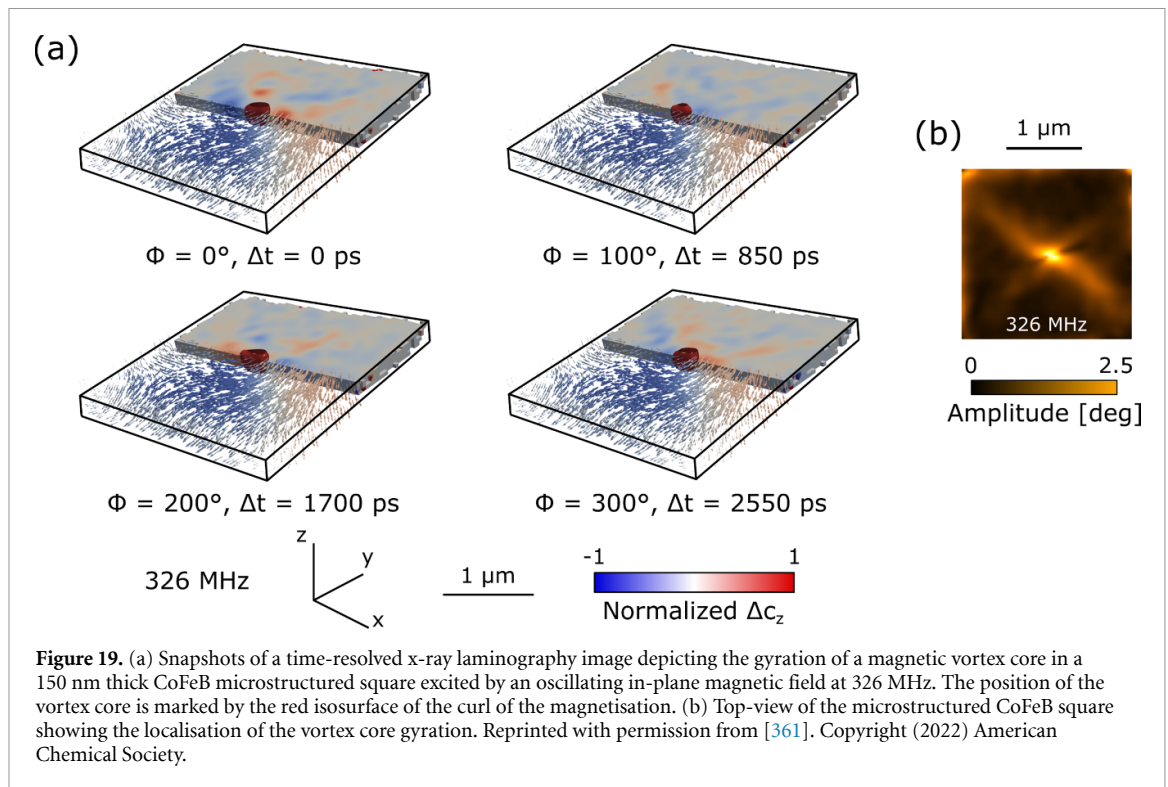


**Figure 18.** Overview over the timescales reachable with different time-resolving techniques of x-ray magnetisation imaging. Signal and detector constraints still limit the acquisition time of x-ray images to the second regime. Time-resolution on the relevant ns, ps and fs scale is only achieved with pump-probe methods, restricting the experiments to repeatable dynamics. New imaging methods, new sources, and new detectors open new research perspectives (shaded areas and text boxes), including investigating stochastic (non-repeatable) dynamics up to the MHz regime in the future. (from left to right) Reprinted figure with permission from [342], Copyright 2014 by the American Physical Society. Reprinted figure with permission from [347], Copyright 2019 by the American Physical Society. Reprinted figure with permission from [336], Copyright (2014) by the American Physical Society. Reproduced from [311], with permission from Springer Nature.

Another exciting development is the ability to probe stochastic dynamics, e.g. magnetic fluctuations in thermal equilibrium, rather than deterministic dynamics accessible in pump-probe experiments. X-ray photon correlation spectroscopy (XPCS) analyses the magnetisation fluctuations in Fourier space to measure characteristic timescales, typically, in the second to millisecond regime at SR sources [333, 358, 359]. Higher time-resolution is expected with the event of faster 2D detectors. Recently, XPCS found application also for studying nanosecond fluctuations of a skyrmion lattice at an FEL [360]. While XPCS can dissect the dynamics for different lengthscales, a real-space image is not retrieved. Coherent correlation imaging (CCI) overcomes this limitation, allowing video-like imaging of stochastic processes in a finite ensemble of possible states, where the photons required to generate an image can be collected at different temporal realisations of a state [311]. Exploiting this random repetition of states allows to stay away from destructive probe-beam fluences to generate high-resolution image frames in the video.

The x-ray source developments towards bright coherent beams do also take us into the third dimension: so far, almost all research with the imaging techniques described here has detected the projected magnetisation in thin magnetic films, but in the recent years studies of magnetisation vector within 3D structures have started to become possible [309]. It is a tremendous challenge to add temporal resolution on relevant temporal scales to 3D magnetisation imaging given that tomography scans are required, but first steps have been taken, e.g. via laminography (figure 19) [361]. Considering that the evolution of the vector field of the magnetisation within a sample volume is almost a terra incognita today at pretty much any time scale, there is large potential especially for material science applications, with bright sources allowing to use hard x-rays with high penetration in spite of weak magnetic contrast [309]. While PTY has shown great advances in static 3D imaging, the lack of 2D detectors recording the time-of-arrival of the x-ray photons in the soft x-ray energy range with a temporal resolution below the repetition rate of the light source is, however, still a critical issue for the development of time-resolved PTY imaging.

Regarding imaging of ultrafast magnetisation dynamics with temporal resolution in the fs up to tens of ps regime, the use of HHG in a laser laboratory is an exciting perspective. Given that HHG beams have a very high degree of coherence, they are suitable for time-resolved FTH, CDI, STXM and PTY. For the study of optically induced magnetisation dynamics, they combine large flexibility in generating synchronised excitation pulses at various wavelengths with the access option outside of a large-scale facility. Time-resolved magnetisation imaging has been performed via holographically assisted CDI exploiting Co M-edge (21 nm wavelength) XMCD, reaching 16 nm spatial resolution in static imaging of ferromagnetic domains [310]. From an imaging perspective, it is highly desirable to be able to work with XMCD-active resonances with much higher magnetic contrast in the future, such as the  $N_{5,4}$  edges of the lanthanides, the  $L_{3,2}$  edges of the 3d metals or even the lanthanide  $M_{4,5}$  edges. However, it is a tremendous challenge to achieve the photon flux and polarisation control at these photon energies required for imaging with HHG. Nevertheless, we note that first HHG-based time-resolved resonant magnetic scattering experiments at about 150 eV photon energy



have been successfully carried out recently [362]. HHG-based time-resolved imaging can be expected to move to progressively higher photon energies in the next years, allowing to reach further core-level resonances of interest in ultrafast magnetism.

### Current and future use of the technique for material science

The dominant contributions from time-resolved high-resolution imaging of magnetisation dynamics to materials science are expected to lie in the field of (quasi) 2D materials, thin films and heterostructured systems based on such ingredients. Given that many existing and proposed functional devices for storage, computation and sensing are realised in thin-film technologies, and given that generating a particular device function typically goes along with a lateral patterning of components on the nanometre scale together with (transient) local changes of the magnetisation state, this focus on ‘2D materials science’ is hardly a restriction. We envision that the following (partially overlapping) areas will benefit, in particular, from progress in time-resolved magnetisation imaging on the nanometre scale:

**All-optical switching** of magnetisation states via application of single or multiple light pulses is studied intensively as it is a potentially very fast and energy-efficient means of controlling magnetisation, e.g. for data-storage purposes. Considering spatio-temporal dynamics, the fs dynamics can be expected to occur on the shortest lengthscale of neighbouring atoms, with spin textures on the nanometre or even mesoscale exhibiting slower dynamics in the ps or even nanosecond regime. The ultrafast processes involved include transport of hot electrons and spins, also in the lateral dimension [336].

Intrinsically very fast dynamics on the interatomic scale is present in **AFM and multiferroic systems**: here, contrast can be generated by XMLD probing of the magnitude and direction of the Néel vector. Furthermore, the insensitivity of the imaging techniques discussed here to external electric and magnetic fields enables studies of device prototypes in operation. The atomic selectivity in resonant imaging is an additional asset to understand how dynamics relates to function in these compounds. Many of these systems are grown epitaxially on single-crystalline substrates, and we note that significant progress has been made in locally thinning such samples to soft x-ray transparency [305, 363].

Atomic selectivity is of great relevance also in the area of **2D magnets and heterostructures built up from 2D components**. The ability to stack vdW bonded 2D systems onto each other has fuelled the imagination of how to realise spin-based devices in such geometries. In the direction normal to the individual 2D constituents, spin currents may be generated, e.g. by the spin Hall effect or exploiting spin-momentum locking in suitable materials, possibly via optimised optical pulses [364]. The effect of spin

injection into different layers, including, e.g. spin accumulation at interfaces or the generation of spin–orbit torques can nicely be ‘dissected’ via imaging the local magnetisation within the constituent layers separately.

Focusing more on lateral dynamics, **magnonic devices** in addition to spintronic devices can be expected to be understood in more detail on the basis of time-resolved studies with nanometre spatial resolution [347, 348]. In passing, we note that x-ray-based reciprocal-space techniques may also be very valuable to address a frequency-momentum region not accessible by visible light.

In recent years, **topological spin structures in real space** and their dynamics have attracted considerable interest. So far, only few aspects of magnetic skyrmion dynamics have been accessible to study via time-resolved imaging [308, 343, 344], while some insight, in particular, in the optically triggered generation and decay of such structures have been enabled by time-resolved reciprocal-space studies with x-rays [365]. In particular to understand dynamics such as nucleation, topological transformation, or interaction with defects, time-resolved imaging of the transformations of magnetisation textures will be key to an understanding of the fundamental processes, hopefully paving the way to also exert control in applications. This is also true for the vast field of domain-wall motion, where lateral heterogeneity is critical to the dynamics and where x-ray imaging already has made unique contributions to understand changes in the spin structure during motion [366].

Understanding **rare failure events** (or even **rare success events**) at the odd time when something happens differently than it typically does, is a classical case for time resolved imaging—but, so far, has been elusive for processes at the nanometre level. Given the technological importance of understanding such events and optimising devices, true ‘video-mode’ imaging rather than pump-probe experiments may be very instrumental. As far as the dose required for magnetic imaging with a single exposure per frame goes, microsecond dynamics of magnetisation processes should become accessible at high-repetition rate at x-ray FELs. However, studies may remain limited by dose-driven damage. The observation of **stochastic dynamics** including (thermal) fluctuations is a related field as far as the need for video-mode imaging goes. There is hope that combined correlation and imaging approaches like CCI can be pushed to the relevant spatio-temporal range for the respective problem [311].

Finally, the observation of **magnetisation dynamics in 3D** is an experimental challenge, with first steps recently taken [309, 361]. In 3D, topologically nontrivial structures different from the ones that can exist in 2D are possible, with a hopfion being one example. It would be very interesting to observe such objects and their dynamics, which will be challenging even when only considering the size of the required data sets. While pump-probe schemes are most promising in this respect, the application of suitable pump events may represent another formidable challenge depending on the material or device to be studied.

## 9. Scanning Electron Microscopy with Polarization Analysis

*Robert Frömter and Mathias Kläui*

Institute of Physics, Johannes Gutenberg University Mainz, 55099 Mainz, Germany

### Current state of the technique

Scanning Electron Microscopy with Polarization Analysis (SEMPA, or spin-SEM) [367] is a powerful magnetic imaging technique that was developed already in the 1980s in Japan [368], USA [369], Germany [370], and Switzerland [371]. The method relies on measuring the spin polarisation of the SEs that are locally emitted from the surface of an FM while scanning it with a focused primary electron beam. The vector of spin polarisation of the SE is oriented opposite to the local magnetisation of the FM within the very small information depth of less than 1 nm from the surface. The spin polarisation of the SE is measured either by high-energy Mott scattering [368, 371] or by low-energy electron diffraction (SPLEED) at a W(001) single crystal [372].

By simultaneously measuring four symmetric scattering directions, two components of the spin-polarisation asymmetry are obtained simultaneously, usually corresponding to the two components of the magnetisation within the surface plane of the sample, as illustrated in figure 20. The missing out-of-plane component can be detected either permanently, by a 90° electrostatic deflection of the detector axis, or subsequently, by an electromagnetic spin rotator [373], or simply by tilting the sample [374, 375]. Due to the required vacuum environment, a wide range of sample temperatures can be easily accessed. Under optimum conditions a spatial resolution of 3 nm has been demonstrated imaging bit boundaries in a perpendicular recording medium [376]. Beam damage is not an issue on metallic samples for the used parameters, as most of the primary-beam energy is distributed into a large volume beneath the surface during the cascade process.

Owing to the very high surface sensitivity of SEMPA, dedicated UHV systems are needed to preserve the SE spin polarisation at the sample and to operate the SPLEED detector. While the early systems were custom built, a full-featured commercial UHV SEMPA instrument has been available from Scienta Omicron GmbH since the early 2000s, making the technique available to a wider community.

Exploiting the single-electron detection scheme of the SPLEED detector and adding a fast time-to-digital converter, it is possible to subdivide the recorded single-electron counts into successive time frames that are synchronised with a periodic external excitation of the sample, like an alternating current in a microstrip underneath giving rise to a modulated Oersted field. In this way, a movie of periodically driven magnetisation dynamics can be recorded at full duty cycle, i.e. without sacrificing counting efficiency [377, 378]. Even without further in-vacuum optimisations of the spin detector, a time resolution of 700 ps has been demonstrated [377].

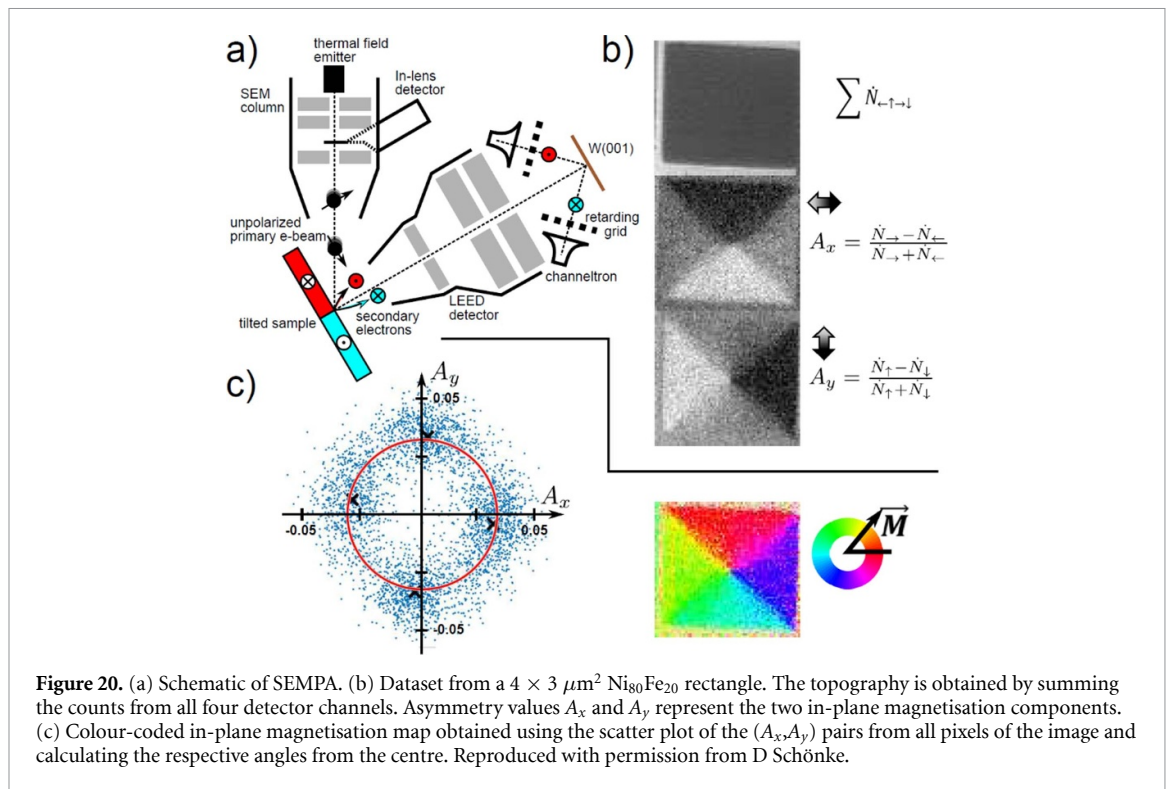
### Future advances of the technique

Even though SEMPA may appear as a mature method, there are consistently new technological developments in fast electronics, electron optics, and detectors, leading to new advanced capabilities. Most notable is the already mentioned addition of time resolution. The presently obtained value of 700 ps is not limited by the detection electronics or the microchannel plate electron multiplier—both would allow for a time resolution well below 100 ps. It is determined by the time jitter of the electrons passing through the SPLEED detector, where especially the passages at very low kinetic energy at the sample surface and at the retarding grid contribute. By further miniaturisation and using a faster filter, the total flight time and thus the jitter can be significantly reduced in an optimised detector design.

The above-stated maximum documented spatial resolution of 3 nm appears very competitive with other methods. However, it has been obtained using a unique UHV-SEM column that is equipped with a specifically optimised spherical and chromatic aberration corrector, allowing for 2 nm beam diameter at a working distance of 10 mm, energy of 20 keV and beam current of 1 nA [376]. Systems without aberration correction routinely obtain magnetic resolutions of 10–15 nm, which is a consequence of the required relatively large working distance and the high desired beam current, both increasing the lens aberrations. With today's progress in machine learning for image analysis, the focusing and astigmatism correction of the electron beam can be automatised, allowing for routine operation under optimised beam conditions, even on low-non-magnetic contrast samples.

There are also new developments concerning spin detectors. A significantly higher electron reflectivity at a similar spin sensitivity can be obtained when using the exchange interaction in scattering from an Fe layer at 6.3 eV instead of spin-orbit coupling in the W-SPLEED at 104.5 eV. Such a detector is presently marketed by FOCUS GmbH for SEMPA under the brand FERRUM [379]. We are presently characterising the prototype of an SPLEED detector using Au/Ir(001) that also promises improved sensitivity and service time,





while maintaining the two-component vector imaging capability. An improved detector efficiency enables higher spatial resolution, the study of lower-contrast samples, or reduced measurement times.

Given the reported progress in imaging on ambient-condition-stable samples, such as using capped surfaces or magnetic oxides, it could be interesting to revisit the potential for SEMPA as an add-on feature to standard SEM systems. While a Mott detector will operate under such reduced vacuum conditions (typically  $5 \times 10^{-7}$  mbar), earlier attempts suffered from the build-up of cracked hydrocarbons on the sample surface, greatly reducing the SE spin polarisation over extended exposure times [380]. A hydrocarbon-free vacuum and *in-situ* UV ozone cleaning may be applicable to mitigate such problems.

A shortcoming of SEMPA for certain applications, such as imaging of spin structures that are stabilised by external fields, is the apparent limitation to zero magnetic field. However, it has been demonstrated that in-plane fields of up to 100 mT can be applied during SEMPA imaging with appropriate corrections of the disturbing effects [381]. Locally confined out-of-plane fields of some mT can very likely be applied by using microcoils, as shown in in ref [308], and compensated for in a similar way as with the in-plane fields.

### Current and future use of the technique for material science

The main challenge in carrying out successful SEMPA investigations is the very low efficiency of all presently available detectors for the spin of free electrons. The contrast-to-noise ratio in the SEMPA images is dominated by the Poisson noise of the single-electron-counting statistics. Therefore, it increases linearly with the spin polarisation of the detected SE and the spin sensitivity of the detector, but it increases only with the square root of the measurement time and the primary beam current [372]. The latter is usually at variance with spatial resolution due to lens aberrations, so it cannot be arbitrarily increased. To obtain the same image quality at half the spin polarisation requires four times longer acquisition times. As a consequence, over the first decades many experiments were carried out on *in-situ* prepared uncapped ferromagnetic 3d-metal layers, where measurement times of several minutes were sufficient for good contrast, owing to the high spin polarisation of the low-energy SE of up to 50% for the case of iron. There are still today relevant scientific questions that can be addressed in this way, like the measurement of the DMI strength in epitaxial Co/Ir(111) [374], or studying the metamagnetic transition of epitaxial FeRh(001) films [382].

To image the magnetisation of *ex-situ* prepared samples two methods were used from early on: *in-situ* argon-ion sputter cleaning and *in-situ* dusting [383, 384] with an ultrathin iron layer or a combination of both. Both methods can in principle restore the full spin-polarisation of a 3d FM, however they will also modify the magnetic properties at the surface. While this is usually no problem with bulk samples, it may strongly influence ultrathin films and multilayers; which can even be exploited for depth profiling [385]. Sputter cleaning has been successfully applied for instance in a variety of hard-disk-drive related studies [376,

386], multiferroic materials [387, 388], exchange bias systems [389], and chiral domain walls in ferrimagnetic alloys [390].

Many magnetic systems that are in the focus of present-day physics and technology provide SE with less spin polarisation than the 3d metals and their alloys, or are too sensitive to allow for sputter cleaning or dusting. Here, at present the only way is to increase the measurement time, which—due to issues with stability—usually requires drift correction based on automatic pattern matching. The correction can be carried out either on the fly, or post-acquisition on an image stack, and is usually done based on the non-magnetic sum image. In this way, measurement times up to several days for a single SEMPA image have been realised. Drift correction is also a necessary procedure for time-resolved SEMPA, as the measurement time scales linearly with the number of frames in a movie.

The ability to measure systems with reduced SE spin-polarisation enables access to studying systems without any in-vacuum pretreatment, directly after inserting from atmosphere, provided the FM surface has been covered during deposition with a protective layer thin enough to retain some SE spin-polarisation and tight enough to prevent oxidation. E.g. a monolayer of graphene will protect a clean Ni(111) surface, while retaining 1/3 of the original spin polarisation from Ni [391]. In a systematic study using a sputter-deposited Pt wedge on 1.1 nm Co we could show that a cap layer of 1.1 nm Pt is sufficient to block oxidation, while retaining 1/4 of the SE spin polarisation from Co [392]. This allows, e.g. for studying chiral domain walls in sputtered multilayers [375, 393], as well as more recently topological spin structures such as merons and antimerons in layered synthetic antiferromagnets [394]. Due to its surface sensitivity, SEMPA is especially well suited to study synthetic antiferromagnets, whereas depth-integrating methods obtain only very little contrast due to the mutual compensation of the layers. By virtue of its bilayer-stepped interface, even the antiferromagnet Mn<sub>2</sub>Au could be imaged via its coupling to a permalloy layer that was protected by 2 nm of highly spin-transparent SiN<sub>x</sub> [395].

There are even magnetic materials where the surface is intrinsically stable against oxidation and that can be directly imaged in SEMPA, like the half-metallic perovskite manganite La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> [396].

Being set up in a standard UHV lab environment, SEMPA offers unlimited access times in contrast to synchrotron-based x-ray methods. The capability for imaging untreated, *ex-situ* prepared, capped samples is a strong new feature, since it allows for direct interchange and comparison with other spatially resolving as well as integrating methods. A novel route to understand magnetic systems better, is thus the study of systems using complementary techniques. As previously shown, a comparison of magnetic imaging with for instance NV microscopy that detects the stray field can give additional insights into a sample's physics [397], as both techniques determine complementary quantities and have different depth and lateral sensitivities. In particular, given the very strong surface sensitivity, SEMPA will be a valuable tool in such studies where one combines SEMPA with for instance a more bulk sensitive or a 3D-imaging method.

## Acknowledgements

This work has received funding from the ERC under the European Union's Horizon 2020 Research and Innovation Programme (Synergy Grant No. 856538, Project '3D-MAGiC'). It has also been supported by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation), Grant TRR 173–268565370 (Project A01, B02) and Grant SPP 2137 Skyrmionics—403502522.

## 10. Transmission electron microscopy techniques

### 10.1. Lorentz transmission electron microscopy

Fehmi Sami Yasin<sup>1</sup>, Benjamin J McMorran<sup>2</sup>, Shinichiro Seki<sup>3</sup> and Xiuzhen Yu<sup>1</sup>

<sup>1</sup> Center for Emergent Matter Science (CEMS), RIKEN, Wako, Japan

<sup>2</sup> Department of Physics, University of Oregon, Eugene, OR, United States of America

<sup>3</sup> Department of Applied Physics and Institute of Engineering Innovation, University of Tokyo, Tokyo, Japan

#### Current state of the technique

For decades, LTEM has been used to study magnetic configurations in thin magnets. In LTEM (right-hand-side (RHS), figure 21), a plane wave of electrons is transmitted through the magnetic field or vector potential ( $\mathbf{A}(\mathbf{r}_\perp, z)$ ) distribution, imparting a phase (Lorentz deflection in the classical picture) onto the passing electron wavefronts of the form  $\phi_m(\mathbf{r}_\perp) = \frac{e}{\hbar} \int_L \mathbf{A}(\mathbf{r}_\perp, z) \cdot d\mathbf{r}$  [398]. The distorted wavefronts form measurable contrast after propagating some distance, revealed by defocussing ( $\Delta f$ ) the image. This allows for phase retrieval using the transport-of-intensity (TIE) equation and an over-, under-, and in-focus image [399], although recent improvements require only a single defocussed image on some samples [400].

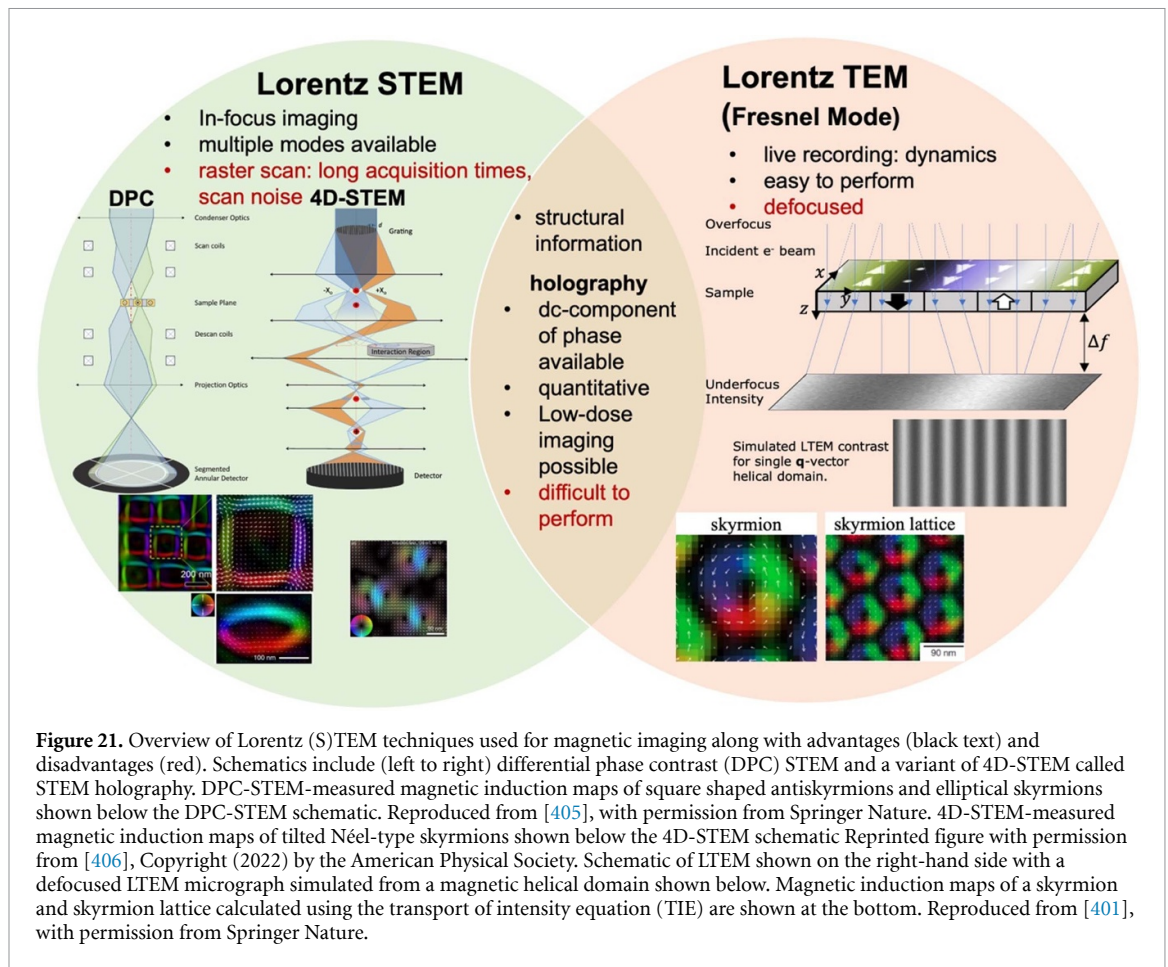
Consider the magnetic field distribution in a proper screw spin propagation with period  $\lambda_{hel}$  along the  $y$ -axis within a flat specimen with thickness  $t$ , which can be expressed as  $\mathbf{B}_\perp = B_0 \sin\left(\frac{2\pi}{\lambda_{hel}}y\right) \hat{x}$ . The phase difference between the origin and  $y$ ,  $\Delta\phi_m(\mathbf{r}_\perp) = \frac{e}{\hbar} \iint \nabla \times \mathbf{A} \cdot d\mathbf{S}$  (here  $\mathbf{S} = \hat{x}$ ) can be rewritten using Stokes theorem to obtain  $\phi_m(y) = \frac{e}{\hbar} \int_{-\frac{t}{2}}^{+\frac{t}{2}} \int_0^y \mathbf{B}_\perp dy dz = -\frac{e\lambda_{hel}}{\hbar} B_0 t \cos\left(\frac{2\pi}{\lambda_{hel}}y\right) \hat{x}$ . The simulated image near the bottom RHS of figure 21 illustrates a defocussed image from the above field distribution. This powerful technique and phase retrieval algorithm were used to observe magnetic skyrmions in real space for the first time [401] (bottom RHS of figure 21).

STEM mode offers an alternative optical setup, and two types of LSTEM imaging are outlined on the left-hand-side of figure 21 including differential phase contrast STEM (DPC-STEM) [402] and interferometric 4D-STEM [403, 404]. By measuring beam shifts across a segmented annular detector or imaging interferograms on a pixelated grid detector (e.g. charged couple device), DPC-STEM and interferometric STEM directly measure the in-plane magnetic induction and phase, respectively. DPC-STEM-measured in-plane magnetic induction maps of an antiskyrmion square lattice and elliptically shaped skyrmion present in the recently discovered  $(\text{Fe}_{0.63}\text{Ni}_{0.3}\text{Pd}_{0.07})_3\text{P}$  [405] are shown below the DPC-STEM schematic while an interferometric 4D-STEM-measured magnetic induction map of Néel-type skyrmions in a room temperature 2D polar magnet [406] is shown below the 4D-STEM schematic. Recent advances in detector acquisition speed and dynamic range have resulted in an explosion of LTEM work studying spin texture dynamics [407–409], and the rise of Lorentz 4D-STEM magnetic imaging. 4D-STEM was recently used by Nguyen *et al* to resolve antiferromagnetic magnetic ordering in  $\text{Fe}_2\text{As}$  [410].

#### Future advances of the technique

LTEM has been a crucial tool for imaging magnetic materials, with several promising directions that could further establish LTEM as an indispensable technique for magnetism research. One major advancement is the real-time imaging of dynamic magnetic phenomena such as magnetic domain wall motion and spin texture transformations. This is achieved by holding the LTEM image at a constant defocus, rendering magnetic contrast on the detector that may be observed via live recording while applying an external stimulus such as magnetic field, electric current or heating. Recent developments in direct electron detector technology have significantly improved the frame rates and electron-dose efficiency. For example, the second-generation Electron Microscope Pixel Array Detector (EMPAD-G2) can achieve tens of thousands of frames per second (fps) and offers dynamic ranges capable of single electron detection up to 180 pA per pixel [411]. Optically pulsed electron sources and fast beam blankers (choppers) [412] enable ultrafast TEM (UTEM) [413] and even ultrafast STEM, also known as 5D-STEM [414]. These types of advancements enable real-time imaging of dynamic magnetic events while applying external stimuli. Additionally, advancements in TEM holder technology, including holders combining several features such as heating/cooling, electric/thermal current, and the addition of piezo probes for applying controlled mechanical pressure on samples, have expanded the possibilities for varied magnetic responses and experimental designs that better represent real-world device conditions. For example, a recent experimental design allows for *in-situ* magnetic imaging and electronic transport measurements, enabling the direct correlation between the anomalous Hall effect and real-space spin textures for the first time [415].

These detector advancements have also enabled 4D-STEM holography and PTY techniques that have, in some cases, remained dormant for over 40 years [404]. When configured in Lorentz mode, these techniques



**Figure 21.** Overview of Lorentz (S)TEM techniques used for magnetic imaging along with advantages (black text) and disadvantages (red). Schematics include (left to right) differential phase contrast (DPC) STEM and a variant of 4D-STEM called STEM holography. DPC-STEM-measured magnetic induction maps of square shaped antiskyrmions and elliptical skyrmions shown below the DPC-STEM schematic. Reproduced from [405], with permission from Springer Nature. 4D-STEM-measured magnetic induction maps of tilted Néel-type skyrmions shown below the 4D-STEM schematic. Reprinted figure with permission from [406], Copyright (2022) by the American Physical Society. Schematic of LTEM shown on the right-hand side with a defocused LTEM micrograph simulated from a magnetic helical domain shown below. Magnetic induction maps of a skyrmion and skyrmion lattice calculated using the transport of intensity equation (TIE) are shown at the bottom. Reproduced from [401], with permission from Springer Nature.

allow for multi-modal imaging, simultaneously mapping the magnetic field, electrostatic field, strain, and polarisation. PTY was recently used to beat the usual diffraction limit and phase resolution in Lorentz mode [416]. This unprecedented ability is expected to advance further with the adoption of phase-structured electron beams using phase plates, diffraction holograms, and programmable apertures as well as the development of TEMs optimised for magnetic imaging with aberration correction and improved sample positioning for higher resolving and magnifying power in LTEM.

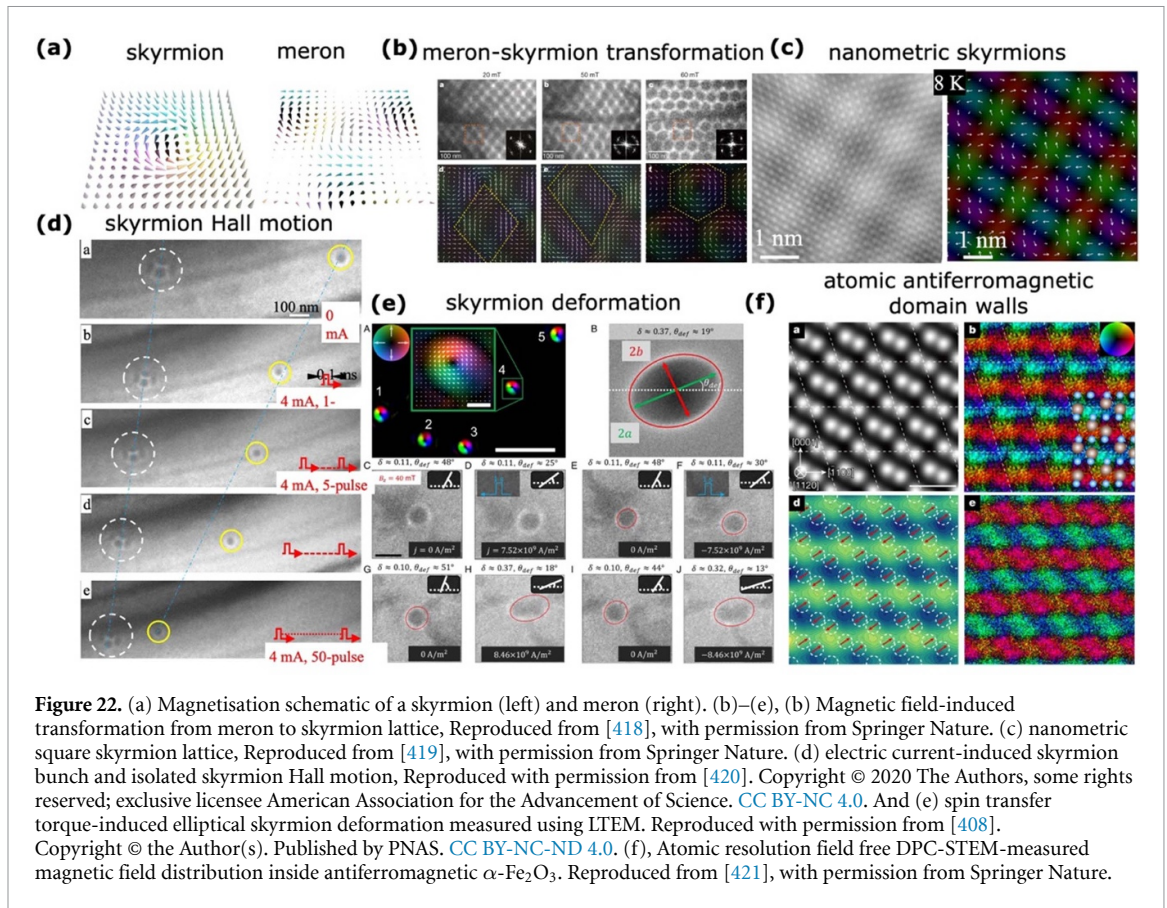
Another promising advancement in LTEM is the progress in 3D imaging using tomography holders and reconstruction algorithms. Traditional 3D magnetic imaging requires the acquisition of 2D projections of the 3D magnetic field at many tilt angles which can be used to back project the full 3D magnetic field. It is typically labour-intensive, requiring fine alignments of the FOV for every tilt angle along two perpendicular tilt axes. However, machine learning-driven software has begun automating much of this process, reducing the human effort involved in data acquisition. Additionally, the development of multibeam techniques such as multibeam electron diffraction [417] may enable the imaging of 3D dynamic magnetic phenomena when configured in Lorentz mode. With simultaneous acquisition of multiple tilt angles, these techniques allow for real-time viewing during the application of external stimuli, further enhancing the understanding of 3D magnetic features.

Finally, the integration of AI into LTEM experiments has the potential to revolutionise the user experience and enable the design of previously unthinkable experiments. AI can be utilised for real-time electronic aberration correction, sample and beam drift correction, and real-time magnetic induction mapping, to name a few operative procedures. In post-processing, AI can assist in sifting through large data sets which are often generated in 3D LTEM imaging and 4D-STEM experiments and help correlate the magnetic real space imaging with other analytical techniques such as transport measurements, strain mapping, x-ray microscopy, or pump-probe ultrafast microscopy, providing a comprehensive understanding of the magnetic structure within materials and its response to external stimuli.

### Current and future use of the technique for material science

As mentioned earlier, one of the most powerful capabilities available in LTEM is the real space measurement of spin textures and their response to external stimuli in live time. One example is the recent discovery and





real space observation of magnetic (anti)merons [418], spin textures pictured in figure 22(a) that twist around half the unit sphere from their cores to their border, granting them half integer topological charge,  $N_{(anti)meron} = \frac{1}{4\pi} \iint \mathbf{n} \cdot \left( \frac{\partial \mathbf{n}}{\partial x} \times \frac{\partial \mathbf{n}}{\partial y} \right) dx dy = \pm 1/2$ . They contrast with magnetic skyrmions, whirling spin textures also pictured in figure 22(a) that traverse the entire unit sphere once from core to edge, with a corresponding integer topological charge,  $N_{Skyrmion} = -1$ . Using LTEM, Yu *et al* could track the merons as they warped in real space while under an increasing magnetic field towards the final skyrmion lattice state illustrated in figure 22(b).

As more magnetic imaging-specific TEMs are developed with an eye towards improving the magnifying power and correcting more lens aberrations, real space observations of nanometric spin textures such as the square skyrmion lattice shown in figure 22(c) [419] and their dynamics are expected to become more commonplace. Advancements in detector speed and sensitivity will also enable previously unobtainable measurements of magnetic spin texture dynamics such as the electric current driven motion of magnetic skyrmions and their bunches shown in figure 22(d) [420]. With current technology, LTEM is limited to acquiring still images between current pulse applications since the spin texture motion is many orders of magnitude faster than the camera readout speeds. While such frame rates may not be obtainable in the near future, and indeed the most promising path may lie in ultrafast pump-probe LTEM methods currently under development, magnetic thin plates have already shown under certain conditions to host spin texture responses slow enough and/or in a steady state deformation to view via live recording. Such an electric current-induced skyrmion response was recently measured and quantitatively characterised [408]. LTEM clearly shows the spin transfer torque-driven elliptical deformation and stretching of skyrmion domain walls, illustrated in figure 22(e).

One magnetic imaging-specific STEM has already been developed at the University of Tokyo. MARS, or magnetic field free atomic resolution STEM, enables the direct observation of atoms in a magnetic field free environment. This achievement was made possible by the design of a new objective lens (OL) that maintains the resolving power of the microscope while isolating the sample from any magnetic fields and including aberration correction. Using this TEM configured for DPC-STEM, Shibata *et al* resolved the atomic magnetic field distribution inside AFM haematite ( $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>), shown in figure 22(f) [421].

2D materials have also seen a surge of interest due to their emergent properties such as superconductivity and magnetism that are measurable using LTEM [422]. While magnetic thin plates generally avoid electron



beam knock-on damage, 2D materials composed of lighter elements may not. To avoid such radiation damage, LTEM imaging may be performed at lower beam energies.

In conclusion, magnetic imaging using LTEM is an active and exciting field with several promising directions for future advancements. Real-time imaging of dynamic magnetic phenomena, 3D imaging, integration of AI during experiments and post-processing, and improvements in resolution and aberration correction are some of the areas that hold great potential for unlocking new insights into the behaviour of magnetic materials. These advancements promise to enable researchers to design and conduct experiments that were not previously feasible, opening new avenues for understanding and manipulating magnetic structures at the nanoscale.

### **Acknowledgements**

This work was supported in part by Grants-In-Aid for Scientific Research (A) (Grant No. 19H00660) from the Japan Society for the Promotion of Science and CREST program (Grant Nos. JPMJCR1874, JPMJCR20T1) from the JST, Japan.

## 10.2. Electron holography

Axel Lubk<sup>1,2</sup> and Daniel Wolf<sup>1</sup>

<sup>1</sup> Leibniz Institute for Solid State and Materials Research Dresden, Dresden, Germany

<sup>2</sup> Institute of Solid State and Materials Physics, TU Dresden, 01062 Dresden, Germany

### Current state of the technique

In the general context of magnetic imaging, EH refers to a set of TEM techniques that reconstruct the Aharonov–Bohm phase shift  $\varphi_{\text{mag}}$ , imprinted on the beam electrons' (beam direction  $z$ ) quantum wave function during transmission through a thin magnetic sample. Today's electron microscopes allow the recording of 2D phase maps  $\varphi_{\text{mag}}(x, y)$  in various magnetic-field-free imaging modi (referred to as L(S)TEM, see section 10.1), namely off-axis EH (figure 23(a)), differential defocus and focal series inline EH, as well as DPC that provide a spatial resolution approaching the sub-nm regime [421] and a phase resolution beyond  $2\pi/1000$  [423]. Under kinematical scattering conditions these maps can be quantitatively evaluated in terms of projected in-plane magnetic flux densities via their directional derivatives, i.e.  $\partial_x \varphi_{\text{mag}} = e/\hbar \int_{-\infty}^{\infty} B_y dz$  and  $\partial_y \varphi_{\text{mag}} = -e/\hbar \int_{-\infty}^{\infty} B_x dz$ . Differential defocus EH (also referred to as TIE method) retrieves that information from two mutually defocused Lorentz LTEM images [424]. Focal series EH analyses a whole set of defocused images, while off-axis EH utilises the interference between the electron wave passing both the object as well as a (field)-free reference region [425]. DPC deviates from the above interferometric techniques in that local deflections of a focused electron beam are measured while scanning the beam over the sample (STEM), which may then be translated into the Aharonov–Bohm phase exploiting the semiclassical relationship between the wave front (phase) and classical trajectories (deflection). EH has been mainly employed for studying nanoscale magnetic field configurations governed by micromagnetics, including, e.g. nanoparticles [426], nanowires [427], domain walls [427, 428], skyrmions [429, 430], but also assembled 3D MRAM devices [431]. That includes the *in-situ* application of additional stimuli notably magnetic fields (mostly in  $z$ -direction, typically up to several 100 mT) and temperature (ranging from 4 K to 1000 K). Combining EH and electron tomography enables the reconstruction of magnetic vector fields in 3D (figure 23(b)) [427, 432]. To reconstruct the flux density vector field  $B$  being subject to  $\text{div } B = 0$  at least two tilt series around mutually perpendicular axes are necessary. This approach has been previously used to study 3D magnetisation textures such as transverse and longitudinal vortices in magnetic nanowires [427, 432] as well as Bloch skyrmion tubes in FeGe helimagnets [433].

### Future advances of the technique

Over the last decade EH techniques have seen drastically improved precision and accuracy by employing high-brightness cold-field emission guns, improved aberration-corrected field free electron optics, automated long-term hologram series acquisition involving drift correction, improved detector technology, such as direct electron detection, and increased overall instrumental stability [423]. This quest will continue, e.g. with the arrival of next generation detectors and the increasing implementation of fully-fledged auto-alignment of the whole TEM column. Considering that spatial resolution in EH magnetic imaging is noise-limited to a large extent (the magnetic phase shift is small compared to the dominating electrostatic contribution) this will further push the resolution to atomic length scales provided by the optics of the microscope.

Notwithstanding, the trend towards improved *in-situ* capabilities, better integration with other experiments and methods, and increasing the time-resolution [434] will probably be the dominating theme in the next couple of years. Another emerging trend concerns the increasing implementation of automated evaluation of characteristic magnetic quantities, such as exchange stiffness, magnetocrystalline anisotropies, or topological charge, profiting from modern machine learning algorithms. The remainder of this section attempts a short outlook into these impeding advances of EH.

*In-situ* TEM techniques combine the optical power and versatility of state-of-the-art transmission electron microscopes with miniaturised instrumentation, which allow to manipulate and stimulate the magnetic sample within the small space available in TEM. With the advent of MEMS, 3D printing and other nanofabrication and micromachining technologies the last decade has seen a fast progress and proliferation of both commercial and custom-made sample holders (see e.g. [435] for the latter) for *in-situ* TEM that are compatible with the sample stages of the main TEM manufacturers and hence usable by a large community. These holders allow the application of, e.g. temperature, currents, electric and magnetic fields, light, strain, liquids, and gases. While all these capabilities may be fruitfully combined with EH to study transient and response phenomena (e.g. magnetostriction, magnetoionics) particularly strong impact may be expected from advances concerning the application of magnetic vector-fields, cryogenic temperatures, and sample rotation around three independent axes, in particular if these are combined. The reason is the rich phase

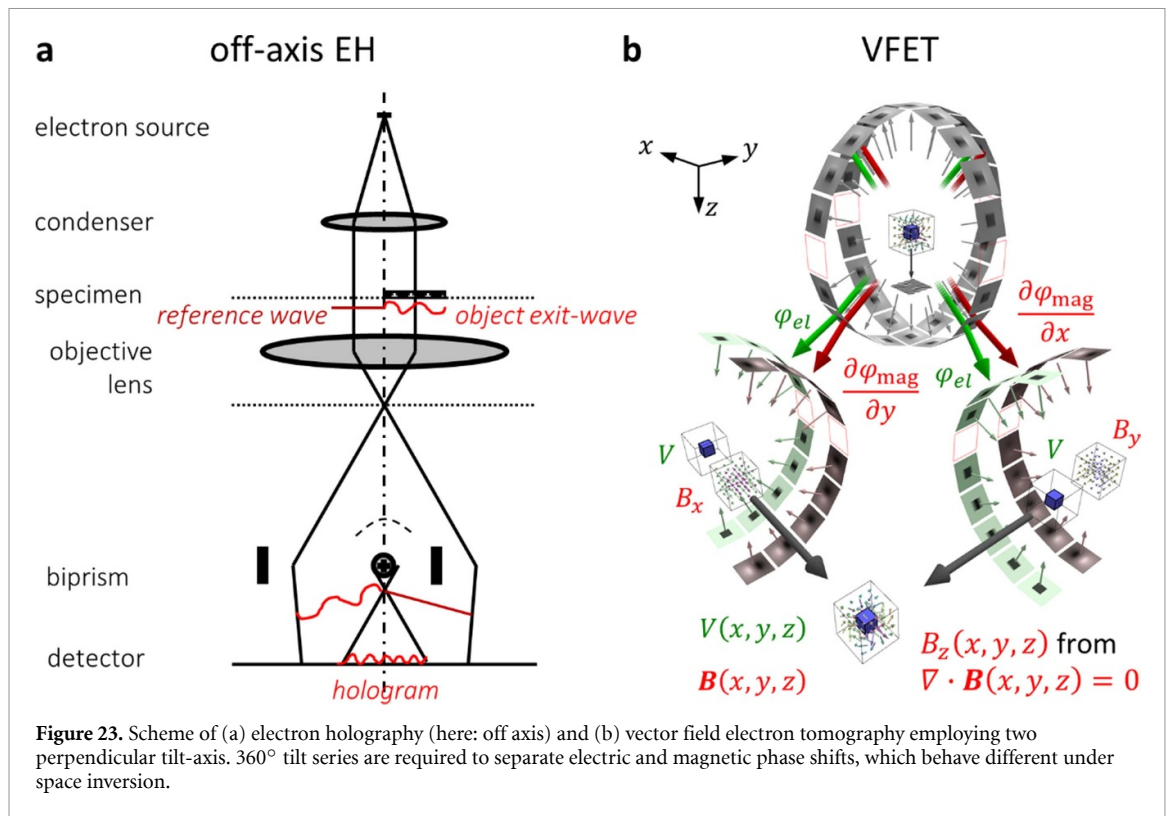
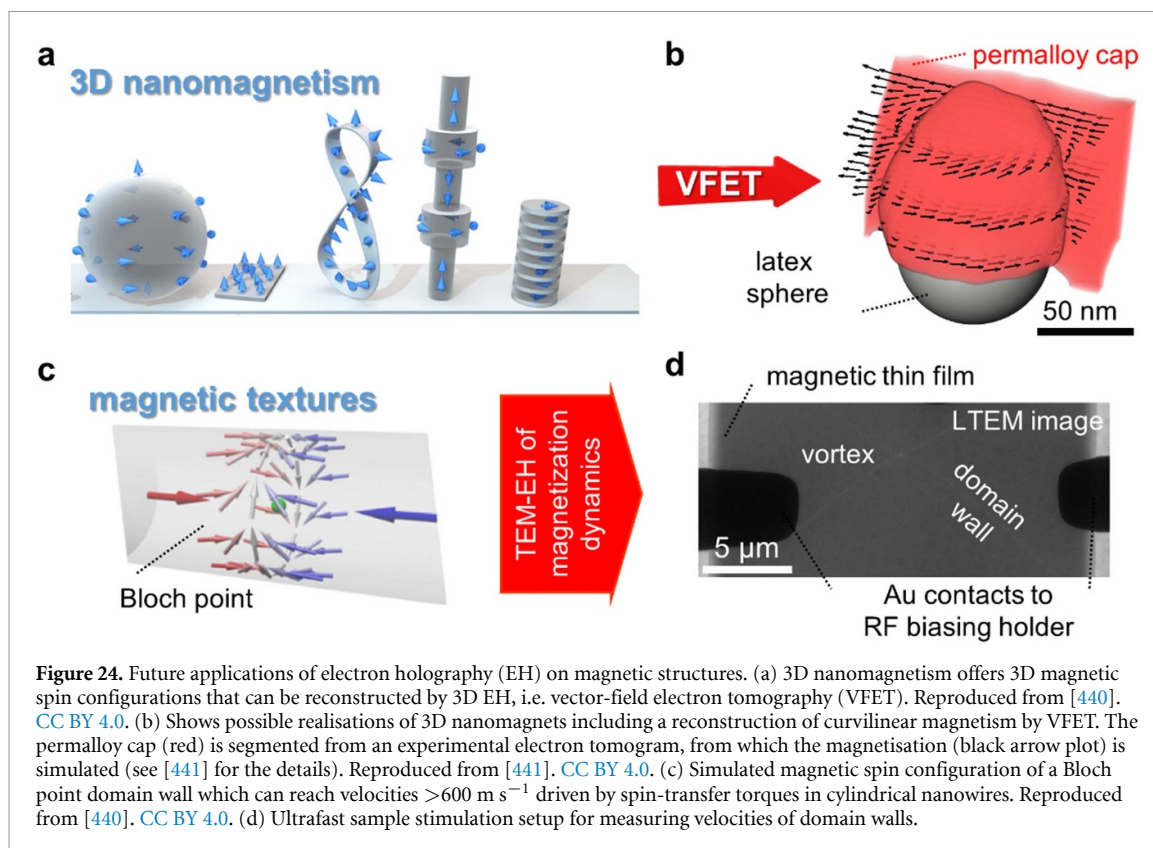


diagram and phenomenology of mesoscale frustrated 3D magnetic configurations, such as magnetic skyrmions, which become accessible by such *in-situ* TEM studies. Magnetic vector-field TEM holders, which incorporate a beam deflection correction, are currently under development by various labs as are continuous flow liquid helium sample holders, which facilitate long-term cooling and counter-heating in wide temperature range starting from approximately 5 K. Here the cooling would also mitigate some beam damage effects (i.e. radiolysis) hampering the study of beam sensitive magnetic materials. Three-axis 360° rotation holders as required for optimal magnetic vector-field tomography that require redesign of the OL in order to gain more space around the sample are also in their early development stages.

A second related development focuses on a better integration of EH with other experiments in the sense that all measurements are carried out at the same TEM sample, preferably *in-situ* but also *ex-situ*. Notable example for *in-situ* experiments is the combination of EH with current and entropy transport measurements employing modified electric biasing holders [436]. Such experiments allow a direct correlation of observed magnetic field configurations with longitudinal and transverse magnetotransport and hence the zoo of Hall and Nernst effects. Possible *ex-situ* registered measurements pertain, for instance, to XMCD based imaging methods facilitating the mapping of larger sample regions.

Significant improvements of the time-resolution of EH are closely linked to the ongoing development and proliferation of a variety of UTEM techniques, e.g. employing laser-pulsed electron guns, ultrafast beam deflection or very fast detectors in combination with ultrafast sample manipulation, e.g. via short current pulses. For example, the rotation of a magnetic vortex state in a thin permalloy film has been recently imaged using laser-pulsed ultrafast LTEM and a radio-frequency biasing sample holder, which has been synchronised with the gun [437]. Emerging ultrafast cavity and MEMS blankers fitted in existing TEM columns will further broaden the accessibility of such experiments operating in the sub-ns time resolution domain.

Although EH has been proven a useful technique for resolving magnetic field configuration on the nanoscale it typically requires comparison to simulations (e.g. micromagnetic modelling) to reveal the magnetisation texture or structure underlying the magnetic flux density. The foremost reason is that EH is blind with respect to magnetic field strength  $H$ , corresponding to the irrotational part of the magnetisation in magnetostatics. Consequently, to reveal the whole magnetisation distribution,  $H$  has to be provided by additional means, such as simulations. Therefore, ongoing research efforts elaborate on combining EH,



micromagnetic (or other) modelling, and machine learning to retrieve the full magnetisation, particular micromagnetic energy terms or characteristic magnetic parameters such as the exchange stiffness.

### Current and future use of the technique for material science

EH will remain a mainstay in the study of nanoscale magnetic field configurations determined by the competition of symmetric and antisymmetric exchange, magnetocrystalline anisotropy, Zeeman energy and dipolar interactions. Here the previously described advances will allow a significant extension of the accessible phase diagram spanned by temperature, external magnetic field, etc, particularly benefitting the study of the large class of anisotropic magnets, low temperature magnetic phases, multiferroics and topological magnetic materials. Improved vector field reconstruction incorporating simulations will address a larger class of topological non-trivial magnetic solitons, such as skyrmion tubes, braids [438], chiral bobbars [439] and Bloch points [440], including their topological charges, notably those of Néel-type exhibiting significant magnetic charges and hence field strengths  $H$ . Improved vector-field electron tomography will also help in revealing the impact of effective magnetic interactions such as anisotropies and antisymmetric exchange emerging in curved magnetic thin films (figures 24(a,b)) [441]. Combination of EH and transport measurements, i.e. the measurement of longitudinal and transversal transport on well-determined magnetisation structures, has the potential to shed new light on the origins and amplitude of anomalous and topological magnetic and thermal Hall effects. The combination of cryogenic EH and electron tomography as well as other TEM techniques such as spectroscopy may also spark a revival of high-resolution studies of magnetic field vortices in type-II superconductors such as high  $T_c$  oxides.

Further advances in signal and spatial resolution bear the prospect of studying 2D magnetic materials as well as magnetic configurations emerging at surfaces and interfaces, e.g. in complex oxides, where symmetry breaking and a stronger localisation of electrons can lead to an increased tendency for magnetic order, strong uniaxial anisotropies as well as antisymmetric exchange, resulting in a wealth of complex magnetic states. Further resolution improvement is also instrumental for the application of EH techniques on antiferromagnets facilitating microscopic insights into antiferromagnetic order and field distribution. That notably includes characterisation of uncompensated spins at lattice defects or interfaces. A bit further down the road waits the detection of magnetic monopoles in spin ice materials, e.g. of the pyrochlore class.

Finally, development of ultrafast magnetic imaging (both LTEM and EH) will open the door to the quantitative study of transient and non-equilibrium magnetic states and hence the wide field of spintronics. That includes the characterisation of the motion and pinning of magnetic domain walls in nanostructures, such as Bloch points (figures 24(c,d) or transverse vortex domain walls, of various types of skyrmions, by also spin thermalisation at magnetic phase transitions, etc.

### Acknowledgments

The authors acknowledge financial support of the DFG by the Collaborative Research Center SFB 1143 (Project-id 247310070).



## Data availability statement

No new data were created or analysed in this study.

## ORCID iDs

D V Christensen  <https://orcid.org/0000-0003-0048-7595>  
U Staub  <https://orcid.org/0000-0003-2035-3367>  
T R Devidas  <https://orcid.org/0000-0002-6855-7505>  
B Kalisky  <https://orcid.org/0000-0002-1270-2670>  
A Huck  <https://orcid.org/0000-0002-2354-5922>  
D A Broadway  <https://orcid.org/0000-0002-7375-8766>  
K Wagner  <https://orcid.org/0000-0003-4252-0584>  
P Maletinsky  <https://orcid.org/0000-0003-1699-388X>  
T van der Sar  <https://orcid.org/0000-0002-6197-4808>  
C R Du  <https://orcid.org/0000-0001-8063-7711>  
D Collomb  <https://orcid.org/0000-0001-5591-8802>  
S Bending  <https://orcid.org/0000-0002-4474-2554>  
A Oral  <https://orcid.org/0000-0001-8972-4862>  
H J Hug  <https://orcid.org/0000-0002-1281-4138>  
A-O Mandru  <https://orcid.org/0000-0001-8778-8571>  
V Neu  <https://orcid.org/0000-0001-9170-1280>  
H W Schumacher  <https://orcid.org/0000-0001-7100-682X>  
S Sievers  <https://orcid.org/0000-0002-6848-8984>  
H Saito  <https://orcid.org/0009-0000-9554-3590>  
A A Khajetoorians  <https://orcid.org/0000-0002-6685-9307>  
N Hauptmann  <https://orcid.org/0000-0003-4264-2493>  
S Baumann  <https://orcid.org/0000-0002-3821-494X>  
A Eichler  <https://orcid.org/0000-0001-6757-3442>  
C L Degen  <https://orcid.org/0000-0003-2432-4301>  
J McCord  <https://orcid.org/0000-0003-0237-6450>  
M Vogel  <https://orcid.org/0000-0001-6253-5965>  
M Fiebig  <https://orcid.org/0000-0003-4998-7179>  
P Fischer  <https://orcid.org/0000-0002-9824-9343>  
A Hierro-Rodriguez  <https://orcid.org/0000-0001-6600-7801>  
S Finizio  <https://orcid.org/0000-0002-1792-0626>  
S S Dhesi  <https://orcid.org/0000-0003-4966-0002>  
C Donnelly  <https://orcid.org/0000-0002-9942-2419>  
F Büttner  <https://orcid.org/0000-0002-6204-9948>  
O Kfir  <https://orcid.org/0000-0003-1253-9372>  
W Hu  <https://orcid.org/0000-0003-4936-1483>  
S Zayko  <https://orcid.org/0000-0001-9826-8627>  
S Eisebitt  <https://orcid.org/0000-0001-7608-5061>  
B Pfau  <https://orcid.org/0000-0001-9057-0346>  
R Frömter  <https://orcid.org/0000-0002-3399-5628>  
M Kläui  <https://orcid.org/0000-0002-4848-2569>  
F S Yasin  <https://orcid.org/0000-0001-9382-7565>  
B J McMorran  <https://orcid.org/0000-0001-7207-1076>  
S Seki  <https://orcid.org/0000-0001-6094-8643>  
X Yu  <https://orcid.org/0000-0003-3136-7289>  
A Lubk  <https://orcid.org/0000-0003-2698-8806>  
D Wolf  <https://orcid.org/0000-0001-5000-8578>  
N Pryds  <https://orcid.org/0000-0002-5718-7924>  
D Makarov  <https://orcid.org/0000-0002-7177-4308>  
M Poggio  <https://orcid.org/0000-0002-5327-051X>

## References

- [1] Clarke J and Braginski A I 2004 *The SQUID Handbook: Fundamentals and Technology of SQUIDs and SQUID Systems* vol 1 (WILEY-VCH Verlag GmbH & Co. KGaA)
- [2] Persky E, Sochnikov I and Kalisky B 2022 Studying quantum materials with scanning SQUID microscopy *Annu. Rev. Condens. Matter Phys.* **13** 385–405
- [3] José Martínez-Pérez M and Koelle D 2017 NanoSQUIDs: basics & recent advances *Phys. Sci. Rev.* **2** 20175001
- [4] Huber M E, Koshnick N C, Bluhm H, Archuleta L J, Azua T, Björnsson P G, Gardner B W, Halloran S T, Lucero E A and Moler K A 2008 Gradiometric micro-SQUID susceptometer for scanning measurements of mesoscopic samples *Rev. Sci. Instrum.* **79** 053704
- [5] Finkler A, Segev Y, Myasoedov Y, Rappaport M L, Ne'eman L, Vasyukov D, Zeldov E, Huber M E, Martin J and Yacoby A 2010 Self-aligned nanoscale SQUID on a tip *Nano Lett.* **10** 1046–9
- [6] Vasyukov D *et al* 2013 A scanning superconducting quantum interference device with single electron spin sensitivity *Nat. Nanotechnol.* **8** 639–44
- [7] Wyss M, Bagani K, Jetter D, Marchiori E, Vervelaki A, Gross B, Ridderbos J, Gliga S, Schönenberger C and Poggio M 2022 Magnetic, thermal, and topographic imaging with a nanometer-scale SQUID-on-lever scanning probe *Phys. Rev. Appl.* **17** 034002
- [8] Kirtley J R 2010 Fundamental studies of superconductors using scanning magnetic imaging *Rep. Prog. Phys.* **73** 126501
- [9] Persky E, Vardi N, Shperber Y and Kalisky B 2018 Improving the sensitivity of scanning probe microscopy with mechanical vibrations *Appl. Phys. Lett.* **113** 173101
- [10] Foroughi F, Mol J-M, Müller T, Kirtley J R, Moler K A and Bluhm H 2018 A micro-SQUID with dispersive readout for magnetic scanning microscopy *Appl. Phys. Lett.* **112** 252601
- [11] Cui Z, Kirtley J R, Wang Y, Kratz P A, Rosenberg A J, Watson C A, Gibson G W, Ketchen M B and Moler K A 2017 Scanning SQUID sampler with 40-ps time resolution *Rev. Sci. Instrum.* **88** 083703
- [12] Halbertal D *et al* 2016 Nanoscale thermal imaging of dissipation in quantum systems *Nature* **539** 407–10
- [13] Griesmar J, Rodriguez R H, Benzoni V, Pillet J-D, Smirr J-L, Lafont F and Girit Ç Ö 2021 Superconducting on-chip spectrometer for mesoscopic quantum systems *Phys. Rev. Res.* **3** 043078
- [14] Müller B, Karrer M, Limberger F, Becker M, Schröppel B, Burkhardt C J, Kleiner R, Goldobin E and Koelle D 2019 Josephson junctions and SQUIDs created by focused helium-ion-beam irradiation of  $\text{YBa}_2\text{Cu}_3\text{O}_7$  *Phys. Rev. Appl.* **11** 044082
- [15] Bert J A, Kalisky B, Bell C, Kim M, Hikita Y, Hwang H Y and Moler K A 2011 Direct imaging of the coexistence of ferromagnetism and superconductivity at the  $\text{LaAlO}_3/\text{SrTiO}_3$  interface *Nat. Phys.* **7** 767–71
- [16] Christensen D V *et al* 2019 Strain-tunable magnetism at oxide domain walls *Nat. Phys.* **15** 269–74
- [17] Wang X R *et al* 2015 Imaging and control of ferromagnetism in  $\text{LaMnO}_3/\text{SrTiO}_3$  heterostructures *Science* **349** 716–9
- [18] Lachman E O *et al* 2015 Visualization of superparamagnetic dynamics in magnetic topological insulators *Sci. Adv.* **1** e1500740
- [19] Tschirhart C L *et al* 2021 Imaging orbital ferromagnetism in a moiré Chern insulator *Science* **372** 1323–7
- [20] Wang Z-C *et al* 2021 Colossal magnetoresistance without mixed valence in a layered phosphide crystal *Adv. Mater.* **33** 2005755
- [21] Frenkel Y, Haham N, Shperber Y, Bell C, Xie Y, Chen Z, Hikita Y, Hwang H Y, Salje E K H and Kalisky B 2017 Imaging and tuning polarity at  $\text{SrTiO}_3$  domain walls *Nat. Mater.* **16** 1203–8
- [22] Tsuei C C and Kirtley J R 2000 Pairing symmetry in cuprate superconductors *Rev. Mod. Phys.* **72** 969–1016
- [23] Kirtley J R, Tsuei C C, Ariando A, Verwijs C J M, Harkema S and Hilgenkamp H 2006 Angle-resolved phase-sensitive determination of the in-plane gap symmetry in  $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$  *Nat. Phys.* **2** 190–4
- [24] Embon L *et al* 2017 Imaging of super-fast dynamics and flow instabilities of superconducting vortices *Nat. Commun.* **8** 82
- [25] Kremen A, Khan H, Loh Y L, Baturina T I, Trivedi N, Frydman A and Kalisky B 2018 Imaging quantum fluctuations near criticality *Nat. Phys.* **14** 1205–10
- [26] Bachmann M D *et al* 2019 Spatial control of heavy-fermion superconductivity in  $\text{CeIrIn}_5$  *Science* **366** 221–6
- [27] Jarjour A, Ferguson G M, Schaefer B T, Lee M, Loh Y L, Trivedi N and Nowack K C 2023 Superfluid response of an atomically thin, gate-tuned van der Waals superconductor *Nat. Commun.* **14** 2055
- [28] Sochnikov I *et al* 2015 Nonsinusoidal current-phase relationship in Josephson junctions from the 3D topological insulator  $\text{HgTe}$  *Phys. Rev. Lett.* **114** 066801
- [29] Spanton E M, Deng M, Vaitiekėnas S, Krogstrup P, Nygård J, Marcus C M and Moler K A 2017 Current-phase relations of few-mode InAs nanowire Josephson junctions *Nat. Phys.* **13** 1177–81
- [30] Persky E, Bjørlig A V, Feldman I, Almoalem A, Altman E, Berg E, Kimchi I, Ruhman J, Kanigel A and Kalisky B 2022 Magnetic memory and spontaneous vortices in a van der Waals superconductor *Nature* **607** 692–6
- [31] Iguchi Y, Shi R A, Kihou K, Lee C-H, Barkman M, Benfenati A L, Grinenko V, Babaev E and Moler K A 2023 Superconducting vortices carrying a temperature-dependent fraction of the flux quantum *Science* **380** 1244–7
- [32] Nowack K C *et al* 2013 Imaging currents in  $\text{HgTe}$  quantum wells in the quantum spin Hall regime *Nat. Mater.* **12** 787–91
- [33] Ferguson G M, Xiao R, Richardella A R, Low D, Samarth N and Nowack K C 2023 Direct visualization of electronic transport in a quantum anomalous Hall insulator *Nat. Mater.* **22** 1100–5
- [34] Kalisky B *et al* 2013 Locally enhanced conductivity due to the tetragonal domain structure in  $\text{LaAlO}_3/\text{SrTiO}_3$  heterointerfaces *Nat. Mater.* **12** 1091–5
- [35] Persky E *et al* 2021 Non-universal current flow near the metal-insulator transition in an oxide interface *Nat. Commun.* **12** 3311
- [36] Aharon-Steinberg A *et al* 2022 Direct observation of vortices in an electron fluid *Nature* **607** 74–80
- [37] Uri A *et al* 2020 Nanoscale imaging of equilibrium quantum Hall edge currents and of the magnetic monopole response in graphene *Nat. Phys.* **16** 164–70
- [38] Uri A *et al* 2020 Mapping the twist-angle disorder and Landau levels in magic-angle graphene *Nature* **581** 47–52
- [39] Wang X, Laav M, Volotsenko I, Frydman A and Kalisky B 2022 Visualizing current in superconducting networks *Phys. Rev. Appl.* **17** 024073
- [40] Anahory Y *et al* 2016 Emergent nanoscale superparamagnetism at oxide interfaces *Nat. Commun.* **7** 12566
- [41] Chernobrod B M and Berman G P 2005 Spin microscope based on optically detected magnetic resonance *J. Appl. Phys.* **97** 014903
- [42] Taylor J M *et al* 2008 High-sensitivity diamond magnetometer with nanoscale resolution *Nat. Phys.* **4** 810–816
- [43] Degen C L 2008 Scanning magnetic field microscope with a diamond single-spin sensor *Appl. Phys. Lett.* **92** 243111
- [44] Steinert S, Dolde F, Neumann P, Aird A, Naydenov B, Balasubramanian G, Jelezko F and Wrachtrup J 2010 High sensitivity magnetic imaging using an array of spins in diamond *Rev. Sci. Instrum.* **81** 043705
- [45] Barry J F *et al* 2020 Sensitivity optimization for NV-diamond magnetometry *Rev. Mod. Phys.* **92** 015004

- [46] Delaney P, Greer J C and Larsson J A 2010 Spin-polarization mechanisms of the nitrogen-vacancy center in diamond *Nano Lett.* **10** 610–4
- [47] Lekavicius I *et al* 2022 Orders of magnitude improvement in coherence of silicon-vacancy ensembles in isotopically purified 4H-SiC *PRX Quantum* **3** 010343
- [48] Osterkamp C, Mangold M, Lang J, Balasubramanian P, Teraji T, Naydenov B and Jelezko F 2019 Engineering preferentially-aligned nitrogen-vacancy centre ensembles in CVD grown diamond *Sci. Rep.* **9** 5786
- [49] Rembold P, Oshnik N, Müller M M, Montangero S, Calarco T and Neu E 2020 Introduction to quantum optimal control for quantum sensing with nitrogen-vacancy centers in diamond *AVS Quantum Sci.* **2** 024701
- [50] Hart C A, Schloss J M, Turner M J, Scheidegger P J, Bauch E and Walsworth R L 2021 N-V–diamond magnetic microscopy using a double quantum 4-Ramsey protocol *Phys. Rev. Appl.* **15** 044020
- [51] Webb J L *et al* 2022 High-speed wide-field imaging of microcircuitry using nitrogen vacancies in Diamond *Phys. Rev. Appl.* **17** 064051
- [52] Wolf T *et al* 2015 Subpicotesla diamond magnetometry *Phys. Rev. X* **5** 041001
- [53] Stürner F M, Liu Y, Colard P-O, Markham M and Jelezko F 2021 Magnetometry based on the excited-state lifetimes of a single nitrogen-vacancy center in diamond *Appl. Phys. Lett.* **119** 134001
- [54] Bertoldo F, Ali S, Manti S and Thygesen K S 2022 Quantum point defects in 2D materials—the QPOD database *npj Comput. Mater.* **8** 56
- [55] Haryuyama M, Okigawa Y, Okada M, Nakajima H, Okazaki T, Kato H, Makino T and Yamada T 2023 Charge stabilization of shallow nitrogen-vacancy centers using graphene/diamond junctions *Appl. Phys. Lett.* **122** 141601
- [56] Kehayias P *et al* 2022 Measurement and simulation of the magnetic fields from a 555 timer integrated circuit using a quantum diamond microscope and finite-element analysis *Phys. Rev. Appl.* **17** 014021
- [57] Tétienne J-P, Dontschuk N, Broadway D A, Stacey A, Simpson D A and Hollenberg L C L 2017 Quantum imaging of current flow in graphene *Sci. Adv.* **3** e1602429
- [58] Ku M J H *et al* 2020 Imaging viscous flow of the Dirac fluid in graphene *Nature* **583** 537–41
- [59] Zhou B B *et al* 2020 Spatiotemporal mapping of a photocurrent vortex in monolayer MoS<sub>2</sub> using diamond quantum sensors *Phys. Rev. X* **10** 011003
- [60] McLaughlin N J *et al* 2021 Strong correlation between superconductivity and ferromagnetism in an Fe-chalcogenide superconductor *Nano Lett.* **17** 7277–83
- [61] Bertelli I, Simon B G, Yu T, Aarts J, Bauer G E W, Blanter Y M and van der Sar T 2021 Imaging spin-wave damping underneath metals using electron spins in diamond *Adv. Quantum Technol.* **4** 2100094
- [62] Scheidegger P J, Diesch S, Palm M L and Degen C L 2022 Scanning nitrogen-vacancy magnetometry down to 350 mK *Appl. Phys. Lett.* **120** 224001
- [63] Hall L T *et al* 2012 High spatial and temporal resolution wide-field imaging of neuron activity using quantum NV-diamond *Sci. Rep.* **2** 401
- [64] Sengottuvel S, Mrózek M, Sawczak M, Głowacki M J, Ficek M, Gawlik W and Wojciechowski A M 2022 Wide-field magnetometry using nitrogen-vacancy color centers with randomly oriented micro-diamonds *Sci. Rep.* **12** 17997
- [65] Liu W *et al* 2022 Silicon-vacancy nanodiamonds as high performance near-infrared emitters for live-cell dual-color imaging and thermometry *Nano Lett.* **22** 2881–8
- [66] Maletinsky P, Hong S, Grinolds M S, Hausmann B, Lukin M D, Walsworth R L, Loncar M and Yacoby A 2012 A robust scanning diamond sensor for nanoscale imaging with single nitrogen-vacancy centres *Nat. Nanotechnol.* **7** 320–4
- [67] Simpson D A, Tétienne J P, McCoe J M, Ganesan K, Hall L T, Petrou S, Scholten R E and Hollenberg L C L 2016 Magneto-optical imaging of thin magnetic films using spins in diamond *Sci. Rep.* **6** 22797
- [68] Lesik M *et al* 2019 Magnetic measurements on micrometer-sized samples under high pressure using designed NV centers *Science* **366** 1359–62
- [69] Liu G-Q, Feng X, Wang N, Li Q and Liu R 2019 Coherent quantum control of nitrogen-vacancy center spins near 1000 kelvin *Nat. Commun.* **10** 1344
- [70] Broadway D, Lillie S, Scholten S, Rohner D, Dontschuk N, Maletinsky P, Tétienne J-P and Hollenberg L 2020 Improved current density and magnetization reconstruction through vector magnetic field measurements *Phys. Rev. Appl.* **14** 024076
- [71] Dubois A E E, Broadway D A, Stark A, Tschudin M A, Healey A J, Huber S D, Tétienne J-P, Greplova E and Maletinsky P 2022 Untrained physically informed neural network for image reconstruction of magnetic field sources *Phys. Rev. Appl.* **18** 064076
- [72] Chang K, Eichler A, Rhensius J, Lorenzelli L and Degen C L 2017 Nanoscale imaging of current density with a single-spin magnetometer *Nano Lett.* **17** 2367–73
- [73] Scholten S C, Healey A J, Robertson I O, Abrahams G J, Broadway D A and Tétienne J-P 2021 Widefield quantum microscopy with nitrogen-vacancy centers in diamond: strengths, limitations, and prospects *J. Appl. Phys.* **130** 150902
- [74] Batzer M, Shields B, Neu E, Widmann C, Giese C, Nebel C and Maletinsky P 2020 Single crystal diamond pyramids for applications in nanoscale quantum sensing *Opt. Mater. Express* **10** 492–500
- [75] Chrostoski P, Sadeghpour H R and Santamore D H 2018 Electric noise spectra of a near-surface nitrogen-vacancy center in diamond with a protective layer *Phys. Rev. Appl.* **10** 064056
- [76] Abendroth J M, Herb K, Janitz E, Zhu T, Völker L A and Degen C L 2022 Single-nitrogen–vacancy NMR of amine-functionalized diamond surfaces *Nano Lett.* **22** 7294
- [77] Zhang Q *et al* 2021 High-fidelity single-shot readout of single electron spin in diamond with spin-to-charge conversion *Nat. Commun.* **12** 1529
- [78] Irber D M, Poggiali F, Kong F, Kieschnick M, Lüthmann T, Kwiatkowski D, Meijer J, Du J, Shi F and Reinhard F 2021 Robust all-optical single-shot readout of nitrogen-vacancy centers in diamond *Nat. Commun.* **12** 532
- [79] Hrubesch F M, Braunbeck G, Stutzmann M, Reinhard F and Brandt M S 2017 *Phys. Rev. Lett.* **118** 037601
- [80] Siyushev P, Nesladek M, Bourgeois E, Gulka M, Hruba J, Yamamoto T, Trupke M, Teraji T, Isoya J and Jelezko F 2019 *Science* **363** 728
- [81] Hopper D A, Lauigan J D, Huang T-Y and Bassett L C 2020 Real-time charge initialization of diamond nitrogen-vacancy centers for enhanced spin readout *Phys. Rev. Appl.* **13** 024016
- [82] Zhang J *et al* 2023 Efficient real-time spin readout of nitrogen-vacancy centers based on Bayesian estimation (arXiv:2302.06310)
- [83] Craigie K, Gauger E M, Altmann Y and Bonato C 2021 Resource-efficient adaptive Bayesian tracking of magnetic fields with a quantum sensor *J. Phys.: Condens. Matter* **33** 195801

- [84] Rollo M, Finco A, Tanos R, Fabre F, Devolder T, Robert-Philip I and Jacques V 2021 Quantitative study of the response of a single NV defect in diamond to magnetic noise *Phys. Rev. B* **103** 235418
- [85] Lin S *et al* 2022 Diamond surface electric-field noise detection using shallow nitrogen-vacancy centers *Phys. Rev. B* **106** 165406
- [86] Thiel L, Wang Z, Tschudin M A, Rohner D, Gutiérrez-Lezama I, Ubrig N, Gibertini M, Giannini E, Morpurgo A F and Maletinsky P 2019 Probing magnetism in 2D materials at the nanoscale with single-spin microscopy *Science* **364** 973–6
- [87] Gross I *et al* 2017 Real-space imaging of non-collinear antiferromagnetic order with a single-spin magnetometer *Nature* **549** 252–6
- [88] Appel P, Shields B J, Kosub T, Hedrich N, Hübner R, Faßbender J, Makarov D and Maletinsky P 2019 Nanomagnetism of magnetoelectric granular thin-film antiferromagnets *Nano Lett.* **19** 1682–7
- [89] Hedrich N, Wagner K, Pylypovskiy O V, Shields B J, Kosub T, Sheka D D, Makarov D and Maletinsky P 2021 Nanoscale mechanics of antiferromagnetic domain walls *Nat. Phys.* **17** 574–7
- [90] Gaurav Rana K *et al* 2020 Olivier Boule room-temperature skyrmions at zero field in exchange-biased ultrathin films *Phys. Rev. Appl.* **13** 044079
- [91] Tétienne J-P *et al* 2015 The nature of domain walls in ultrathin ferromagnets revealed by scanning nanomagnetometry *Nat. Commun.* **6** 6733
- [92] Thiel L, Rohner D, Ganzhorn M, Appel P, Neu E, Mueller B, Kleiner R, Koelle D and Maletinsky P 2016 Quantitative nanoscale vortex imaging using a cryogenic quantum magnetometer *Nat. Nanotechnol.* **11** 677–81
- [93] Finco A *et al* 2022 Imaging topological defects in a noncollinear antiferromagnet *Phys. Rev. Lett.* **128** 187201
- [94] Rondin L, Tétienne J-P, Hingant T, Roch J-F, Maletinsky P and Jacques V 2014 Magnetometry with nitrogen-vacancy defects in diamond *Rep. Prog. Phys.* **77** 056503
- [95] Du C *et al* 2017 Control and local measurement of the spin chemical potential in a magnetic insulator *Science* **357** 195–8
- [96] Flebus B and Tserkovnyak Y 2018 Quantum-impurity relaxometry of magnetization dynamics *Phys. Rev. Lett.* **121** 187204
- [97] Zhou T X *et al* 2021 A magnon scattering platform *Proc. Natl Acad. Sci.* **118** e2019473118
- [98] Andrich P, de Las Casas C F, Liu X, Bretscher H L, Berman J R, Heremans F J, Nealey P F and Awschalom D D 2017 Long-range spin wave mediated control of defect qubits in nanodiamonds *npj Quantum Inf.* **3** 28
- [99] Bertelli I, Carmiggelt J J, Yu T, Simon B G, Pothoven C C, Bauer G E W, Blanter Y M, Aarts J and van der Sar T 2020 Magnetic resonance imaging of spin-wave transport and interference in a magnetic insulator *Sci. Adv.* **6** eabd3556
- [100] Simon B G, Kurdi S, Carmiggelt J J, Borst M, Katan A J and van der Sar T 2022 Filtering and imaging of frequency-degenerate spin waves using nanopositioning of a single-spin sensor *Nano Lett.* **22** 9198–204
- [101] Wang G, Liu Y-X, Schloss J M, Alsid S T, Braje D A and Cappellaro P 2022 Sensing of arbitrary-frequency fields using a quantum mixer *Phys. Rev. X* **12** 021061
- [102] Carmiggelt J J, Bertelli I, Mulder R W, Teepe A, Elyasi M, Simon B G, Bauer G E W, Blanter Y M and van der Sar T 2023 Broadband microwave detection using electron spins in a hybrid diamond-magnet sensor chip *Nat. Commun.* **14** 490
- [103] Wolfe C S, Bhallamudi V P, Wang H L, Du C H, Manuilov S, Teeling-Smith R M, Berger A J, Adur R, Yang F Y and Hammel P C 2014 Off-resonant manipulation of spins in diamond via precessing magnetization of a proximal ferromagnet *Phys. Rev. B* **89** 180406
- [104] Koerner C, Dreyer R, Wagener M, Liebing N, Bauer H G and Woltersdorf G 2022 Frequency multiplication by collective nanoscale spin-wave dynamics *Science* **375** 1165–9
- [105] McCullian B A, Thabt A M, Gray B A, Melendez A L, Wolf M S, Safonov V L, Pelekhov D V, Bhallamudi V P, Page M R and Hammel P C 2020 Broadband multi-magnon relaxometry using a quantum spin sensor for high frequency ferromagnetic dynamics sensing *Nat. Commun.* **11** 5229
- [106] Wang H *et al* 2022 Noninvasive measurements of spin transport properties of an antiferromagnetic insulator *Sci. Adv.* **8** eabg8562
- [107] Finco A *et al* 2021 Imaging non-collinear antiferromagnetic textures via single spin relaxometry *Nat. Commun.* **12** 767
- [108] McLaughlin N J, Hu C, Huang M, Zhang S, Lu H, Yan G Q, Wang H, Tserkovnyak Y, Ni N and Du C R 2022 Quantum imaging of magnetic phase transitions and spin fluctuations in intrinsic magnetic topological nanoflakes *Nano Lett.* **22** 5810–7
- [109] Huang M *et al* 2022 Wide field imaging of van der Waals ferromagnet Fe<sub>3</sub>GeTe<sub>2</sub> by spin defects in hexagonal boron nitride *Nat. Commun.* **13** 5369
- [110] Chatterjee S, Rodriguez-Nieva J F and Demler E 2019 Diagnosing phases of magnetic insulators via noise magnetometry with spin qubits *Phys. Rev. B* **99** 104425
- [111] Song T *et al* 2021 Direct visualization of magnetic domains and moiré magnetism in twisted 2D magnets *Science* **374** 1140–4
- [112] Pelliccione M, Jenkins A, Overtchayapong P, Reetz C, Emmanouilidou E, Ni N and Bleszynski Jayich A C 2016 Scanned probe imaging of nanoscale magnetism at cryogenic temperatures with a single-spin quantum sensor *Nat. Nanotechnol.* **11** 700–5
- [113] Dingle R, Stormer H L, Gossard A C and Wiegmann W 1978 Electron mobilities in modulation-doped semiconductor heterojunction superlattices *Appl. Phys. Lett.* **33** 665–7
- [114] Chang A M, Hallen H D, Harriott L, Hess H F, Kao H L, Kwo J, Miller R E, Wolfe R, van der Ziel J and Chang T Y 1992 Scanning Hall probe microscopy *Appl. Phys. Lett.* **61** 1974–6
- [115] Oral A, Bending S J and Henini M 1996 Real-time scanning Hall probe microscopy *Appl. Phys. Lett.* **69** 1324–6
- [116] Vervaeke K, Simoen E, Borghs H and Moshchalkov V V 2009 Size dependence of microscopic Hall sensor detection limits *Rev. Sci. Instrum.* **80** 074701
- [117] Sandhu A, Kurosawa K, Dede M and Oral A 2004 50 nm Hall sensors for room temperature scanning Hall probe microscopy *Jpn. J. Appl. Phys.* **43** 777–8
- [118] Sandhu A, Sanbonsugi H, Shibasaki I, Abe M and Handa H 2004 High sensitivity InSb ultra-thin film micro-hall sensors for bioscreening applications *Jpn. J. Appl. Phys.* **43** L868–L870
- [119] Collomb D, Li P and Bending S 2021 Frontiers of graphene-based Hall-effect sensors *J. Phys.: Condens. Matter* **33** 243002
- [120] Sonusen S, Karci O, Dede M, Aksoy S and Oral A 2014 Single layer graphene hall sensors for scanning Hall probe microscopy (SHPM) in 3–300 K temperature range *Appl. Surf. Sci.* **308** 414–8
- [121] Collomb D, Li P and Bending S 2019 Nanoscale graphene Hall sensors for high-resolution ambient magnetic imaging *Sci. Rep.* **9** 14424
- [122] Collomb D, Zhang M, Yuan W and Bending S J 2021 Imaging of strong nanoscale vortex pinning in GdBaCuO high-temperature superconducting tapes *Nanomaterials* **11** 1082
- [123] Hicks C W, Luan L, Moler K A, Zeldov E and Shtrikman H 2007 Noise characteristics of 100nm scale GaAs/Al<sub>x</sub>Ga<sub>1-x</sub>As scanning Hall probes *Appl. Phys. Lett.* **90** 133512



- [124] Orr J M S, Gilbertson A M, Fearn M, Croad O W, Storey C J, Buckle L, Emeny M T, Buckle P D and Ashley T 2008 Electronic transport in modulation-doped InSb quantum well heterostructures *Phys. Rev. B* **77** 165334
- [125] Behet M, Bekaert J, De Boeck J and Borghs G 2000 InAs/Al<sub>0.2</sub>Ga<sub>0.8</sub>Sb quantum well Hall effect sensors *Sens. Actuators* **81** 13–17
- [126] Kazakova O, Gallop J C, Cox D C, Brown E, Cuenat A and Suzuki K 2008 Optimization of 2DEG InAs/GaSb Hall sensors for single particle detection *IEEE Trans. Magn.* **44** 4480–3
- [127] Bando M, Ohashi T, Dede M, Akram R, Oral A, Park S Y, Shibasaki I, Handa H and Sandhu A 2009 High sensitivity and multifunctional micro-Hall sensors fabricated using InAlSb/InAsSb/InAlSb heterostructures *J. Appl. Phys.* **105** 07E909
- [128] Pross A, Crisan A I, Bending S J, Mosser V and Konczykowski M 2005 Second-generation quantum-well sensors for room-temperature scanning Hall probe microscopy *J. Appl. Phys.* **97** 096105
- [129] Chenaud B *et al* 2016 Sensitivity and noise of micro-Hall magnetic sensors based on InGaAs quantum wells *J. Appl. Phys.* **119** 024501
- [130] Tang C C, Li M Y, Li L J, Chi C C and Chen J C 2011 Characteristics of a sensitive micro-Hall probe fabricated on chemical vapor deposited graphene over the temperature range from liquid-helium to room temperature *Appl. Phys. Lett.* **99** 112107
- [131] Xu H, Zhang Z, Shi R, Liu H, Wang Z, Wang S and Peng L-M 2013 Batch-fabricated high-performance graphene Hall elements *Sci. Rep.* **3** 1207
- [132] Xu H, Huang L, Zhang Z, Chen B, Zhong H and Peng L-M 2013 Flicker noise and magnetic resolution of graphene Hall sensors at low frequency *Appl. Phys. Lett.* **103** 112405
- [133] Huang L, Zhang Z, Chen B, Ma X, Zhong H and Peng L-M 2014 Ultra-sensitive graphene Hall elements *Appl. Phys. Lett.* **104** 183106
- [134] Panchal V, Cedergren K, Yakimova R, Tzalenchuk A, Kubatkin S and Kazakova O 2012 Small epitaxial graphene devices for magnetosensing applications *J. Appl. Phys.* **111** 07E509
- [135] Novoselov K S, Jiang Z, Zhang Y, Morozov S V, Stormer H L, Zeitler U, Maan J C, Boebinger G S, Kim P and Geim A K 2007 Room-temperature quantum Hall effect in graphene *Science* **315** 1379
- [136] Dai T, Xu H, Chen S and Zhang Z 2022 High performance Hall sensors built on chemical vapor deposition-grown bilayer graphene *ACS Omega* **7** 25644–9
- [137] Mayorov A S *et al* 2011 Micrometer-scale ballistic transport in encapsulated graphene at room temperature *Nano Lett.* **11** 2396–9
- [138] Schaefer B T, Wang L, Jarjour A, Watanabe K, Taniguchi T, McEuen P L and Nowack K C 2020 Magnetic field detection limits for ultraclean graphene Hall sensors *Nat. Commun.* **11** 4163
- [139] Murali R, Yang Y, Brenner K, Beck T and Meindl J D 2009 Breakdown current density of graphene nanoribbons *Appl. Phys. Lett.* **94** 243114
- [140] Asad M, Jeppson K O, Vorobiev A, Bonmann M and Stake J 2021 Enhanced high-frequency performance of top-gated graphene FETs due to substrate-induced improvements in charge carrier saturation velocity *IEEE Trans. Electron Devices* **68** 899–902
- [141] Brook A J, Bending S J, Pinto J, Oral A, Ritchie D, Beere H, Springthorpe A and Henini M 2003 Micromachined III–V cantilevers for AFM-tracking scanning Hall probe microscopy *J. Micromech. Microeng.* **13** 124–8
- [142] Schweinböck T, Weiss D, Lipinski M and Eberl K 2000 Scanning Hall probe microscopy with shear force distance control *J. Appl. Phys.* **87** 6496–8
- [143] Dede M, Ürkmen K, Girişen Ö, Atabak M, Oral A, Farrer I and Ritchie D 2008 Scanning Hall probe microscopy (SHPM) using quartz crystal AFM feedback *J. Nanosci. Nanotechnol.* **8** 619–22
- [144] Chong B K, Zhou H, Mills G, Donaldson L and Weaver J M R 2001 Scanning Hall probe microscopy on an atomic force microscope tip *J. Vac. Sci. Technol. A* **19** 1769–72
- [145] Oral A, Kaval A, Dede M, Masuda H, Okamoto A, Shibasaki I and Sandhu A 2002 Room-temperature scanning Hall probe microscope (RT-SHPM) imaging of garnet films using new high-performance InSb sensors *IEEE Trans. Magn.* **38** 2438–40
- [146] Dauber J, Sagade A A, Oellers M, Watanabe K, Taniguchi T, Neumaier D and Stampfer C 2015 Ultra-sensitive Hall sensors based on graphene encapsulated in hexagonal boron nitride *Appl. Phys. Lett.* **106** 193501
- [147] Holleis S, Anna Thomas A, Shipulin I A, Hühne R, Steiger-Thirfeld A, Bernardi J and Eisterer M 2022 Magnetic granularity in PLD-grown Fe(Se,Te) films on simple RABiTS templates *Supercond. Sci. Technol.* **35** 074001
- [148] Li Y, Chen S, Paidpilli M, Jain R, Goel C and Selvamannickam V 2022 A reel-to-reel scanning Hall probe microscope for characterizing long REBCO conductor in magnetic fields up to 5 Tesla *IEEE Trans. Appl. Supercond.* **32** 9000206
- [149] Sandra J S, Paidpilli M, Sun S, Suarez-Villagran M, Yerraguravagari V, Wosik J and Selvamannickam V 2022 Applications-oriented development of buffer architecture for REBCO films on nonmetallic substrates *IEEE Trans. Appl. Supercond.* **33** 6602506
- [150] Fert A, Cros V and Sampaio J 2013 Skyrmions on the track *Nat. Nanotechnol.* **8** 152–6
- [151] Wang L *et al* 2018 Ferroelectrically tunable magnetic skyrmions in ultrathin oxide heterostructures *Nat. Mater.* **17** 1087–94
- [152] Panchal V, Iglesias-Freire Ó, Lartsev A, Yakimova R, Asenjo A and Kazakova O 2013 Magnetic scanning probe calibration using graphene Hall sensor *IEEE Trans. Magn.* **49** 3520–3
- [153] Pereira J M B, Pacheco C J, Arenas M P, Araujo J F D F, Pereira G R and Bruno A C 2017 Novel scanning dc-susceptometer for characterization of heat-resistant steels with different states of aging *J. Magn. Magn. Mater.* **442** 311–8
- [154] Meyer E, Bennowitz R and Hug H J 2021 Magnetic force microscopy *Scanning Probe Microscopy. The Lab on a Tip* 2nd edn, ed E Meyer, R Bennowitz and H J Hug (Springer) pp 109–211
- [155] Jaafar M, Iglesias-Freire Ó, Serrano-Ramon L, Ibarra M R, de Teresa J M and Asenjo A 2011 Distinguishing magnetic and electrostatic interactions by a Kelvin probe force microscopy-magnetic force microscopy combination *Beilstein J. Nanotechnol.* **2** 552–60
- [156] Meng K-Y *et al* 2019 Observation of nanoscale skyrmions in SrIrO<sub>3</sub>/SrRuO<sub>3</sub> bilayers *Nano Lett.* **19** 3169–75
- [157] Dietz C, Herruzo E T, Lozano J R and Garcia R 2011 Nanomechanical coupling enables detection and imaging of 5 nm superparamagnetic particles in liquid *Nanotechnology* **22** 125708
- [158] Deussner S, Suess D, Abert C, Bruckner F, Fähler S, Heistracher P, Reichel L and Neu V 2022 Rigorous single-period micromagnetic model of stripe domains: comparison with analytics and experiment *Phys. Rev. B* **106** 064404
- [159] Feng Y, Vaghefi P M, Vranjkovic S, Penedo M, Kappenberger P, Schwenk J, Mandra A-O and Hug H J 2022 Magnetic force microscopy contrast formation and field sensitivity *J. Magn. Magn. Mater.* **551** 169073
- [160] Zhao X, Schwenk J, Mandra A-O, Penedo M, Bacani M and Hug H J 2018 Magnetic force microscopy with frequency-modulated capacitive tip-sample distance control *New J. Phys.* **20** 013018
- [161] Hu X *et al* 2020 Round robin comparison on quantitative nanometer scale magnetic field measurements by magnetic force microscopy *J. Magn. Magn. Mater.* **511** 166947



- [162] Feng Y, Mandru A-O, Yıldırım O and Hug H J 2022 Quantitative magnetic force microscopy: transfer-function method revisited *Phys. Rev. Appl.* **18** 024016
- [163] Sakar B, Liu Y, Sievers S, Neu V, Lang J, Osterkamp C, Markham M L, Öztürk O, Jelezko F and Schumacher H W 2021 Quantum calibrated magnetic force microscopy *Phys. Rev. B* **104** 214427
- [164] Cao Y, Zhao Y, Kumar P, Yoshimura S and Saito H 2018 Magnetic domain imaging of a very rough fractured surface of Sr ferrite magnet without topographic crosstalk by alternating magnetic force microscopy with a sensitive FeCo-GdO<sub>x</sub> superparamagnetic tip *J. Appl. Phys.* **27** 224503
- [165] Freitag N, Reiche C F, Neu V, Devi P, Burkhardt U, Felser C, Wolf D, Lubk A, Büchner B and Mühl T 2023 Simultaneous magnetic field and field gradient mapping of hexagonal MnNiGa by quantitative magnetic force microscopy *Commun. Phys.* **6** 11
- [166] Li Z, Lu W, Song Y, Wang Y, Chen A, Yan B, Yoshimura S and Saito H 2016 Quantitatively probing the magnetic behavior of individual nanoparticles by an AC field-modulated magnetic force microscopy *Sci. Rep.* **6** 22467
- [167] Makarova M V, Suzuki K, Kon H, Dubey P, Sonobe H, Matsumura T and Saito H 2023 High-frequency magnetic field energy imaging of magnetic recording head by alternating magnetic force microscopy (A-MFM) with superparamagnetic tip *Appl. Sci.* **13** 4843
- [168] Singh B, Ravishankar R, Otálora J A, Soldatov I, Schäfer R, Karnausenko D, Neu V and Schmidt O G 2022 Direct imaging of field-driven nanoscale domain wall oscillations in Landau structures *Nanoscale* **14** 13667
- [169] Rugar D, Mamin H J, Guethner P, Lambert S E, Stern J E, McFadyen I and Yogi T 1990 Magnetic force microscopy: general principles and application to longitudinal recording media *J. Appl. Phys.* **68** 1169–16
- [170] Nozaki Y, Ohta M, Taharazako S, Tateishi K, Yoshimura S and Matsuyama K 2007 Magnetic force microscopy study of microwave-assisted magnetization reversal in submicron-scale ferromagnetic particles *Appl. Phys. Lett.* **91** 082510
- [171] Parkin S S P, Hayashi M and Thomas L 2008 Magnetic domain-wall racetrack memory *Science* **320** 190–4
- [172] Soumyanarayanan A *et al* 2017 Tunable room-temperature magnetic skyrmions in Ir/Fe/Co/Pt multilayers *Nat. Mater.* **16** 898–904
- [173] Chesnel K *et al* 2018 Morphological stripe-bubble transition in remanent magnetic domain patterns of Co/Pt multilayer films and its dependence on Co thickness *Phys. Rev. B* **98** 224404
- [174] Shinjo T, Okuno T, Hassdorf R, Shigeto K and Ono T 2000 Magnetic vortex core observation in circular dots of permalloy *Science* **289** 930
- [175] May A, Saccone M and van den Berg A 2021 Magnetic charge propagation upon a 3D artificial spin-ice *Nat. Commun.* **12** 3217
- [176] Banuazizi S A H, Houshang A, Awad A A, Mohammadi J, Åkerman J and Belova L M 2022 Magnetic force microscopy of an operational spin nano-oscillator *Microsyst. Nanoeng.* **8** 65
- [177] Schmid I, Marioni M A, Kappenberger P, Romer S, Parlinska-Wojtan M, Hug H J, Hellwig O, Carey M J and Fullerton E E 2010 Exchange bias and domain evolution at 10 nm scales *Phys. Rev. Lett.* **105** 197201
- [178] Heigl M, Vogler C, Mandru A O, Zhao X, Hug H J, Suess D and Albrecht M 2020 Microscopic origin of magnetization reversal in nanoscale exchange-coupled ferri/ferromagnetic bilayers: implications for high energy density permanent magnets and spintronic devices *ACS Appl. Nano Mater.* **3** 9218–25
- [179] Mandru A-O, Yıldırım O, Tomasello R, Heistracher P, Penedo M, Giordano A, Suess D, Finocchio G and Hug H-J 2020 Coexistence of distinct skyrmion phases observed in hybrid ferromagnetic/ferrimagnetic multilayers *Nat. Commun.* **11** 6365
- [180] Wu Y *et al* 2022 Van der Waals interface hosting two groups of magnetic skyrmions *Adv. Mater.* **34** 2110583
- [181] Weber R, Mertin W and Kubalek E 2000 Voltage-influence of biased interconnection line on integrated circuit-internal current contrast measurements via magnetic force microscopy *Microelectron. Reliab.* **40** 1389–94
- [182] Nabaei V, Rajkumar R K, Manzin A, Kazakova O and Tzalenchuk A 2013 Optimization of Hall bar response to localized magnetic and electric fields *J. Appl. Phys.* **14** 064504
- [183] Geng Y, Das H, Wysocki A L, Wang X, Cheong S-W, Mostovoy M, Fennie C J and Wu W 2014 Direct visualization of magnetoelectric domains *Nat. Mater.* **13** 163–7
- [184] Geng W-R *et al* 2022 Magneto–electric–optical coupling in multiferroic BiFeO<sub>3</sub>-based films *Adv. Mater.* **34** 2106396
- [185] Casiraghi A, Corte-León H, Vafaei M, Garcia-Sanchez F, Durin G, Pasquale M, Jakob G, Kläui M and Kazakova O 2019 Individual skyrmion manipulation by local magnetic field gradients *Commun. Phys.* **2** 145
- [186] Sievers S, Braun K-F, Eberbeck D, Gustafsson S, Olsson E, Schumacher H W and Siegner U 2012 Quantitative measurement of the magnetic moment of individual magnetic nanoparticles by magnetic force microscopy *Small* **8** 2675–9
- [187] Stenning K D, Gartside J C, Dion T, Vanstone A, Arroo D M and Branford W R 2021 Magnonic bending, phase shifting and interferometry in a 2D reconfigurable nanodisk crystal *ACS Nano* **15** 674
- [188] Kaiser U, Schwarz A and Wiesendanger R 2007 Magnetic exchange force microscopy with atomic resolution *Nature* **446** 522
- [189] Hauptmann N, Haldar S, Hung T-C, Jolie W, Gutzeit M, Wegner D, Heinze S and Khajetoorians A A 2020 Quantifying exchange forces of a spin spiral on the atomic scale *Nat. Commun.* **11** 1197
- [190] Lepadatu S, Wessely O, Vanhaverbeke A, Allenspach R, Potenza A, Marchetto H, Charlton T R, Langridge S, Dhessi S S and Marrows C H 2010 Domain-wall spin-torque resonators for frequency-selective operation *Phys. Rev. B* **81** 060402
- [191] Kikuchi H and Sumida C 2021 The effect of dc bias current on domain wall resonance property of thin-film magnetoimpedance element *IEEE Trans. Magn.* **57** 4400305
- [192] Marioni M A, Penedo M, Bačani M, Schwenk J and Hug H J 2018 Halbach effect at the nanoscale from chiral spin textures *Nano Lett.* **18** 2263–7
- [193] Zhao X, Mandru A O, Vogler C, Marioni M A, Suess D and Hug H J 2019 Magnetization reversal of strongly exchange-coupled double nanolayers for spintronic devices *ACS Appl. Nano Mater.* **2** 7478–87
- [194] Moser A, Xiao M, Kappenberger P, Takano K, Weresin W, Ikeda Y, Do H and Hug H J 2005 High-resolution magnetic force microscopy study of high-density transitions in perpendicular recording media *J. Magn. Magn. Mater.* **287** 298–302
- [195] Maccariello D, Legrand W, Reyren N, Garcia K, Bouzehouane K, Collin S, Cros V and Fert A 2018 Electrical detection of single magnetic skyrmions in metallic multilayers at room temperature *Nat. Nanotechnol.* **13** 1–6
- [196] Straver E W J, Hoffman J E, Auslaender O M, Rugar D and Moler K A 2008 Controlled manipulation of individual vortices in a superconductor *Appl. Phys. Lett.* **93** 172514
- [197] Hauptmann N, Dupé M, Hung T-C, Lemmens A K, Wegner D, Dupé B and Khajetoorians A A 2018 Revealing the correlation between real-space structure and chiral magnetic order at the atomic scale *Phys. Rev. B* **97** 100401
- [198] Jeon S, Xie Y, Li J, Wang Z, Bernevig B A and Yazdani A 2017 Distinguishing a Majorana zero mode using spin-resolved measurements *Science* **358** 772–6
- [199] Meier F, Zhou L, Wiebe J and Wiesendanger R 2008 Revealing magnetic interactions from single-atom magnetization curves *Science* **320** 82–86

- [200] Paul W, Yang K, Baumann S, Romming N, Choi T, Lutz C and Heinrich A 2017 Control of the millisecond spin lifetime of an electrically probed atom *Nat. Phys.* **13** 403–7
- [201] Baumann S, Paul W, Choi T, Lutz C P, Ardavan A and Heinrich A J 2015 Electron paramagnetic resonance of individual atoms on a surface *Science* **350** 417–20
- [202] Yang K, Paul W, Phark S-H, Willke P, Bae Y, Choi T, Esat T, Ardavan A, Heinrich A J and Lutz C P 2019 Coherent spin manipulation of individual atoms on a surface *Science* **366** 509–12
- [203] Wiesendanger R 2009 Spin mapping at the nanoscale and atomic scale *Rev. Mod. Phys.* **81** 1495–550
- [204] Romming N, Hanneken C, Menzel M, Bickel J E, Wolter B, von Bergmann K, Kubetzka A and Wiesendanger R 2013 Writing and deleting single magnetic skyrmions *Science* **341** 636–9
- [205] Kamber U *et al* 2020 Self-induced spin glass state in elemental and crystalline neodymium *Science* **368** eaay6757
- [206] Khajetoorians A A, Wiebe J, Chilian B, Lounis S, Blügel S and Wiesendanger R 2012 Atom-by-atom engineering and magnetometry of tailored nanomagnets *Nat. Phys.* **8** 497–503
- [207] Loth S, Eitzkorn M, Lutz C P, Eigler D M and Heinrich A J 2010 Measurement of fast electron spin relaxation times with atomic resolution *Science* **329** 1628–30
- [208] Hauptmann N, Gerritsen J W, Wegner D and Khajetoorians A A 2017 Sensing noncollinear magnetism at the atomic scale combining magnetic exchange and spin-polarized imaging *Nano Lett.* **17** 5660–5
- [209] Schmidt R, Lazo C, Kaiser U, Schwarz A, Heinze S and Wiesendanger R 2011 Quantitative measurement of the magnetic exchange interaction across a vacuum gap *Phys. Rev. Lett.* **106** 257202
- [210] Cocker T L, Peller D, Yu P, Repp J and Huber R 2016 Tracking the ultrafast motion of a single molecule by femtosecond orbital imaging *Nature* **539** 263–7
- [211] Hänze M, McMurtrie G, Baumann S, Malavolti L, Coppersmith S N and Loth S 2022 Quantum stochastic resonance of individual Fe atoms *Sci. Adv.* **7** eabg2616
- [212] Willke P, Bae Y, Yang K, Lado J L, Ferrón A, Choi T, Ardavan A, Fernández-Rossier J, Heinrich A J and Lutz C P 2018 Hyperfine interaction of individual atoms on a surface *Science* **362** 336–9
- [213] Choi T, Paul W, Rolf-Pissarczyk S, Macdonald A J, Natterer F D, Yang K, Willke P, Lutz C P and Heinrich A J 2017 Atomic-scale sensing of the magnetic dipolar field from single atoms *Nat. Nanotechnol.* **12** 420–4
- [214] Pielmeier F and Giessibl F J 2013 Spin resolution and evidence for superexchange on NiO(001) observed by force microscopy *Phys. Rev. Lett.* **110** 266101
- [215] Gross L, Mohn F, Moll N, Liljeroth P and Meyer G 2009 The chemical structure of a molecule resolved by atomic force microscopy *Science* **325** 1110–4
- [216] Sellies L, Spachtholz R, Bleher S, Eckrich J, Scheuerer P and Repp J 2023 Single-molecule electron spin resonance by means of atomic force microscopy *Nature* **624** 64–68
- [217] Bistrizter R and MacDonald A H 2011 Moiré bands in twisted double-layer graphene *Proc. Natl Acad. Sci.* **108** 12233–7
- [218] Mishra S *et al* 2021 Observation of fractional edge excitations in nanographene spin chains *Nature* **598** 287–92
- [219] Jolie W *et al* 2019 Tomonaga-Luttinger liquid in a box: electrons confined within MoS<sub>2</sub> mirror-twin boundaries *Phys. Rev. X* **9** 011055
- [220] Tokura Y, Yasuda K and Tsukazaki A 2019 Magnetic topological insulators *Nat. Rev. Phys.* **1** 126–43
- [221] Kiraly B, Knol E J, van Weerdenburg W M J, Kappen H J and Khajetoorians A A 2021 An atomic Boltzmann machine capable of self-adaption *Nat. Nanotechnol.* **16** 414–20
- [222] Veldman L M, Farinacci L, Rejali R, Broekhoven R, Gobeil J, Coffey D, Ternes M and Otte A F 2021 Free coherent evolution of a coupled atomic spin system initialized by electron scattering *Science* **372** 964–8
- [223] Sidles J A 1991 Noninductive detection of single-proton magnetic resonance *Appl. Phys. Lett.* **58** 2854
- [224] Sidles J A 1992 Folded Stern-Gerlach experiment as a means for detecting nuclear magnetic resonance in individual nuclei *Phys. Rev. Lett.* **68** 1124
- [225] Rugar D, Budakian R, Mamin H J and Chui B W 2004 Single spin detection by magnetic resonance force microscopy *Nature* **430** 329
- [226] Degen C L, Poggio M, Mamin H J, Rettner C T and Rugar D 2009 Nanoscale magnetic resonance imaging *Proc. Natl Acad. Sci.* **106** 1313
- [227] Haas S *et al* 2022 Nuclear magnetic resonance diffraction with subangstrom precision *Proc. Natl Acad. Sci.* **119** e2209213119
- [228] Budakian R *et al* 2024 *Nanotechnology* (in press) (<https://doi.org/10.1088/1361-6528/ad4b23>)
- [229] Ajoy A, Bissbort U, Lukin M D, Walsworth R L and Cappellaro P 2015 Atomic-scale nuclear spin imaging using quantum-assisted sensors in diamond *Phys. Rev. X* **5** 011001
- [230] Kost M, Cai J and Plenio M B 2015 Resolving single molecule structures with nitrogen-vacancy centers in diamond *Sci. Rep.* **5** 11007
- [231] Shi F *et al* 2015 Single-protein spin resonance spectroscopy under ambient conditions *Science* **347** 1135
- [232] Janitz E, Herb K, Völker L A, Huxter W S, Degen C L and Abendroth J M 2022 Diamond surface engineering for molecular sensing with nitrogen-vacancy centers *J. Mater. Chem. C* **10** 13533
- [233] Tao Y, Navaretti P, Hauert R, Grob U, Poggio M and Degen C L 2015 Permanent reduction of dissipation in nanomechanical Si resonators by chemical surface protection *Nanotechnology* **26** 465501
- [234] Heritier M, Pachlatko R, Tao Y, Abendroth J M, Degen C L and Eichler A 2021 Spatial correlation between fluctuating and static fields over metal and dielectric substrates *Phys. Rev. Lett.* **127** 216101
- [235] Eichler A 2022 Ultra-high-Q nanomechanical resonators for force sensing *Mater. Quantum Technol.* **2** 043001
- [236] Oppliger J and Natterer F D 2020 Sparse sampling for fast quasiparticle-interference mapping *Phys. Rev. Res.* **2** 023117
- [237] Abobeih M H, Randall J, Bradley C E, Bartling H P, Bakker M A, Degen M J, Markham M, Twitche D J and Taminiau T H 2019 Atomic-scale imaging of a 27-nuclear-spin cluster using a quantum sensor *Nature* **576** 411
- [238] Cujia K S, Herb K, Zopes J, Abendroth J M and Degen C L 2022 Parallel detection and spatial mapping of large nuclear spin clusters *Nat. Commun.* **13** 1260
- [239] Bucher D B, Glenn D R, Park H, Lukin M D and Walsworth R L 2020 Hyperpolarization-enhanced NMR spectroscopy with femtomole sensitivity using quantum defects in diamond *Phys. Rev. X* **10** 021053
- [240] Vinante A, Wijts G, Usenko O, Schinkelshoek L and Oosterkamp T H 2011 Magnetic resonance force microscopy of paramagnetic electron spins at millikelvin temperatures *Nat. Commun.* **2** 572
- [241] McCord J 2015 Progress in magnetic domain observation by advanced magneto-optical microscopy *J. Phys. D: Appl. Phys.* **48** 333001

- [242] Schäfer R and McCord J 2021 Magneto-optical microscopy *Magnetic Measurement Techniques for Materials Characterization* (Springer International Publishing) pp 171–229
- [243] Hubert A and Schäfer R 1998 *Magnetic Domains: The Analysis of Magnetic Microstructures* 1st edn (Springer)
- [244] Meguro S, Akahane K and Saito S 2016 Photographic observation of magnetic domain structure with three-dimensional local magnetization direction *Jpn. J. Appl. Phys.* **55** 07MF04
- [245] Soldatov I V and Schäfer R 2017 Selective sensitivity in Kerr microscopy *Rev. Sci. Instrum.* **88** 073701
- [246] Holländer R B, Müller C, Lohmann M, Mozooni B and McCord J 2017 Component selection in time-resolved magneto-optical wide-field imaging for the investigation of magnetic microstructures *J. Magn. Magn. Mater.* **432** 283–90
- [247] Büscher P and Reimer L 1993 Imaging and colour coding of magnetic domains by Kerr scanning optical microscopy *Scanning* **15** 123–9
- [248] Freeman M R and Hiebert W K 2001 Stroboscopic microscopy of magnetic dynamics *Spin Dynamics in Confined Magnetic Structures I* (Springer) pp 93–126
- [249] Hashimoto Y *et al* 2014 Ultrafast time-resolved magneto-optical imaging of all-optical switching in GdFeCo with femtosecond time-resolution and a  $\mu\text{m}$  spatial-resolution *Rev. Sci. Instrum.* **85** 063702
- [250] Vogel M, Hoffmann A and te Velthuis S G E 2020 Super resolution for magneto-optical microscopy *US Patent* 11474283
- [251] Keatley P S, Loughran T H J, Hendry E, Barnes W L, Hicken R J, Childress J R and Katine J A 2017 A platform for time-resolved scanning Kerr microscopy in the near-field *Rev. Sci. Instrum.* **88** 1–26
- [252] Jovičević-Klug M, Jovičević-Klug P, McCord J and Podgornik B 2021 Investigation of microstructural attributes of steel surfaces through magneto-optical Kerr effect *J. Mater. Res. Technol.* **11** 1245–59
- [253] McCormick E J, Newburger M J, Luo Y K, McCreary K M, Singh S, Martin I B, Cichewicz E J, Jonker B T and Kawakami R K 2018 Imaging spin dynamics in monolayer  $\text{WS}_2$  by time-resolved Kerr rotation microscopy *2D Mater.* **5** 011010
- [254] Shi S *et al* 2019 All-electric magnetization switching and Dzyaloshinskii–Moriya interaction in  $\text{WTe}_2$ /ferromagnet heterostructures *Nat. Nanotechnol.* **14** 945–9
- [255] Jiang W *et al* 2015 Blowing magnetic skyrmion bubbles *Science* **349** 283–6
- [256] Higo T *et al* 2018 Large magneto-optical Kerr effect and imaging of magnetic octupole domains in an antiferromagnetic metal *Nat. Photon.* **12** 79–78
- [257] Xu J *et al* 2019 Imaging antiferromagnetic domains in nickel oxide thin films by optical birefringence effect *Phys. Rev. B* **100** 1–8
- [258] Lahtinen T H E, Tuomi J O and van Dijken S 2011 Pattern transfer and electric-field-induced magnetic domain formation in multiferroic heterostructures *Adv. Mater.* **23** 3187–91
- [259] Müller C, Durdaut P, Holländer R B, Kittmann A, Schell V, Meyners D, Höft M, Quandt E and McCord J 2022 Imaging of love waves and their interaction with magnetic domain walls in magnetoelectric magnetic field sensors *Adv. Electron. Mater.* **8** 2200033
- [260] Dreyer R, Schäffer A F, Bauer H G, Liebing N, Berakdar J and Woltersdorf G 2022 Imaging and phase-locking of non-linear spin waves *Nat. Commun.* **13** 4939
- [261] Jangid R *et al* 2023 Evidence of extreme domain wall speeds under ultrafast optical excitation (arXiv:2303.16131)
- [262] Shen Y R 2002 *The Principles of Nonlinear Optics* (Wiley)
- [263] Franken P A, Hill A E, Peters C W and Weinreich G 1961 Generation of optical harmonics *Phys. Rev. Lett.* **7** 118
- [264] Birss R R 1966 *Symmetry and Magnetism* (North-Holland Publishing Company)
- [265] Fiebig M 2023 *Nonlinear Optics on Ferromagnetic Materials* (Wiley)
- [266] Wadhawan V K 2000 *Introduction to Ferromagnetic Materials* (Gordon and Breach)
- [267] Reif J, Zink J C, Schneider C-M and Kirschner J 1991 Effects of surface magnetism on optical second harmonic generation *Phys. Rev. Lett.* **67** 2878
- [268] Cheong S-W, Fiebig M, Wu W, Chapon L and Kiryukhin V 2020 Seeing is believing: visualization of antiferromagnetic domains *npj Quantum Mater.* **5** 3
- [269] van Aken B B, Rivera J P, Schmid H and Fiebig M 2008 Anisotropy of antiferromagnetic  $180^\circ$  domains in  $\text{LiNiPO}_4$  and  $\text{LiCoPO}_4$  *Phys. Rev. Lett.* **101** 157202
- [270] Tunyagi A R, Ulex M and Betzler K 2003 Noncollinear optical frequency doubling in strontium barium niobate *Phys. Rev. Lett.* **90** 243901
- [271] Zayats A V 2007 Apertureless near-field microscopy of second-harmonic generation *Tip Enhancement (Advances in Nano-Optics and Nano-Photonics)* ed S Kawata and V M Shalaev (Elsevier)
- [272] Denev S A, Lummen T T A, Barnes E, Kumar A and Gopalan V 2011 Probing ferroelectrics using optical second harmonic generation *J. Am. Ceram. Soc.* **94** 2699
- [273] Zu R, Wang B, He J, Wang -J-J, Weber L, Chen L-Q and Gopalan V 2022 Analytical and numerical modeling of optical second harmonic generation in anisotropic crystals using #SHAARP package *npj Comput. Mater.* **8** 246
- [274] De Luca G, Strkalj N, Manz S, Bouillet C, Fiebig M and Trassin M 2017 Nanoscale design of polarization in ultrathin ferroelectric heterostructures *Nat. Commun.* **8** 1419
- [275] Ahn Y, Guo X, Xue R, Qu K, Sun K, Mandrus D and Zhao L 2023 Electric quadrupole second-harmonic generation revealing dual magnetic orders in a magnetic Weyl semimetal *Nat. Photon.* **18** 26–31
- [276] García J D, Gómez F, Velázquez V and Rodríguez B A 2022 Bulk second-order optical nonlinearities in centrosymmetric materials *J. Opt. Soc. Am. B* **39** 2582
- [277] Song Q *et al* 2022 Evidence for a single-layer van der Waals multiferroic *Nature* **602** 601
- [278] Jungwirth T, Sinova J, Manchon A, Marti X, Wunderlich J and Felser C 2018 The multiple directions of antiferromagnetic spintronics *Nat. Phys.* **14** 200
- [279] Zhao L, Torchinsky D H, Chu H, Ivanov V, Lifshitz R, Flint R, Qi T, Cao G and Hsieh D 2016 Evidence of an odd-parity hidden order in a spin-orbit coupled correlated iridate *Nat. Phys.* **12** 32–36
- [280] Nakamura S, Katsumi K, Terai H and Shimano R 2020 Nonreciprocal terahertz second-harmonic generation in superconducting NbN under supercurrent injection *Phys. Rev. Lett.* **125** 097004
- [281] Juraschek D M and Spaldin N A 2019 Orbital magnetic moments of phonons *Phys. Rev. Mater.* **3** 064405
- [282] Scholl A, Ohldag H, Nolting F, Stöhr J and Padmore H A 2002 X-ray photoemission electron microscopy, a tool for the investigation of complex magnetic structures (invited) *Rev. Sci. Instrum.* **73** 1362–6
- [283] Stoehr J and Siegmund H C 2006 *Magnetism: From Fundamentals to Nanoscale Dynamics* (Springer)
- [284] Kuch W, Schäfer R, Fischer P and Hillebrecht F 2015 *Magnetic Microscopy of Layered Structures* (Springer)
- [285] Fischer P 2021 *Nanomagnetic Materials* ed A Yamaguchi, A Hirohata and B J H Stadler (Elsevier) pp 199–303

- [286] Acremann Y, Chembrolu V, Strachan J P, Tyliczszak T and Stoehr J 2007 Software defined photon counting system for time resolved x-ray experiments *Rev. Sci. Instrum.* **78** 014702
- [287] Stoll H *et al* 2004 High-resolution imaging of fast magnetization dynamics in magnetic nanostructures *Appl. Phys. Lett.* **84** 3328–30
- [288] Choe S-B, Acremann Y, Scholl A, Bauer A, Doran A, Stohr J and Padmore H A 2004 Vortex core-driven magnetization dynamics *Science* **304** 420
- [289] Streubel R, Kronast F, Fischer P, Parkinson D, Schmidt O G and Makarov D 2015 Retrieving three-dimensional spin textures by full-field magnetic soft x-ray tomography *Nat. Commun.* **6** 7612
- [290] Witte K *et al* 2020 From 2D STXM to 3D imaging: soft x-ray laminography of thin specimens *Nano Lett.* **20** 1305–14
- [291] Finizio S, Watts B and Raabe J 2021 Why is my image noisy? A look into the terms contributing to a time-resolved x-ray microscopy image *J. Synchrotron Radiat.* **28** 1146–58
- [292] Raftrey D, Hierro-Rodriguez A, Fernandez-Pacheco A and Fischer P 2022 The road to 3-dim nanomagnetism: steep curves and architected crosswalks *J. Magn. Magn. Mater.* **563** 169899
- [293] Hierro-Rodriguez A, Quirós C, Sorrentino A, Alvarez-Prado L M, Martín J I, Alameda J M, McVitie S, Pereiro E, Vélez M and Ferrer S 2020 Revealing 3D magnetization of thin films with soft x-ray tomography: magnetic singularities and topological charges *Nat. Commun.* **11** 6382
- [294] Levitan A L, Keskinbora K, Sanli U T, Weigand M and Comin R 2020 Single-frame far-field diffractive imaging with randomized illumination *Opt. Express* **28** 37103–17
- [295] Chen G, Ma T, N'Diaye A T, Kwon H, Won C, Wu Y and Schmid A K 2013 Tailoring the chirality of magnetic domain walls by interface engineering *Nat. Commun.* **4** 2671
- [296] Tusche C, Krasnyuk A and Kirschner J 2015 Spin resolved bandstructure imaging with a high resolution momentum microscope *Ultramicroscopy* **159** 520–9
- [297] Fischer P, Sanz-Hernández D, Streubel R and Fernández-Pacheco A 2020 Launching a new dimension with 3D magnetic nanostructures *APL Mater.* **8** 010701
- [298] Šmejkal L, Sinova J and Jungwirth T 2022 Emerging research landscape of altermagnetism *Phys. Rev. X* **12** 040501
- [299] Carra P, Thole B T, Altarelli M and Wang X 1993 x-ray circular dichroism and local magnetic fields *Phys. Rev. Lett.* **70** 694–7
- [300] Thole B T, Carra P, Sette F and van der Laan G 1992 x-ray circular dichroism as a probe of orbital magnetization *Phys. Rev. Lett.* **68** 1943–6
- [301] Mayr S *et al* 2021 Spin-wave emission from vortex cores under static magnetic bias fields *Nano Lett.* **21** 1584–90
- [302] Seaberg M H *et al* 2017 Nanosecond x-ray photon correlation spectroscopy on magnetic skyrmions *Phys. Rev. Lett.* **119** 067403
- [303] Kukreja R, Bonetti S, Chen Z, Backes D, Acremann Y, Katine J A, Kent A D, Dürr H A, Ohldag H and Stöhr J 2015 X-ray detection of transient magnetic moments induced by a spin current in Cu *Phys. Rev. Lett.* **115** 096601
- [304] Donnelly C *et al* 2020 Time-resolved imaging of three-dimensional nanoscale magnetization dynamics *Nat. Nanotechnol.* **15** 356–60
- [305] Mayr S, Finizio S, Reuteler J, Stutz S, Dubs C, Weigand M, Hrabec A, Raabe J and Wintz S 2021 *Crystals* **11** 546
- [306] Eisebitt S, Lüning J, Schlotter W F, Lörger M, Hellwig O, Eberhardt W and Stöhr J 2004 Lensless imaging of magnetic nanostructures by x-ray spectro-holography *Nature* **432** 885–8
- [307] Tripathi A, Mohanty J, Dietze S H, Shpyrko O G, Shipton E, Fullerton E E, Kim S S and McNulty I 2011 Dichroic coherent diffractive imaging *Proc. Natl Acad. Sci.* **108** 13393–8
- [308] Büttner F *et al* 2015 Dynamics and inertia of skyrmionic spin structures *Nat. Phys.* **11** 225–8
- [309] Donnelly C, Guizar-Sicairos M, Scagnoli V, Gliga S, Holler M, Raabe J and Heyderman L J 2017 Three-dimensional magnetization structures revealed with x-ray vector nanotomography *Nature* **547** 328–31
- [310] Zayko S, Kfir O, Heigl M, Lohmann M, Sivis M, Albrecht M and Ropers C 2021 Ultrafast high-harmonic nanoscopy of magnetization dynamics *Nat. Commun.* **12** 6337
- [311] Klose C *et al* 2023 Coherent correlation imaging for resolving fluctuating states of matter *Nature* **614** 256–61
- [312] Scherz A *et al* 2007 Phase imaging of magnetic nanostructures using resonant soft x-ray holography *Phys. Rev. B* **76** 214410
- [313] Donnelly C *et al* 2016 High-resolution hard x-ray magnetic imaging with dichroic ptychography *Phys. Rev. B* **94** 064421
- [314] Di Pietro Martinez *et al* 2023 Three-dimensional tomographic imaging of the magnetization vector field using Fourier transform holography *Phys. Rev. B* **107** 094425
- [315] Neethirajan J N *et al* 2023 Soft x-ray phase nano-microscopy of micrometre-thick magnets (arXiv:2309.14969)
- [316] Battistelli R *et al* 2024 Coherent x-ray magnetic imaging with 5 nm resolution *Optica* **11**, 2 234–7
- [317] Rana A *et al* 2023 Three-dimensional topological magnetic monopoles and their interactions in a ferromagnetic meta-lattice *Nat. Nanotechnol.* **18** 227–32
- [318] Shi X *et al* 2016 Soft x-ray ptychography studies of nanoscale magnetic and structural correlations in thin SmCo<sub>5</sub> films *Appl. Phys. Lett.* **108** 094103
- [319] Zhu X, Hitchcock A P, Bazylinski D A, Denes P, Joseph J, Lins U, Marchesini S, Shiu H-W, Tyliczszak T and Shapiro D A 2016 Measuring spectroscopy and magnetism of extracted and intracellular magnetosomes using soft x-ray ptychography *Proc. Natl Acad. Sci.* **113** E8219–E8227
- [320] Kfir O *et al* 2017 Nanoscale magnetic imaging using circularly polarized high-harmonic radiation *Sci. Adv.* **3** eaao4641
- [321] Ferrand P, Baroni A, Allain M and Chamard V 2018 Quantitative imaging of anisotropic material properties with vectorial ptychography *Opt. Lett.* **43** 763–6
- [322] Song Q *et al* 2020 Ptychography retrieval of fully polarized holograms from geometric-phase metasurfaces *Nat. Commun.* **11** 2651
- [323] Mutzafi M, Shechtman Y, Eldar Y C, Cohen O and Segev M 2015 Sparsity-based ankylography for recovering 3D molecular structures from single-shot 2D scattered light intensity *Nat. Commun.* **6** 7950
- [324] Wu L, Juhás P, Yoo S and Robinson I 2021 Complex imaging of phase domains by deep neural networks *IUCr* **8** 12–21
- [325] Cherukara M J, Zhou T, Nashed Y, Enfedaque P, Hexemer A, Harder R J and Holt M V 2020 AI-enabled high-resolution scanning coherent diffraction imaging *Appl. Phys. Lett.* **117** 044103
- [326] Chan H, Nashed Y S G, Kandel S, Hruszkewycz S O, Sankaranarayanan S K R S, Harder R J and Cherukara M J 2021 Rapid 3D nanoscale coherent imaging via physics-aware deep learning *Appl. Phys. Rev.* **8** 021407
- [327] Birch M T *et al* 2020 Real-space imaging of confined magnetic skyrmion tubes *Nat. Commun.* **11** 1726
- [328] Grelrier M *et al* 2023 X-ray holography of skyrmionic cocoons in aperiodic magnetic multilayers *Phys. Rev. B* **107** L220405
- [329] Li W *et al* 2019 Anatomy of skyrmionic textures in magnetic multilayers *Adv. Mater.* **31** 1807683



- [330] Donnelly C, Metlov K L, Scagnoli V, Guizar-Sicarios M, Holler M, Bingham N S, Raabe J, Heyderman L J, Cooper N R and Gliga S 2021 Experimental observation of vortex rings in a bulk magnet *Nat. Phys.* **17** 316–21
- [331] Kim M G *et al* 2018 Imaging antiferromagnetic antiphase domain boundaries using magnetic Bragg diffraction phase contrast *Nat. Commun.* **9** 5013
- [332] Kim M G *et al* 2022 Real-space observation of fluctuating antiferromagnetic domains *Sci. Adv.* **8** eabj9493
- [333] Chen X M *et al* 2019 Spontaneous magnetic superdomain wall fluctuations in an artificial antiferromagnet *Phys. Rev. Lett.* **123** 197202
- [334] Zhang B, Gardner D F, Seaberg M D, Shanblatt E R, Kapteyn H C, Murnane M M and Adams D E 2015 High contrast 3D imaging of surfaces near the wavelength limit using tabletop EUV ptychography *Ultramicroscopy* **158** 98–104
- [335] McCarter M R *et al* 2023 Antiferromagnetic real-space configuration probed by dichroism in scattered x-ray beams with orbital angular momentum *Phys. Rev. B* **107** L060407
- [336] von Korff Schmising C *et al* 2014 Imaging ultrafast demagnetization dynamics after a spatially localized optical excitation *Phys. Rev. Lett.* **112** 217203
- [337] Applied Solid State Science 1979 *Magnetic Domain Walls in Bubble Materials II* ed A P Malozemoff and J C Slonczewski (Academic) (<https://doi.org/10.1016/B978-0-12-002951-8.50001-7>)
- [338] Thiele A A 1973 Steady-state motion of magnetic domains *Phys. Rev. Lett.* **30** 230–3
- [339] Yoshimura Y, Kim K-J, Taniguchi T, Tono T, Ueda K, Hiramatsu R, Moriyama T, Yamada K, Nakatani Y and Ono T 2016 Soliton-like magnetic domain wall motion induced by the interfacial Dzyaloshinskii–Moriya interaction *Nat. Phys.* **12** 157–61
- [340] Reichhardt C, Reichhardt C J O and Milošević M V 2022 Statics and dynamics of skyrmions interacting with disorder and nanostructures *Rev. Mod. Phys.* **94** 035005
- [341] Im M-Y, Han H-S, Jung M-S, Yu Y-S, Lee S, Yoon S, Chao W, Fischer P, Hong J-I and Lee K-S 2019 Dynamics of the Bloch point in an asymmetric permalloy disk *Nat. Commun.* **10** 593
- [342] Pfau B, Günther C M, Guehrs E, Hauet T, Hennen T, Eisebitt S and Hellwig O 2014 Influence of stray fields on the switching-field distribution for bit-patterned media based on pre-patterned substrates *Appl. Phys. Lett.* **105** 132407
- [343] Woo S *et al* 2017 Spin-orbit torque-driven skyrmion dynamics revealed by time-resolved x-ray microscopy *Nat. Commun.* **8** 15573
- [344] Litzius K *et al* 2017 Skyrmion Hall effect revealed by direct time-resolved x-ray microscopy *Nat. Phys.* **13** 170
- [345] Baumgartner M *et al* 2017 Spatially and time-resolved magnetization dynamics driven by spin–orbit torques *Nat. Nanotechnol.* **12** 980
- [346] Van Waeyenberge B *et al* 2006 Magnetic vortex core reversal by excitation with short bursts of an alternating field *Nature* **444** 461
- [347] Dieterle G *et al* 2019 Coherent excitation of heterosymmetric spin waves with ultrashort wavelengths *Phys. Rev. Lett.* **122** 117202
- [348] Baumgaertl K, Grafe J, Che P, Mucchietto A, Forster J, Trager N, Bechtel M, Weigand M, Schutz G and Grundler D 2020 Nanoimaging of ultrashort magnon emission by ferromagnetic grating couplers at GHz frequencies *Nano Lett.* **20** 7281
- [349] Wang T H *et al* 2012 Femtosecond single-shot imaging of nanoscale ferromagnetic order in Co/Pd multilayers using resonant x-ray holography *Phys. Rev. Lett.* **108** 267403
- [350] Schneider M *et al* 2020 Ultrafast demagnetization dominates fluence dependence of magnetic scattering at Co M edges *Phys. Rev. Lett.* **125** 127201
- [351] Higley D J *et al* 2019 Femtosecond x-ray induced changes of the electronic and magnetic response of solids from electron redistribution *Nat. Commun.* **10** 5289
- [352] Hagstrom N Z *et al* 2022 Megahertz-rate ultrafast x-ray scattering and holographic imaging at the European XFEL *J. Synchrotron Radiat.* **29** 1454
- [353] Bonetti S *et al* 2015 Microwave soft x-ray microscopy for nanoscale magnetization dynamics in the 5–10 GHz frequency range *Rev. Sci. Instrum.* **86** 093703
- [354] Finizio S, Mayr S and Raabe J 2020 Time-of-arrival detection for time-resolved scanning transmission x-ray microscopy imaging *J. Synchrotron Radiat.* **27** 1320
- [355] Finizio S, Bailey J, Olsthoorn B and Raabe J 2022 Periodogram-based detection of unknown frequencies in time-resolved scanning transmission x-ray microscopy *ACS Nano* **16** 21071
- [356] Gerlinger K *et al* 2023 Pump-probe x-ray microscopy of photo-induced magnetization dynamics at MHz repetition rates *Struct. Dyn.* **10** 024301
- [357] Donnelly C, Finizio S, Gliga S, Holler M, Odstrcil M, Mayr S, Scagnoli V, Heyderman L J, Guizar-Sicarios M and Raabe J 2020 Imaging three-dimensional magnetization dynamics with laminography *Nat. Nanotechnol.* **15** 356
- [358] Shpyrko O G *et al* 2007 Direct measurement of antiferromagnetic domain fluctuations *Nature* **447** 7140
- [359] Konings S, Schüßler-Langeheine C, Ott H, Weschke E, Schierle E, Zabel H and Goedkoop J B 2011 Magnetic domain fluctuations in an antiferromagnetic film observed with coherent resonant soft x-ray scattering *Phys. Rev. Lett.* **106** 077402
- [360] Seaberg M H *et al* 2021 Spontaneous fluctuations in a magnetic Fe/Gd skyrmion lattice *Phys. Rev. Res.* **3** 033249
- [361] Finizio S, Donnelly C, Mayr S, Hrabec A and Raabe J 2022 Three-dimensional vortex gyration dynamics unraveled by time-resolved soft x-ray laminography with freely selectable excitation frequencies *Nano Lett.* **22** 1971
- [362] Fan G *et al* 2022 Ultrafast magnetic scattering on ferrimagnets enabled by a bright Yb-based soft x-ray source *Optica* **9** 399
- [363] Fohler M *et al* 2017 A general approach to obtain soft x-ray transparency for thin films grown on bulk substrates *Rev. Sci. Instrum.* **88** 103701
- [364] Sharma S, Elliott P and Shallcross S 2023 THz induced giant spin and valley currents *Sci. Adv.* **9** eadf3673
- [365] Büttner F *et al* 2021 Observation of fluctuation-mediated picosecond nucleation of a topological phase *Nat. Mater.* **20** 30
- [366] Bisig A *et al* 2013 Correlation between spin structure oscillations and domain wall velocities *Nat. Commun.* **4** 2328
- [367] Oepen H P and Frömter R 2007 Scanning electron microscopy with polarisation analysis *Handbook of Magnetism and Advanced Magnetic Materials* (Wiley) (<https://doi.org/10.1002/9780470022184.hmm309>)
- [368] Koike K and Hayakawa K 1984 Scanning electron microscope observation of magnetic domains using spin-polarized secondary electrons *Jpn. J. Appl. Phys.* **23** L187
- [369] Unguris J, Hembree G G, Celotta R J and Pierce D T 1985 High resolution magnetic microstructure imaging using secondary electron spin polarization analysis in a scanning electron microscope *J. Microsc.* **139** RP1–RP2
- [370] Oepen H P and Kirschner J 1988 Imaging of magnetic microstructures at surfaces *J. Phys. Colloq.* **49** C8–1857
- [371] Allenspach R 1994 Ultrathin films: magnetism on the microscopic scale *J. Magn. Magn. Mater.* **129** 160–85
- [372] Frömter R, Hankemeier S, Oepen H P and Kirschner J 2011 Optimizing a low-energy electron diffraction spin-polarization analyzer for imaging of magnetic surface structures *Rev. Sci. Instrum.* **82** 033704



- [373] Kohashi T, Konoto M and Koike K 2004 A spin rotator for spin-polarized scanning electron microscopy *Rev. Sci. Instrum.* **75** 2003–7
- [374] Kloodt-Twesten F, Kuhrau S, Oepen H P and Frömter R 2019 Measuring the Dzyaloshinskii-Moriya interaction of the epitaxial Co/Ir(111) interface *Phys. Rev. B* **100** 100402
- [375] Lucassen J, Meijer M J, Kurnosikov O, Swagten H J, Koopmans B, Lavrijsen R, Kloodt-Twesten F, Frömter R and Duine R A 2019 Tuning magnetic chirality by dipolar interactions *Phys. Rev. Lett.* **123** 157201
- [376] Koike K 2013 Spin-polarized scanning electron microscopy *Microscopy* **62** 177–91
- [377] Frömter R, Kloodt F, Rößler S, Frauen A, Staeck P, Cavicchia D R, Bocklage L, Röbisch V, Quandt E and Oepen H P 2016 Time-resolved scanning electron microscopy with polarization analysis *Appl. Phys. Lett.* **108** 142401
- [378] Schönte D, Oelsner A, Krautscheid P, Reeve R M and Kläui M 2018 Development of a scanning electron microscopy with polarization analysis system for magnetic imaging with ns time resolution and phase-sensitive detection *Rev. Sci. Instrum.* **89** 083703
- [379] Escher M, Weber N B, Merkel M, Plucinski L and Schneider C M 2011 FERRUM: a new highly efficient spin detector for electron spectroscopy *e-J. Surf. Sci. Nanotechnol.* **9** 340–3
- [380] Gusev S A, Petrov V N and Skorokhodov E V 2010 Problems of implementing SEMPA in experiments *J. Surf. Investig.* **4** 582–7
- [381] Steierl G, Liu G, Iorgov D and Kirschner J 2002 Surface domain imaging in external magnetic fields *Rev. Sci. Instrum.* **73** 4264–9
- [382] Zhou X, Matthes F, Bürgler D E and Schneider C M 2016 Magnetic surface domain imaging of uncapped epitaxial FeRh(001) thin films across the temperature-induced metamagnetic transition *AIP Adv.* **6** 015211
- [383] VanZandt T, Browning R and Landolt M 1991 Iron overlayer polarization enhancement technique for spin-polarized electron microscopy *J. Appl. Phys.* **69** 1564–8
- [384] Hankemeier S, Frömter R, Mikuszeit N, Stickler D, Stillrich H, Pütter S, Vedmedenko E Y and Oepen H P 2009 Magnetic ground state of single and coupled permalloy rectangles *Phys. Rev. Lett.* **103** 147204
- [385] McMorran B J, Cochran A C, Dumas R K, Liu K, Morrow P, Pierce D T and Unguris J 2010 Measuring the effects of low energy ion milling on the magnetization of Co/Pd multilayers using scanning electron microscopy with polarization analysis *J. Appl. Phys.* **107** 09D305
- [386] Kohashi T, Motai K, Matsuyama H and Maruyama Y 2021 Operando observation of magnetism in HDD writing heads by spin-polarized scanning electron microscopy *Microscopy* **70** 436–41
- [387] Casiraghi A, Rincón Domínguez T, Rößler S, Franke K J A, López González D, Hämäläinen S J, Frömter R, Oepen H P and van Dijken S 2015 Influence of elastically pinned magnetic domain walls on magnetization reversal in multiferroic heterostructures *Phys. Rev. B* **92** 054406
- [388] Li P *et al* 2017 Spatially resolved ferroelectric domain-switching-controlled magnetism in  $\text{Co}_{40}\text{Fe}_{40}\text{B}_{20}/\text{Pb}(\text{Mg}_{1/3}\text{Nb}_{2/3})_{0.7}\text{Ti}_{0.3}\text{O}_3$  multiferroic heterostructure *ACS Appl. Mater. Interfaces* **9** 2642–9
- [389] Gaul A, Hankemeier S, Holzinger D, Mücklich N D, Staeck P, Frömter R, Oepen H P and Ehresmann A 2016 Engineered magnetic domain textures in exchange bias bilayer systems *J. Appl. Phys.* **120** 033902
- [390] Seng B *et al* 2021 Direct imaging of chiral domain walls and Néel-type skyrmionium in ferrimagnetic alloys *Adv. Funct. Mater.* **31** 2102307
- [391] Dzemiantsova L V *et al* 2011 Multiscale magnetic study of Ni(111) and graphene on Ni(111) *Phys. Rev. B* **84** 205431
- [392] Kuhrau S, Kloodt-Twesten F, Heyn C, Oepen H P and Frömter R 2018 Cap-layer-dependent oxidation of ultrathin cobalt films and its effect on the magnetic contrast in scanning electron microscopy with polarization analysis *Appl. Phys. Lett.* **113** 172403
- [393] Meijer M J, Lucassen J, Kurnosikov O, Swagten H J, Koopmans B, Lavrijsen R, Kloodt-Twesten F, Frömter R and Duine R A 2020 Magnetic chirality controlled by the interlayer exchange interaction *Phys. Rev. Lett.* **124** 207203
- [394] Bhukta M *et al* 2024 Homochiral antiferromagnetic merons, antimerons and bimerons realized in synthetic antiferromagnets *Nat. Commun.* **15** 1641
- [395] Bommanaboyena S P *et al* 2021 Readout of an antiferromagnetic spintronics system by strong exchange coupling of  $\text{Mn}_2\text{Au}$  and permalloy *Nat. Commun.* **12** 1
- [396] Reeve R M, Mix C, König M, Foerster M, Jakob G and Kläui M 2013 Magnetic domain structure of  $\text{La}_{0.7}\text{Sr}_{0.3}\text{MnO}_3$  thin-films probed at variable temperature with scanning electron microscopy with polarization analysis *Appl. Phys. Lett.* **102** 122407
- [397] Lenz T *et al* 2021 Imaging topological spin structures using light-polarization and magnetic microscopy *Phys. Rev. Appl.* **15** 024040
- [398] Aharonov Y and Bohm D 1959 Significance of electromagnetic potentials in the quantum theory *Phys. Rev.* **115** 485–91
- [399] Roberts A, Nugent K A and Gureyev T E 1995 Phase retrieval with the transport-of-intensity equation: matrix solution with use of Zernike polynomials *J. Opt. Soc. Am. A* **12** 1932–41
- [400] Chess J J, Montoya S A, Harvey T R, Ophus C, Couture S, Lomakin V, Fullerton E E and McMorran B J 2017 Streamlined approach to mapping the magnetic induction of skyrmionic materials *Ultramicroscopy* **177** 78–83
- [401] Yu X Z, Onose Y, Kanazawa N, Park J H, Han J H, Matsui Y, Nagaosa N and Tokura Y 2010 Real-space observation of a two-dimensional skyrmion crystal *Nature* **465** 901–4
- [402] Chapman J N, Batson P E, Waddell E M and Ferrier R P 1978 The direct determination of magnetic domain wall profiles by differential phase contrast electron microscopy *Ultramicroscopy* **3** 203–14
- [403] Leuthner T, Lichte H and Herrmann K - H 1989 STEM-holography using the electron biprism *Phys. Status Solidi* **116** 113–21
- [404] Yasin F S, Harada K, Shindo D, Shinada H, McMorran B J and Tanigaki T 2018 A tunable path-separated electron interferometer with an amplitude-dividing grating beamsplitter *Appl. Phys. Lett.* **113** 233102
- [405] Karube K, Peng L, Masell J, Yu X, Kagawa F, Tokura Y and Taguchi Y 2021 Room-temperature antiskyrmions and sawtooth surface textures in a non-centrosymmetric magnet with  $S_4$  symmetry *Nat. Mater.* **20** 335–40
- [406] Zhang H *et al* 2022 A room temperature polar magnetic metal *Phys. Rev. Mater.* **6** 044403
- [407] Yu X *et al* 2021 Real-space observations of 60-nm skyrmion dynamics in an insulating magnet under low heat flow *Nat. Commun.* **12** 5079
- [408] Yasin F S, Masell J, Karube K, Kikkawa A, Taguchi Y, Tokura Y and Yu X 2022 Real-space determination of the isolated magnetic skyrmion deformation under electric current flow *Proc. Natl Acad. Sci.* **119** e2200958119
- [409] Yasin F S, Masell J, Karube K, Shindo D, Taguchi Y, Tokura Y and Yu X 2023 Heat current-driven topological spin texture transformations and helical  $q$ -vector switching *Nat. Commun.* **14** 1–8
- [410] Nguyen K X, Huang J, Karigerasi M H, Kang K, Cahill D G, Zuo J-M, Schleife A, Shoemaker D P and Huang P Y 2023 Angstrom-scale imaging of magnetization in antiferromagnetic  $\text{Fe}_2\text{As}$  via 4D-STEM *Ultramicroscopy* **247** 113696

- [411] Philipp H T *et al* 2022 Very-high dynamic range, 10,000 frames/second pixel array detector for electron microscopy *Microsc. Microanal.* **28** 425–40
- [412] Zhang L, Hoogenboom J P, Cook B and Kruit P 2019 Photoemission sources and beam blankers for ultrafast electron microscopy *Struct. Dyn.* **6** 51501
- [413] Zewail A H 2006 4D ultrafast electron diffraction, crystallography, and microscopy *Annu. Rev. Phys. Chem.* **57** 65–103
- [414] Shimojima T, Nakamura A and Ishizaka K 2023 Development of five-dimensional scanning transmission electron microscopy *Rev. Sci. Instrum.* **94** 023705
- [415] Pohl D, Lee Y, Krieger D, Beckert S, Schneider S, Rellinghaus B and Thomas A 2023 Probing magnetic properties at the nanoscale: in-situ Hall measurements in a TEM *Sci. Rep.* **13** 14871
- [416] Chen Z, Turgut E, Jiang Y, Nguyen K X, Stolt M J, Jin S, Ralph D C, Fuchs G D and Muller D A 2022 Lorentz electron ptychography for imaging magnetic textures beyond the diffraction limit *Nat. Nanotechnol.* **17** 1165–70
- [417] Hong X, Zeltmann S E, Savitzky B H, Rangel Dacosta L, Müller A, Minor A M, Bustillo K C and Ophus C 2021 Multibeam electron diffraction *Microsc. Microanal.* **27** 129–39
- [418] Yu X Z, Koshihara W, Tokunaga Y, Shibata K, Taguchi Y, Nagaosa N and Tokura Y 2018 Transformation between meron and skyrmion topological spin textures in a chiral magnet *Nature* **564** 95–98
- [419] Khanh N D *et al* 2020 Nanometric square skyrmion lattice in a centrosymmetric tetragonal magnet *Nat. Nanotechnol.* **15** 444–9
- [420] Yu X Z, Morikawa D, Nakajima K, Shibata K, Kanazawa N, Arima T, Nagaosa N and Tokura Y 2020 Motion tracking of 80-nm-size skyrmions upon directional current injections *Sci. Adv.* **6** eaaz9744
- [421] Kohno Y, Seki T, Findlay S D, Ikuhara Y and Shibata N 2022 Real-space visualization of intrinsic magnetic fields of an antiferromagnet *Nature* **602** 234–9
- [422] Powalla L *et al* 2023 Seeding and emergence of composite skyrmions in a van der Waals magnet *Adv. Mater.* **35** 2208930
- [423] Gatel C, Dupuy J, Houdellier F and Hÿtch M J 2018 Unlimited acquisition time in electron holography by automated feedback control of transmission electron microscope *Appl. Phys. Lett.* **113** 133102
- [424] Lubk A, Guzzinati G, Börrnert F and Verbeeck J 2013 Transport of intensity phase retrieval of arbitrary wave fields including vortices *Phys. Rev. Lett.* **111** 173902
- [425] Kovács A and Dunin-Borkowski R E 2018 Magnetic imaging of nanostructures using off-axis electron holography *Handbook of Magnetic Materials* (Elsevier) vol 27 pp 59–153
- [426] Sturm S, Sigleitmeier M, Wolf D, Vogel K, Gratz M, Faivre D, Lubk A, Büchner B, Sturm (Née Rosseeva) E V and Cölfen H 2019 Magnetic nanoparticle chains in gelatin ferrogels: bioinspiration from magnetotactic bacteria *Adv. Funct. Mater.* **29** 1905996
- [427] Andersen I M *et al* 2021 Field tunable three-dimensional magnetic nanotextures in cobalt-nickel nanowires *Phys. Rev. Res.* **3** 033085
- [428] Murakami Y, Kasai H, Kim J J, Mamishin S, Shindo D, Mori S and Tonomura A 2010 Ferromagnetic domain nucleation and growth in colossal magnetoresistive manganite *Nat. Nanotechnol.* **5** 37–41
- [429] Li Z-A, Zheng F, Tavabi A H, Caron J, Jin C, Du H, Kovács A, Tian M, Farle M and Dunin-Borkowski R E 2017 Magnetic skyrmion formation at lattice defects and grain boundaries studied by quantitative off-axis electron holography *Nano Lett.* **17** 1395–401
- [430] Kovács A, Caron J, Savchenko A S, Kiselev N S, Shibata K, Li Z-A, Kanazawa N, Tokura Y, Blügel S and Dunin-Borkowski R E 2017 Mapping the magnetization fine structure of a lattice of Bloch-type skyrmions in an FeGe thin film *Appl. Phys. Lett.* **111** 192410
- [431] Almeida T P, Palomino A, Lequeux S, Boureau V, Fruchart O, Prejbeanu I L, Dieny B and Cooper D 2022 Off-axis electron holography for the direct visualization of perpendicular shape anisotropy in nanoscale 3D magnetic random-access-memory devices *APL Mater.* **10** 061104
- [432] Wolf D *et al* 2019 Holographic vector field electron tomography of three-dimensional nanomagnets *Commun. Phys.* **2** 87
- [433] Wolf D, Schneider S, Rößler U K, Kovács A, Schmidt M, Dunin-Borkowski R E, Büchner B, Rellinghaus B and Lubk A 2022 Unveiling the three-dimensional magnetic texture of skyrmion tubes *Nat. Nanotechnol.* **17** 250–5
- [434] Wagner T, Niermann T, Urban F and Lehmann M 2019 Nanosecond electron holography by interference gating *Ultramicroscopy* **206** 112824
- [435] Diehle P, Kovács A, Duden T, Speen R, Žagar Soderžnik K and Dunin-Borkowski R E 2021 A cartridge-based turning specimen holder with wireless tilt angle measurement for magnetic induction mapping in the transmission electron microscope *Ultramicroscopy* **220** 113098
- [436] Pohl D *et al* 2023 Probing magnetic properties at the nanoscale: in-situ Hall measurements in a TEM (arXiv:2302.06127)
- [437] Möller M, Gaida J H, Schäfer S and Ropers C 2020 Few-nm tracking of current-driven magnetic vortex orbits using ultrafast Lorentz microscopy *Commun. Phys.* **3** 1–7
- [438] Zheng F *et al* 2021 Magnetic skyrmion braids *Nat. Commun.* **12** 1–7
- [439] Zheng F *et al* 2018 Experimental observation of chiral magnetic bobbles in B20-type FeGe *Nat. Nanotechnol.* **13** 451–5
- [440] Fernández-Pacheco A, Streubel R, Fruchart O, Hertel R, Fischer P and Cowburn R P 2017 Three-dimensional nanomagnetism *Nat. Commun.* **8** 15756
- [441] Volkov O M *et al* 2023 Chirality coupling in topological magnetic textures with multiple magnetochiral parameters *Nat. Commun.* **14** 1491